



**Study of the hydrodynamics and morphodynamics of the
Óbidos lagoon, Portugal**

Diogo Silva Mendes

Dissertation to obtain the Master of Science Degree in

Civil Engineering

(Hydraulics and Water Resources)

Supervisors: Dr. André Bustorff Fortunato and Prof. António Alberto Pires Silva

Examination Committee

Chairperson: Prof. António Alexandre Trigo Teixeira

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March 2015

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RESUMO

Na lagoa de Óbidos, as ondas e as correntes de maré induzem rápidas alterações morfológicas levando ao fecho da sua embocadura. As dragagens são frequentes tendo o objetivo de manter a laguna aberta e proteger as construções marginais. No entanto, o efeito das dragagens no comportamento da embocadura é um tema pouco estudado, especialmente o impacto que poderá ter a adição de canais transversais aos tradicionais canais dragados.

No interior desta laguna as ondas induzem alterações significativas nas condições hidrodinâmicas. Através da aplicação de um modelo acoplado constituído por um modelo hidrodinâmico e um modelo de agitação marítima concluiu-se que as elevações da superfície livre dentro da lagoa aumentaram para uma direcção média da onda perpendicular à linha de costa. Foi também investigado o impacto que a interacção completa entre ondas e correntes tem na dinâmica sedimentar da lagoa de Óbidos, concluindo-se que o transporte de sedimentos para o interior aumenta em 30% durante a enchente se esta interacção for tida em conta no modelo numérico.

Os três planos de dragagens foram avaliados através da implementação de um modelo morfodinâmico. Os novos planos de dragagens aumentaram o prisma de maré e reduziram as diferenças de duração entre a vazante e a enchente devido aos canais principais norte e sul. Os canais transversais aumentaram o assoreamento no canal principal sul e promoveram a estabilidade do canal principal norte. No entanto, devido às baixas profundidades e a uma agitação marítima muito energética, as soluções de dragagens são incapazes estabilizar permanentemente a lagoa de Óbidos.

PALAVRAS-CHAVE: Modelação morfodinâmica, Dragagens, Processos induzidos pelas ondas, Embocaduras de maré, Canais transversais, Profundidade de equilíbrio de dragagem

ABSTRACT

In the Óbidos lagoon, waves and tidal currents induce rapid morphological changes driving its inlet to closure. Dredging solutions are frequent in order to reposition the inlet in a central location and to deepen the main channels, therefore, preventing the inlet's closure. However, less is known about the effect of dredging solutions on the inlet's behaviour, especially the addition of transverse channels to the main dredged channels in dredging solutions.

Waves induce important changes in the hydrodynamics inside the Óbidos lagoon. The effect of mean wave direction on the water levels inside the lagoon was studied using a fully-coupled hydrodynamic and wave model. The water levels increased with the mean wave direction perpendicular to the coastline. The impact of the full interaction between waves and currents on the sediment dynamics in the Óbidos lagoon was also investigated. This study concludes that the sediment transport during the flood towards the lagoon increases 30% if the previous interaction was considered by the numerical model.

The three dredging plans were assessed by implementing and exploiting a morphodynamic model. The new dredging plans increased the tidal prism and reduced the differences between the ebb and flood durations mainly due to the southern and northern main channels. The transverse channels increased the accretion of the southern main channel and promoted the stability of the northern main channel. However, due to the shallow depths and the very energetic wave regime, dredging solutions are unable to permanently stabilize the Óbidos lagoon.

KEYWORDS: Morphodynamic model, Dredging, Wave-induced processes, Tidal inlets, Transverse channels, Dredging equilibrium depth

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LIST OF CONTENTS

LIST OF FIGURES.....	I
LIST OF TABLES.....	V
LIST OF SYMBOLS.....	VII
LIST OF ACRONYMS.....	IX
1. Introduction	1
1.1. The importance of tidal inlets: Óbidos lagoon inlet.....	1
1.2. Objectives and general methodology.....	2
1.3. Outline of the dissertation.....	2
2. Tidal inlets brief background review and the case study: the Óbidos lagoon	5
2.1. Tidal inlets	5
2.1.1. Inlet morphological components	5
2.1.2. Inlet classification and stability	6
2.1.3. Morphodynamic models for tidal inlets – some examples	9
2.2. Study site - the Óbidos lagoon.....	12
2.2.1. Geomorphological and hydrodynamic settings of the Óbidos lagoon	12
2.2.2. Long and short time scales of the morphosedimentary evolution.....	14
2.2.3. Coastal management and human interventions in the last decades.....	15
3. The influence of external forcing on the Óbidos lagoon hydrodynamics.....	19
3.1. Brief review on the effect of waves in the hydrodynamics of two Portuguese coastal lagoons.....	19
3.2. Methodology - a two-way coupling between SELFE and WWM-II.....	19
3.2.1. General scheme.....	19
3.2.2. The wave model – WWM-II.....	20
3.2.3. The hydrodynamic model - SELFE	21
3.2.4. Boundary conditions and numerical grid.....	21
3.2.5. Wave climate verification period and statistical error measures	23
3.2.6. Water levels verification period and statistical error measures	24
3.2.7. Idealized scenarios to study the influence of waves on the Óbidos lagoon	24
3.3. Analysis and discussion of the results	25
3.3.1. Comparisons between ADCP measurements and WWM-II simulations.....	25

3.3.2.	Comparisons between tidal gauge measurements and coupled SELFE and WWM-II simulations	27
3.3.3.	Sensitivity analysis of the wave parameters on the Óbidos lagoon hydrodynamics	27
3.3.4.	The wave blocking phenomenon	32
4.	Assessment of three dredging plans for the Óbidos lagoon.....	35
4.1.	Review on seasonal closure, tidal distortions and dredging solutions in wave-dominated inlets	35
4.2.	The morphodynamic model and the implemented methodology.....	37
4.2.1.	Morphodynamic model general outline	37
4.2.2.	The sediment transport and bottom evolution model – SED2D	38
4.2.3.	Sediment model verification period and statistical error measures	39
4.2.4.	Boundary conditions of the sediment model	40
4.2.5.	The simulated dredging plans under study	40
4.3.	The verification of the bathymetry simulated until November 2001	41
4.3.1.	The qualitative and quantitative morphodynamic model simulation	41
4.3.2.	The evolution of the M_2 harmonic constituent amplitude.....	43
4.4.	The impact of dredging plans on the inlet’s behaviour.....	44
4.4.1.	Qualitative evolution of the inlet	44
4.4.2.	Evolution of the tidal prism	45
4.4.3.	Differences in ebb-flood durations	46
4.4.4.	The Óbidos lagoon ability to flush out sediments	48
4.4.5.	The application of empirical stability relationships to the Óbidos lagoon inlet.....	50
4.5.	Summary	52
5.	Conclusions and future research	55
	Bibliographic references	57
	APPENDIX A - Governing equations of the numerical models	i
	APPENDIX B - Velocity fields during the ebb and the changed parameters in SELFE	v
	APPENDIX C - Modified boundary conditions in SED2D.....	vii
	APPENDIX D - Five-month evolution of the dredging plan proposed by Fortunato and Oliveira (2007a) (DP2).....	ix
	APPENDIX E - Five-month evolution of the dredging plan proposed by Fortunato and Oliveira (2007a) with only one transverse channel (DP3)	xi

APPENDIX F - Comparison between the modelled and measured on the wave gauge and WW3 wave parameters for each WW3 station (numbering from right to left and top to bottom).....	xiii
APPENDIX G – Asymptotically evolution of the bottom depth calculated for the three measured bathymetries available (July 2001, November 2001 and April 2002) at the flood delta	xv

LIST OF FIGURES

Figure 1.1 - The Óbidos lagoon social, environmental and commercial importance	1
Figure 2.1 - Morphological units of an inlet (adapted from Boothroyd, 1985)	5
Figure 2.2 - Hubbard <i>et al.</i> (1979) inlet classification (adapted from Hayes and FitzGerald, 2013)	7
Figure 2.3 - Hayes (1979) inlet classification (adapted from Hayes, 1979)	7
Figure 2.4 - Escoffier (1977) stability diagram for inlets (adapted from Escoffier, 1977)	8
Figure 2.5 - Conceptual morphodynamic models (adapted from FitzGerald <i>et al.</i> , 2000)	9
Figure 2.6 - Conceptual morphodynamic models (adapted from FitzGerald <i>et al.</i> , 2000)	10
Figure 2.7 - Morphodynamic numerical scheme (adapted from Bertin <i>et al.</i> , 2010).....	10
Figure 2.8 - A) Orthogonal structured grid; B) non-orthogonal curvilinear structured grid; C) orthogonal curvilinear structured grid; D) unstructured grid (adapted from Versteeg and Malalasekera, 2007)	11
Figure 2.9 – Geographical location of the Óbidos lagoon (green box) and its three main tributaries (A, B and C in the blue box) (adapted from <i>GoogleEarth</i>).....	12
Figure 2.10 – Distinct morphological zones in Óbidos lagoon: the upper (green box) and lower zones (blue box) (adapted from <i>GoogleEarth</i>)	13
Figure 2.11 - Óbidos lagoon evolution in the last 5000 years (adapted from Ferreira <i>et al.</i> 2009)	14
Figure 2.12 – Three possible configurations of the inlet (adapted from Bruneau <i>et al.</i> 2011)	15
Figure 2.13 – Zoom in on the Óbidos lagoon inlet of the area covered by the POOC (shaded zone) (adapted from Consulmar, 2008)	15
Figure 2.14 - Available Óbidos Lagoon morphological evolution since 1999 to present day (adapted from <i>Landstat</i>).....	16
Figure 2.15 - Solution proposed by Fortunato and Oliveira (2007)	17
Figure 3.1 - Numerical scheme for model verification (adapted from Dodet, 2013)	20
Figure 3.2 - A - Spatial values of the Manning coefficient ($m^{-1/3}s$) according to Bruneau <i>et</i> <i>al.</i> (2011); B - Extension of the Manning coefficient seawards.....	21
Figure 3.3 – Reference levels for the Óbidos lagoon	22
Figure 3.4 - Unstructured grid for the numerical model, bathymetry measured in July 2001, wave gauge (black star), tidal gauges (yellow stars), open boundary nodes (white crosses) and WW3 spectra points (white stars) (adapted from <i>Google</i> <i>Earth</i>)	23
Figure 3.5 - Locations of the recorded model stations 1, 2, 3 and 4 (red stars from left to right) for the sensitivity analysis, station for wave blocking phenomenon verification (black star), bathymetry measured in July 2001 and simulated scenarios to study the influence of different wave parameters inside Óbidos lagoon	25
Figure 3.6 - Comparison between the model (red line) and data (blue line) for the time period between 4 July 2001 and 21 October 2001 at the maritime station	26

Figure 3.7 - Model (red lines) and data (blue lines) SSE and MSL (dash lines) for Cais do Arelho (top), Braço da Barrosa (middle) and Bico dos Corvos (bottom) from 18/07/2001 to 06/08/2001 referred to MSL	27
Figure 3.8 - Top: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE for simulations with and without the full interaction between waves and currents for the four stations in Figure 3.5	28
Figure 3.9 - Top: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE when comparing with a simulation without waves	28
Figure 3.10 - Top: Hs at 5 m depth (black star in Figure 3.11); Middle: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE when comparing with a simulation without waves (Station 4)	29
Figure 3.11 - Hs for a JONSWAP spectrum with constant Hs (2,0 m) and Dir (315° N) as forcing parameters	30
Figure 3.12 - Top: Hs at 10 m depth (black star in Figure 3.13); Middle: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE when compared with a simulation without waves (Station 4).....	31
Figure 3.13 - Hs for a JONSWAP spectrum with constant Hs (2,0 m) and Tp (9 s) as forcing parameters	31
Figure 3.14 – Comparison between a simulation with the full wave-current interaction (red lines) and without (blue lines) for the Hs (first box), SSE (second box), depth-average horizontal velocities (third box), U_{rms} (fourth box) and sediment discharge (fifth box) between the third and fourth days on the black star (Figure 3.5).....	32
Figure 4.1 – Gradient of wave radiation stresses (left), barotropic pressure gradient forces (center) and resultant forces (right) at the Óbidos lagoon inlet for shore-normal waves of Hs = 3,0 m (adapted from Bertin <i>et al.</i> , 2009a).....	36
Figure 4.2 – Morphodynamic numerical scheme (adapted from Dodet, 2013).....	37
Figure 4.3 – Spatial distribution of the median grain diameter – d_{50} (m).....	39
Figure 4.4 - Comparison between the initial (left) and the changed (right) bathymetry on the northern margin.....	40
Figure 4.5 – A – Simulated bathymetry at the end of November; B – Measured bathymetry during November; C- Initial bathymetry in July 2001; Statistical error measures.....	42
Figure 4.6 –Morphodynamic evolution of the Óbidos lagoon between July and November 2001 for DP1.....	43

Figure 4.7 – Evolution of the M_2 amplitude in Cais do Arelho from July to November 2001 (adapted from Bertin <i>et al.</i> , 2009a)	44
Figure 4.8 – Qualitative comparison between the DP1 (left) and the DP2 (right) at the end of November 2001	44
Figure 4.9 - Qualitative comparison between the DP1 (left) and the DP3 (right) at the end of November 2001	45
Figure 4.10 – Tidal prism of DP1 (black line), DP2 (red line) and DP3 (blue line) from July to November 2001.....	45
Figure 4.11 – Ratio between the tidal prism of the DP2 (red lines) and the DP3 (blue lines) with the tidal prism of the DP1 and the linear regression analysis for DP2 and DP3 (strong lines) with the respective coefficients.....	46
Figure 4.12 - Difference in ebb-flood durations near Cais do Arelho for the DP1 (black line), the DP2 (red line) and the DP3 (blue line) from July to November 2001	47
Figure 4.13 - Difference in ebb-flood durations inside the lagoon (blue star in Figure 4.8) for the DP1 (black line), the DP2 (red line) and the DP3 (blue line) from July to November 2001.....	47
Figure 4.14 - Differences between the initial and 5 month simulated bathymetries for the DP1 (left) and the DP2 (right) and the sediment volumes that crossed the inlet ($\times 10^4 \text{ m}^3$).....	48
Figure 4.15 – Differences between the initial and 5 month simulated bathymetries for the DP1 (left) and the DP3 (right) and the sediment volumes that crossed the inlet ($\times 10^4 \text{ m}^3$).....	48
Figure 4.16 – Tidal prism-cross-section relationships for the three dredging plans	50
Figure 4.17 – Escoffier (1977) diagram for each dredging plan.....	51
Figure 5.1 – Conclusions scheme of chapter 3	55
Figure 5.2 – Conclusions scheme of chapter 4	56

LIST OF TABLES

Table 2.1 - Bruun and Gerritsen (1960) inlet stability criterion (adapted from Bruun and Gerritsen, 1960).....	8
Table 2.2 – Mean grain diameter (d_{50}) and gradation coefficient (σ_D) in the Óbidos lagoon (adapted from Fortunato <i>et al.</i> , 2011).....	13
Table 3.1 - Wave breaking type analysis based on the surf similarity parameter varying the T_p	29
Table 4.1 – Average differences between ebb and flood durations for two points inside the Óbidos lagoon.....	47
Table 4.2 - Differences between ebb and flood durations for the initial and final bathymetries and for each dredging plan.....	49
Table 4.3 – Ratios between the sediment fluxes integrated in half of the tidal cycle (ebb or flood) and in the inlet cross-section of DP2 and DP3 by DP1 for the initial and final bathymetries.....	50
Table 4.4 - Bruun and Gerritsen (1960) stability criterion application to the three dredging plans.....	52

LIST OF SYMBOLS

A_C	Area of the inlet cross-section	(m ²)
A_{mi}	Area of the grid element i	(m ²)
c	Wave celerity	(m/s)
C_b	Empirical coefficient of the linear JONSWAP parameterization	(m ² /s ³)
C_D	Drag coefficient computed from a Manning formulation	(-)
c_g	Wave group velocity	(m/s)
C_i	Coefficient of the node i for the finite element method formulation	(-)
c_p	Wave phase velocity	(m/s)
C_{PA}	Empirical coefficient of the PA relationship (O'Brien, 1966)	(-)
d_{50}	Median grain diameter	(m)
Dir	Mean wave direction	(°)
E	Variance density of the sea level elevations	(m ² /Hz)
f	Coriolis factor	(-)
h	Water depth	(m)
H	Total water depth	(m)
H_0	Offshore wave height	(m)
H_s	Significant wave height	(m)
k	Wave number	(rad/s)
L	Wavelength	(m)
L_0	Offshore wavelength	(m)
m	Coordinate perpendicular to the wave propagation direction	(-)
M_e	Mobility parameter (Van Rijn, 2007)	(-)
M_{tot}	Gross annual littoral drift	m ³ /year
N	Wave action density	(m ² /(radHz ²))
n	Empirical coefficient of the PA relationship (O'Brien, 1966)	(-)
n_0	Manning coefficient ($n_0=1/K_S$)	(m ^{-1/3} s)
P	Tidal prism	(m ³)
q	Depth-integrated volumetric sediment flux	(m ³ /s/m)
Q	Sediment transport rate integrated in time	(m ³ /m)
Q^*	Modified sediment transport rate integrated in time	(m ³ /m)
q_b	Bedload sediment transport (Van Rijn, 2007a)	(kg/m/s)
q_s	Suspended sediment transport (Van Rijn, 2007b)	(kg/m/s)
q_{tot}	Total sediment transport (Soulsby, 1997)	(kg/m/s)
Q_x	Sediment fluxes along x-direction	(m ³ /m/s)
Q_y	Sediment fluxes along y-direction	(m ³ /m/s)
s	Coordinate along the wave propagation direction	(-)
S_{ij}	Component of the radiation stress tensor in plan i along j direction	(N/m ² /m)

T	Wave period	(s)
T_p	Wave peak period	(s)
u	Depth-average horizontal velocity along the x-axis	(m/s)
U_A	Effective advection velocity	(m/s)
$u_{cr,c}$	Critical velocity for currents	(m/s)
$u_{cr,w}$	Critical velocity for waves	(m/s)
u_e	Effective velocity	(m/s)
U_i	Average flow velocity at the inlet	(m/s)
U_{rms}	Wave orbital velocity	(m/s)
U_w	Peak orbital velocity	(m/s)
v	Depth-average horizontal velocity along the y-axis	(m/s)
V_m	Maximum velocity on the transitional channel (Escoffier, 1977)	(m/s)
V_{tot}	Net balance of sediments through the inlet over a specific time period	(m ³)
V_{tot_neg}	Total volume of sediments eroded inside the lagoon	(m ³)
V_{tot_pos}	Total volume of sediments accreted inside the lagoon	(m ³)
\mathbf{X}	Cartesian coordinate vector (x,y)	(m)

Greek symbols

β	Beach slope for the Iribaren number	(-)
β_{SVR97}	Bed slope (Soulsby, 1997)	(-)
γ	Breaker height to water depth ratio	(-)
Γ_i	Length of the boundary element i	(m)
Δh	Bed update	(m)
ε	Diffusion coefficient for the modified Exner equation	(-)
η	Sea surface elevation	(m)
θ	Wave direction	(°)
λ	Porosity	(-)
μ	Horizontal eddy viscosity	(m ² /s)
ξ_0	Iribaren number or surf similarity parameter	(-)
ρ_0	Water specific mass	(kg/m ³)
σ	Wave relative frequency	(rad/s)
σ_D	Gradation coefficient	(-)
τ_b	Bottom shear stress	(N/m ²)
ω	Wave absolute frequency	(rad/s)
Ω_i	Control volume of the element i	(m ³)

LIST OF ACRONYMS

ADCP	Acoustic Doppler current profiler
BSS	Brier-skill score
DAHV	Depth-averaged horizontal velocities
DP1	Dredging plan proposed by DHI (1997)
DP2	Dredging plan proposed by Fortunato and Oliveira (2007)
DP3	Dredging plan proposed by Fortunato and Oliveira (2007) with only one transverse channel
FIBWC	Full interaction between waves and currents
HZ	Hydrographic zero
ME	Mean error
MSL	Mean sea level
RMSE	Root-mean square error
SI	Scatter index
Skill	Skill error measure
SSE	Sea surface elevations

1. Introduction

SUMMARY: This chapter introduces the importance of tidal inlets, with an increased focus on the social, economic and environmental importance of the Óbidos lagoon. Next, the objectives and the general methodology of this dissertation are drawn and this chapter ends with the dissertation outline.

1.1. The importance of tidal inlets: Óbidos lagoon inlet

The economic and environmental importance of tidal inlets has been growing worldwide. According to a United Nations study, 60% of the world population lives within 60 km of the coast and this proportion will rise to 75% within two decades (United Nations, 2013). Several hazards, such as coastal erosion and sea-level rise are causing an extreme pressure on coastal environments and this provokes an increasing interest for coastal sciences.

The Óbidos lagoon is a coastal lagoon located in the Portuguese western coast and has a great economic, social and environmental importance for the neighbouring regions (Figure 1.1). The influence of waves and currents is evidenced by the morphodynamic changes on the Óbidos lagoon inlet, which occur in monthly or even weekly time scales.



Figure 1.1 - The Óbidos lagoon social, environmental and commercial importance

As the Óbidos lagoon is a small coastal system, its inlet is mainly driven by the sediment transport induced by waves. Due to its fragile equilibrium, along the years, the inlet frequently closes. Dredging solutions are used to reposition the inlet to a more stable position and also to deepen the lagoon's main channels, thereby improving the lagoon's capability to flush out sediments.

Small coastal lagoons, like the Óbidos lagoon, could develop the nearby regions. These coastal lagoons provide tourism and leisure activities, improve the regional markets through navigation routes that could be established, and also contribute to the wildlife by renewable the water mass. However, to keep these coastal systems alive the central government spend substantial financial resources every year, especially in dredging operations. Overall, this dissertation aims at contributing to the coastal management decisions

that are most needed, with a major focus on the hydrodynamics and morphodynamics of the Óbidos lagoon.

1.2. Objectives and general methodology

This dissertation has the following objectives:

- 1) To further investigate the effect of waves on the Óbidos lagoon hydrodynamics.
- 2) To verify the changes induced by the interactions between waves and currents on the sediments of the Óbidos lagoon inlet.
- 3) To assess the morphodynamic changes induced by three different dredging plans on the hydrodynamic behaviour and on the morphological evolution of the Óbidos lagoon.
- 4) To evaluate the impact of adding transverse channels in dredging plans.

To achieve the previous goals the following methodology is carried out:

- 1) A brief review of tidal inlets regarding inlet's morphological units, classifications, stability and conceptual or numerical morphodynamic models to provide the required theoretical background.
- 2) The presentation of the case under study - Óbidos lagoon – with a focus on tidal conditions, wave regime, sediment characteristics, long and short term evolution and previous implemented engineering solutions.
- 3) The validation of a coupled wave-circulation model with the data recorded by tidal gauges and an acoustic Doppler current profiler (ADCP).
- 4) An analysis of the wave parameters effect on the Óbidos lagoon hydrodynamics and the importance of the full interaction between wave and currents in the numerical model to sediment dynamics through idealized wave conditions.
- 5) The calibration and verification of the morphodynamic model with the measured bathymetry.
- 6) The assessment of three different dredging plans during a period of five months with the morphodynamic model.
- 7) The calculation of several parameters, such as tidal prism, differences between ebb and flood durations, sediment fluxes that crossed the inlet, application of empirical stability relationships and differences between initial and final bathymetries.

1.3. Outline of the dissertation

This dissertation is divided into five chapters including this Introduction. In chapter 2 a background review is made. A number of concepts and definitions regarding tidal inlets are introduced like the tidal prism, an inlet's classifications and the role of stability. Associated with previous studies, conceptual and numerical morphodynamic models are also described. Furthermore, the study area – the Óbidos lagoon - is introduced and a full description of the past solutions that were implemented or studied in the last 20 years for the Óbidos lagoon is presented.

Chapter 3 studies the hydrodynamics of the Óbidos lagoon. In this chapter, the coupling between the hydrodynamic and the wave model is verified and further exploited to understand the impact of the mean wave direction on the sea surface elevations (SSE) inside the Óbidos lagoon. The wave blocking phenomenon is also shown in this chapter.

The full morphodynamic model is verified in chapter 4. This verification involves a comparison with measurements, and the good accuracy makes the morphodynamic model an appropriate tool to study the evolution of the Óbidos lagoon inlet with different dredging plans and also, to investigate the impact of transverse channels on the morphodynamics of the Óbidos lagoon inlet. Finally, chapter 5 presents the main conclusions and some future work directions for subsequent studies are proposed.

2. Tidal inlets brief background review and the case study: the Óbidos lagoon

SUMMARY: This chapter introduces the theoretical background needed to perform the study of tidal inlets. The morphological components of tidal inlets are introduced, a review on inlet's classifications and stability is made and some conceptual/numerical morphodynamic models are presented. Next, the Óbidos lagoon relevant characteristics are described, such as, the wave regime, tidal conditions and the median grain size spatial location. This chapter closes with a review on the engineering solutions implemented and studied for this lagoon during the last twenty years.

2.1. Tidal inlets

2.1.1. Inlet morphological components

Tidal currents and waves shape the sediment deposits that constitute the morphological units of an inlet. A tidal inlet corresponds to the transitional channel that connects an estuary or a coastal lagoon to the sea. Hayes *et al.* (1970) distinguished three main morphological units that compose a tidal inlet: the ebb-tidal delta, the inlet throat and the flood-tidal delta (Figure 2.1).

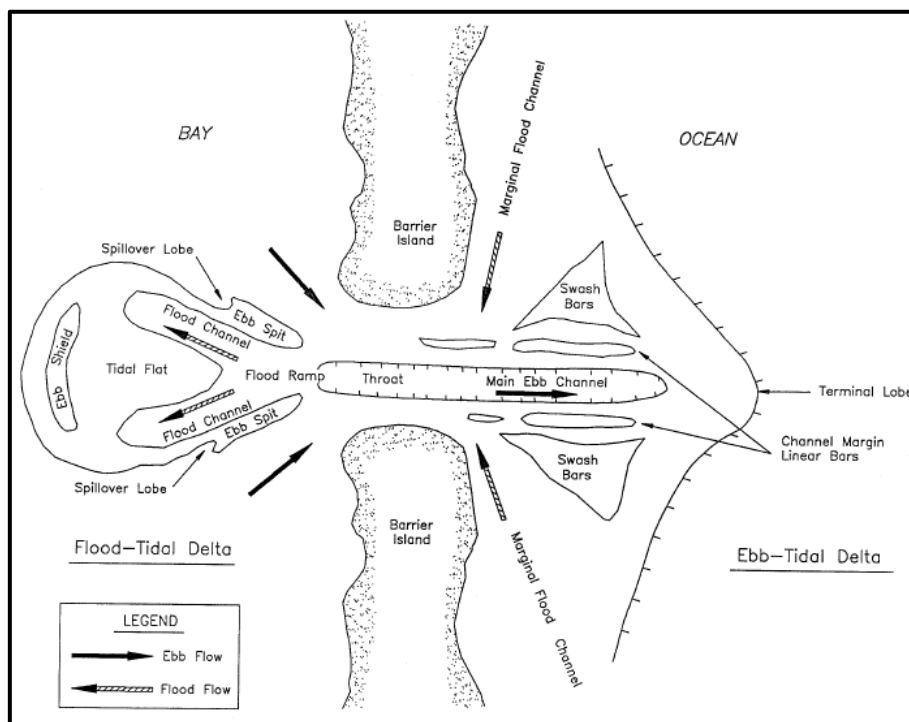


Figure 2.1 - Morphological units of an inlet (adapted from Boothroyd, 1985)

The barrier islands usually migrate in the direction of the dominant transport. Their growth and deposition rate are key factors in order to determine their lateral displacement or even the closure of the transitional channel (Hayes and FitzGerald, 2013).

According to Morang and Parson (2002), the ebb-delta sediments come from the erosion in the transitional channel and the longshore sediment transport. The ebb-tidal delta is composed by: an ebb channel, channel margin linear bars, a terminal lobe, swash bars and flood channels. Ebb currents shape the ebb channel. The interactions between these currents and waves originate the margin linear bars.

When water leaves the ebb channel, the velocity decreases at the end of the terminal lobe and the deposition of suspended sediments occurs. Wave-induced currents originate the movements of swash bars between the terminal lobe and the ebb channel.

The flood-tidal delta is composed by four elements: flood ramp, flood channels, ebb shield and ebb spits. Flood currents generate the flood ramp and the flood channels. The ebb shield induces the divergence of these currents at the upstream end of the flood delta. The effect of both flood and ebb currents across the margins generates ebb spits.

The bedforms found in the inlets are ripples, dunes and antidunes. The oscillatory motion induced by waves in the inlets promotes the appearance of ripples, while dunes and antidunes are due to tidal currents.

To study the tidal inlet morphodynamics it is necessary to introduce the concept of tidal prism. This concept corresponds to the volume of water that crosses an inlet during a flood or ebb cycle (Equation 2.1)

$$P = \int_{t_1}^{t_2} U_i A_c dt \quad [2.1]$$

where P is the tidal prism and U_i is the average flow velocity at the inlet. A_c is the area of the inlet cross-section and t_1 and t_2 are the temporal limits of the half tidal cycle period. Since the inlets have a dynamic behaviour, changing in space and time, their tidal prism is also variable. Therefore, the tidal prism depends on the inlet geometry, the wave forcing, the tidal amplitude and the area of the lagoon (O'Brien, 1966).

According to O'Brien (1966), the tidal prism of a spring tide is roughly proportional to the minimum cross-section of the inlet. Walton and Adams (1976) show that there is a relationship between the tidal prism of an inlet and the size of the ebb delta. FitzGerald (1988) stated that the volume of sand in the flood and ebb deltas is comparable with the volume of sand in the adjacent barrier islands. According to Nummedal and Fischer (1978), the geometry of the inlet entrance and its deltas depends on four factors: the tidal amplitude, the wave energy along the coast, the bathymetry of the sand bar and the area of the coastal lagoon.

2.1.2. Inlet classification and stability

A review of the specialized literature on tidal inlets shows that there are different classifications for this feature. Hubbard *et al.* (1979) proposed one qualitative classification. These authors define three types of inlets (Figure 2.2) characterized by their morphological components:

- 1) tide-dominated inlets
- 2) wave-dominated inlets
- 3) transitional inlets

A tide-dominated inlet is, usually, wide and does not migrate. It exhibits a deep ebb-dominant main channel flanked by long margin linear bars and undeveloped flood deltas. A wave-dominated inlet is small and a very dynamic one. It is characterized by shallow tidal channels, large flood deltas and small ebb deltas. Finally, the transitional inlets display their undistinguishable deltas along the transitional

channel. However, the previous empirical classification disregards the hydrodynamic forces which inlets are subjected to.

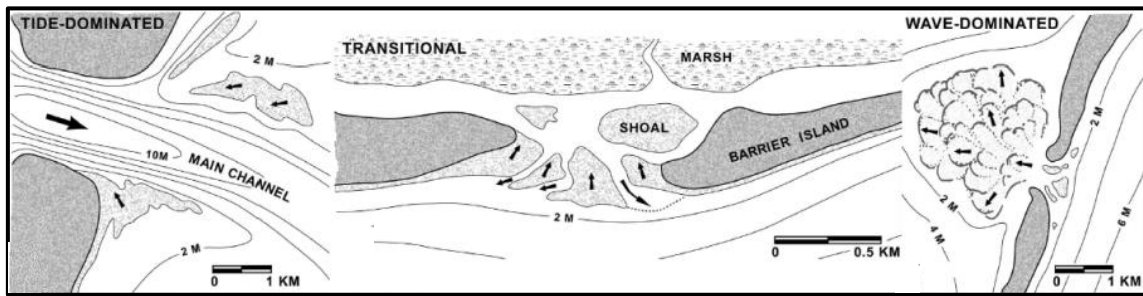


Figure 2.2 - Hubbard *et al.* (1979) inlet classification (adapted from Hayes and FitzGerald, 2013)

One of the most commonly used classifications worldwide is the Hayes (1979) classification. This author considers five categories of inlets based on the average tide range and the mean wave height (Figure 2.3).

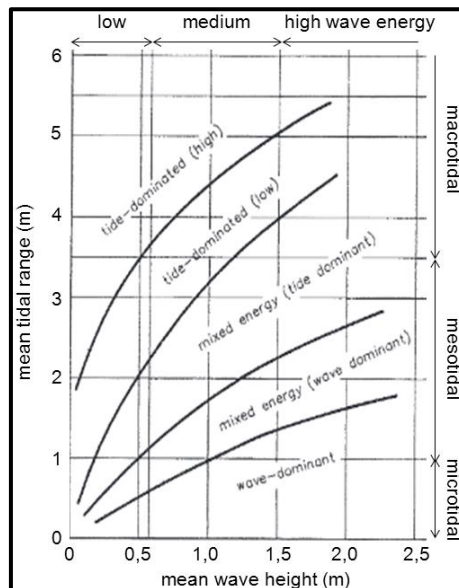


Figure 2.3 - Hayes (1979) inlet classification (adapted from Hayes, 1979)

The type of inlets present in a particular region is estimated through tidal amplitude and wave height maps. However, Nahon *et al.* (2012) claimed that Hayes (1979) classification does not take into account the tidal prism or the longshore sediment transport.

The stability of tidal inlets is controlled by the relative importance of waves and tides (Bruun, 1978). An inlet is considered stable if its location remains unchanged over a certain period and the sediments influx and outflux, through the transitional channel, are similar (Dean, 2002). O'Brien (1966) proposed an empirical relationship between the spring tidal prism and the minimum inlet cross-section (Equation 2.2). This equation became known as the PA relationship.

$$A_C = C_{PA} P^n \quad [2.2]$$

where C_{PA} and n are empirical coefficients. Jarret (1976) studied more than one hundred inlets on the Atlantic and Pacific coasts of the United States and in the Gulf of Mexico. From that study, this author proposed a value for C between $7,76 \times 10^{-6}$ and $5,02 \times 10^{-4}$, and between 0,84 and 1,1 for n . The PA

relationship has been demonstrated empirically and theoretically for tidal inlets in equilibrium. However, Fortunato *et al.* (2014) have shown that the relationship is also valid for inlets that are away from equilibrium.

Bruun and Gerritsen (1960) introduced the concept of tidal inlet global stability based on the ratio of the tidal prism at spring tides and the gross annual littoral drift (M_{tot}) (Table 2.1). Nevertheless, this criterion cannot be used to predict the evolution of an inlet with time-scales equal or lower to a year. According to Dodet (2013) the reason for this lies in the M_{tot} , since it presents a strong inter-annual variability.

Table 2.1 - Bruun and Gerritsen (1960) inlet stability criterion (adapted from Bruun and Gerritsen, 1960)

<i>Inlet Stability Ratings</i>	
$P/M_{tot} \geq 150$	Conditions are relatively good, little bar and good flushing
$100 \leq P/M_{tot} \leq 150$	Conditions become less satisfactory, and offshore bar formation becomes more pronounced
$50 \leq P/M_{tot} \leq 100$	Entrance bar may be rather large, but there is usually a channel through the bar
$20 \leq P/M_{tot} \leq 50$	All inlets are typical “bar-bypassers”
$P/M_{tot} \leq 20$	Descriptive of cases where the entrances become unstable “overflow channels” rather than permanent inlets

Escoffier (1977) considered that the stability of an inlet depends on the balance between settling and erosion forces. This idea was developed analytically and relates the maximum velocity on the transitional channel, V_m , with the area of the inlet cross-section, A_c (Figure 2.4).

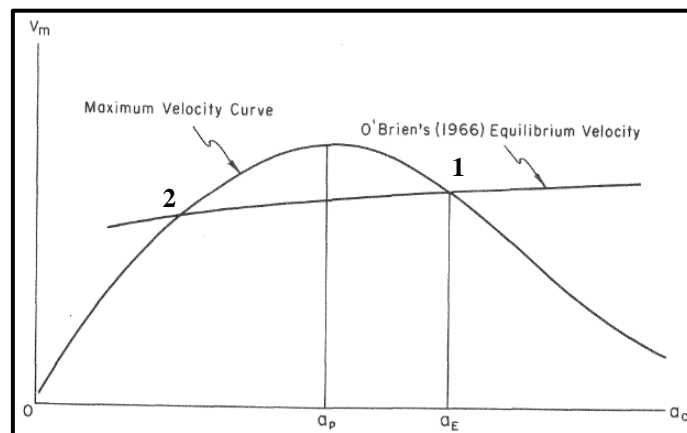


Figure 2.4 - Escoffier (1977) stability diagram for inlets (adapted from Escoffier, 1977)

The tails in the above diagram represent two distinct situations. The maximum velocity tends to zero and the area of the inlet cross-section decreases (left tail) or the maximum velocity decreases as the cross-section area increases (right tail). The first situation is mainly due to friction forces from the bottom and side walls and the second one is the result of mass conservation for an arbitrary tidal prism.

There are two equilibrium points (1 and 2 in Figure 2.4). Point 1 is considered stable because if the area increases the maximum velocity decreases and sediments deposition occurs until the equilibrium state is reached. On the contrary, if the area decreases the maximum velocity will increase eroding sediments, and driving the inlet back to the equilibrium situation. Point 2 is considered unstable because a decrease of the area reduces the velocity inducing the inlet closure.

2.1.3. Morphodynamic models for tidal inlets – some examples

The first models which study tidal inlets were conceptual and based on observations. Brunn and Gerritsen (1959) describe three mechanisms of sand bypassing through an inlet. The transport induced by waves through the terminal lobe, the transport in the channels due to tidal currents and the transitional channel or barrier islands migration.

FitzGerald *et al.* (2000) synthesize the conceptual models for sediment transport and morphological changes for natural or stabilized inlets (Figure 2.5 and Figure 2.6).

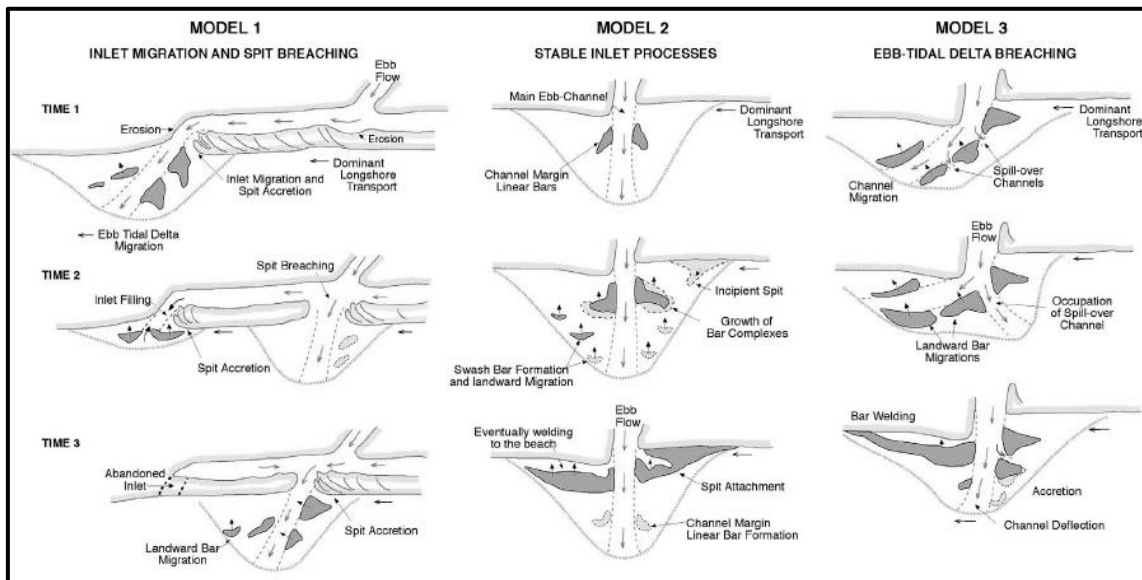


Figure 2.5 - Conceptual morphodynamic models (adapted from FitzGerald *et al.*, 2000)

Model 1 (Figure 2.5) corresponds to a spit formation with the downdrift direction. Its growth will narrow the transitional channel, increasing the flow velocity and the erosion areas. Fragile areas will appear, and the breach creates a new location for the inlet. This model represents the inlet migration, the spit breaching and the downdrift sand bypassing.

Model 2 represents an inlet on stable morphodynamic equilibrium. The landward migration of swash bars will increase the velocity in the transitional channel. This phenomenon promotes erosion and sediments settling in the ebb delta during the ebb tide. This model presents a behavior similar to the one identified by point 1 in Escoffier diagram (Figure 2.4).

The ebb delta breaching (model 3 in Figure 2.5) needs a stable transitional channel. Swash bars are formed in the ebb delta due to longshore sediment transport. Wave-induced currents will move these bars landwards, narrowing the ebb channel. This narrowing causes an overtopping phenomenon, responsible for the ebb delta breaching.

Model 4 (Figure 2.6) is similar to model 3 but mobilizing smaller volumes of sand. The linear bar, in the inner margin of the channel, remains fixed, while the outer margin is deflected due to sediments accumulation. The outside margin breaching occurs during spring tides. Flood-tidal and wave-induced currents move the swash bars downdrift (sand bypassing).

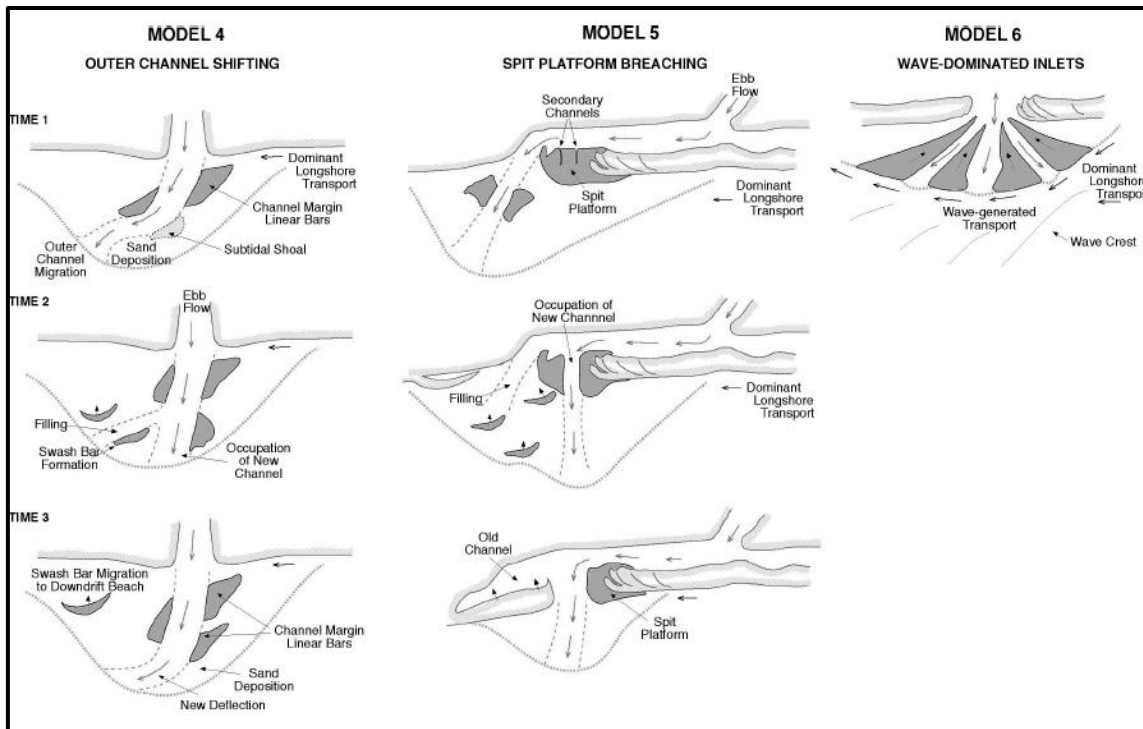


Figure 2.6 - Conceptual morphodynamic models (adapted from FitzGerald *et al.*, 2000)

Model 5 presents the sand bypassing through the breached sandbank. Longshore sediment transport forms this sandbank. Ebb currents create secondary channels in the fragile areas of the spit. This mechanism of sand bypassing is similar to model 1. However, the breaching does not correspond to an extension of the transitional channel.

In model 6, the ebb-delta morphology is largely controlled by waves. These inlets display shallow transitional channels, and due to shallow water and arched geometry of their ebb deltas the waves are the most relevant agent of sediment transport. This transport mechanism increases during floods.

Recently, with the improvements of computers calculations capabilities, morphodynamic numerical models started to be used. The first morphodynamic numerical models are from the 1980's (Bertin *et al.*, 2010). These models simulate the flow, the sediment transport and the bottom evolution. However, their application was restricted to rivers. The initial developments of numerical morphodynamic models for inlets began only in the 1990's. These models are composed of a circulation model, a wave model and a sediment transport and bottom evolution model (Figure 2.7).

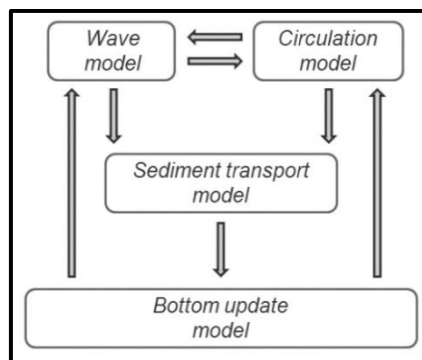


Figure 2.7 - Morphodynamic numerical scheme (adapted from Bertin *et al.*, 2010)

The grid is a fundamental component of numerical model and represents the spatial discretization of the domain in analysis. The grids can be structured or unstructured (Figure 2.8). The grid used in this dissertation is classified as unstructured. Unstructured grids are able to adapt to complex geometries in a better way providing more flexibility for different resolutions.

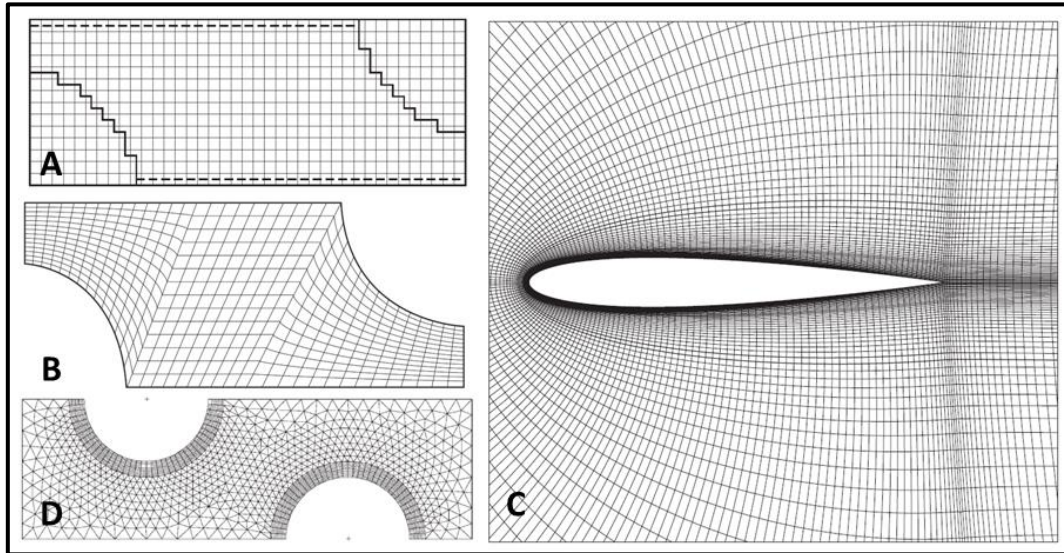


Figure 2.8 - A) Orthogonal structured grid; B) non-orthogonal curvilinear structured grid; C) orthogonal curvilinear structured grid; D) unstructured grid (adapted from Versteeg and Malalasekera, 2007)

Various morphodynamic models have been developed since the 1990's. Mike-21 (Warren and Bach, 1992), Delft3D (Lesser *et al.*, 2004), Morsys2D (Fortunato and Oliveira, 2004a; Bertin *et al.*, 2009b), Telemac&Sisyphe (Villaret *et al.*, 2013) are examples of these developments. All these models can simulate the morphodynamics of coastal zones.

Ranasinghe and Pattiaratchi (1999) applied a morphodynamic model to the Wilson inlet, in Australia. The objective was to understand the physical processes responsible for the closing of the inlet and the landward sediment transport constitutes the main mechanism that leads to the inlet's closure. These authors did not take into account the wave-current interaction due to the coupling incapability between the models.

Cayocca (2001) studied the effect of tides and waves on the long-term morphodynamic changes in the Arcachon inlet, France. This author was able to reproduce the opening of a new channel due to the tide. Long-term morphodynamic simulations forced this author to use a representative tide and wave. The representative tide and wave were obtained based on the annual sediment transport induced by tidal and wave conditions.

Bertin *et al.* (2009a) used the Morsys2D model (Fortunato and Oliveira, 2004a; Bertin *et al.*, 2009b) to simulate the morphodynamics of the Óbidos lagoon inlet. The aim of this study was to understand the relevant physical processes that occur near the inlet. The computational cost forced these authors to use a 2D model.

Villaret *et al.* (2013) coupled Telemac-2D and 3D (the hydrodynamic model) with Sisyphe (the sediment transport and bottom evolution model) and applied them on curved channels, recirculating flows, sand

grading effects in a waterway and wave-induced littoral drift. These authors state that the main limitations are due to the high degree of empiricism inherent in most sediment transport models.

Dodet (2013) applied a morphodynamic model (same as in this dissertation) to the Albufeira Lagoon. This lagoon has an annual cyclic behavior and, in this study, the inlet was artificially opened in April and closed in December. This author shows the importance of two sand transport mechanisms: the swash transport and the infra-gravity wave transport. Until then, morphodynamic models neglected these two mechanisms.

2.2. Study site - the Óbidos lagoon

2.2.1. Geomorphological and hydrodynamic settings of the Óbidos lagoon

The Óbidos Lagoon is a small and shallow coastal system located on the western coast of Portugal (Figure 2.9). The continental shelf in front of it is very narrow and its width ranges from 20 to 30 km (Bertin *et al.*, 2009a). The shore-normal direction is oriented at 315° N in clockwise nautical convention. This lagoon is connected to the sea by a small inlet (less than 100 m wide) in which the morphological changes occur on monthly time scales. This inlet cuts a sandy littoral barrier, which is delimited by two headlands: Cruz do Facho at the northern coast and Rocha do Gronho at the southern coast (Fortunato *et al.*, 2011).



Figure 2.9 – Geographical location of the Óbidos lagoon (green box) and its three main tributaries (A, B and C in the blue box) (adapted from GoogleEarth)

The Óbidos Lagoon drainage basin has 440 km² of area and a surface area of 4,4 km² at the mean sea level (MSL), reaching 6 km² at high tide. Its depth, width and maximum length are 4,5 m, 1800 m and 4,5 km, respectively (Freitas *et al.*, 1992). The average temperature in the Óbidos lagoon is 15° C and the rainfall intensity is between 600 to 650 mm (Freitas, 1989).

The Óbidos lagoon has three main tributaries, Vala do Ameal, Arnóia and Cal (A, B and C in Figure 2.9). However, the freshwater input is generally small, with average flows of 3 m³/s (Vão, 1991). These average flows represent less than 5% of the average tidal prism scaled by the M₂ period (Rego, 2004). The salinity values reach more than 25 in the middle of the lagoon and less than 20 along the rivers Cal and Arnóia and on the northern shore of the lagoon (Freitas, 1989).

According to Fortunato and Portela (2014), the Óbidos lagoon can be divided into four zones due to its distinct morphological behaviours (Figure 2.10).

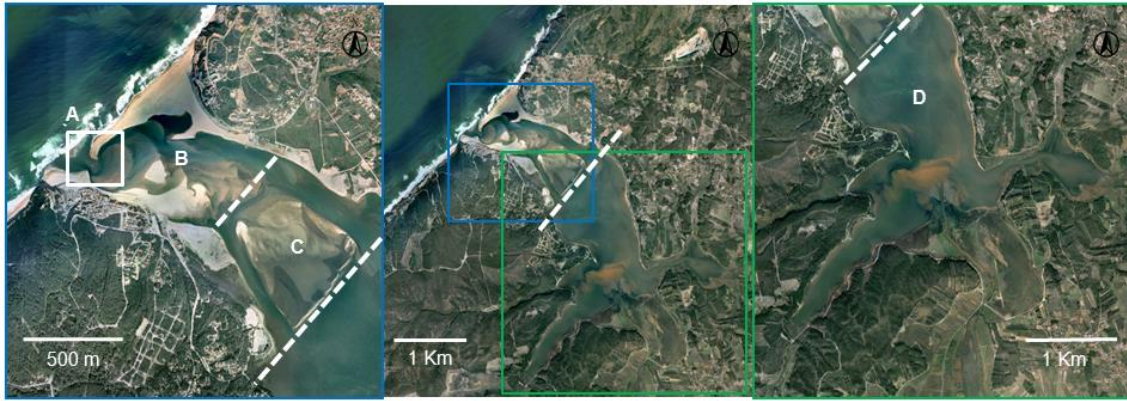


Figure 2.10 – Distinct morphological zones in Óbidos lagoon: the upper (green box) and lower zones (blue box) (adapted from GoogleEarth)

Zone D (Figure 2.10) is the upper region and is composed by a main basin and the three main tributaries. This area is characterized by low velocities, clay bottoms and a tendency for silting (Henriques, 1997).

Zone C, the middle zone, consists of a sandbank and two channels: the southern channel and the northern channel. This zone presents a relatively morphological stability over the past few decades, with a small tendency for silting (Fortunato and Portela, 2014). Flood currents transport sediments towards the lagoon. Furthermore, these sediments accumulated in the sand bank cannot be washed away by the ebb currents. The sediment accumulations in the sandbank and in the lagoon's northern and southern channels (zone C in Figure 2.10) are the reason for the small tendency for silting (Fortunato and Oliveira, 2007).

Dynamic banks, separated by tidal channels, constitute the lower area of the Óbidos lagoon - zone B (Figure 2.10). The length of this area is variable because it depends on the closeness of the inlet to the shores (north and south). The inlet (zone A) is the most dynamic zone of Óbidos lagoon. It displays movements on monthly and weekly time scales.

To characterize the sediments, two distinct zones were taken into account: the upper zone (D) and the lower zone (A, B and C). The median grain diameter (d_{50}) and the gradation coefficient (σ_D) were measured between 26 and 30 August 2010 by a company NEMUS under the supervision of LNEC (Table 2.2).

Table 2.2 – Mean grain diameter (d_{50}) and gradation coefficient (σ_D) in the Óbidos lagoon (adapted from Fortunato *et al.*, 2011)

	<i>Upper zone</i>	<i>Lower zone</i>
d_{50} (mm)	< 0,4	0,4 – 0,8
σ_D	2,0 – 2,2	1,5 – 1,9

The ocean tides are semidiurnal with amplitudes ranging between 2 to 4 m offshore and 1 to 2 m inside the lagoon (Oliveira *et al.* 2006). The most frequent mean wave direction (Dir) is 315° N (nautical convention) and wave peak periods (T_p) vary between 5 and 20 s. Bertin *et al.* (2009a) estimated a significant wave height (H_s) in the maritime summer (from April to September) between 1,5 and 2,5 m and between 2,0 to 6,0 m during maritime winter (from October to March) with an analysis of the offshore wave conditions based on the WaveWatch3 model (hereafter referred to as WW3, Tolman *et al.*,

2002 and further developed by Dodet *et al.*, 2010). Bruneau *et al.* (2011) based on the WW3 outputs also estimated an extreme offshore wave climate and concluded that the wave height exceeds 2,5m 20% of the time. According to the classification defined by Hayes (1979), the Óbidos lagoon inlet can be characterized as a mixed-inlet dominated by tides in the maritime summer and a mixed-inlet dominated by waves in the maritime winter. However, several authors characterized this inlet as a wave-dominated inlet (Oliveira *et al.*, 2006; Bertin *et al.*, 2009a; Bruneau *et al.*, 2011; Bertin *et al.*, 2015) because the wave-induced processes are very intense in the Óbidos lagoon inlet.

2.2.2. Long and short time scales of the morphosedimentary evolution

Ferreira *et al.* (2009) studied the Óbidos lagoon morphological evolution from the Holocene period until the present. The methodology was based on archaeological data, historical documents, cartographic maps, aerial photographs and military maps. The Óbidos lagoon was originated due to the MSL rise, which flooded numerous depressions during the Holocene period (Figure 2.11). After this period, the lagoon suffered a massive sediment accumulation which leads to its current position (Dinis *et al.*, 2006; Freitas *et al.*, 2009).

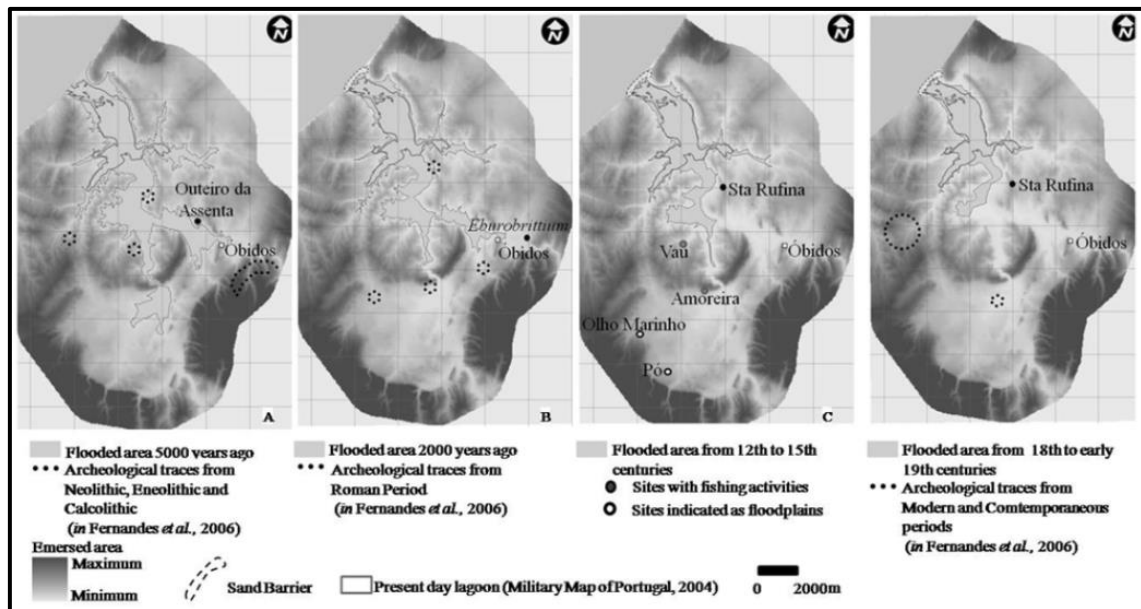


Figure 2.11 - Óbidos lagoon evolution in the last 5000 years (adapted from Ferreira *et al.* 2009)

Dinis *et al.* (2006) concluded that the significant morphosedimentary changes occurred in the Óbidos lagoon are due to anthropogenic interventions since the Medieval Age. Human activities, as intensive agriculture on the margins, influenced by small historical climate changes in the last millennium promoted the sediment accretion in the Óbidos lagoon.

Freitas *et al.* (1992) conducted a quantitative analysis on sediment deposition and erosion in the Óbidos lagoon based on hydrographic surveys between 1917 and 1980. The value of the total deposition was $4,3 \times 10^4$ m³/year, and the erosion was $4,0 \times 10^3$ m³/year. The global average sedimentation rate is 7 mm/year and the erosion is 0,6 mm/year. The same authors state that the Óbidos lagoon can store the vast majority of fluvial sediments, which constitute the largest sediment source of the lagoon.

Fortunato *et al.* (2011) compared the bathymetries of the Óbidos lagoon between 2000 and 2004 in the upper zone (zone D). The comparison area of the previous study was $2,6 \times 10^6 \text{ m}^2$ and these authors concluded that the average sedimentation rate was 49 mm/year. This enhances the increased sedimentation behaviour in the Óbidos lagoon along the recent years.

In what concerns the stability of the Óbidos lagoon inlet, Bruneau *et al.* (2011) studied the morphodynamic evolution of the inlet during one year with a morphodynamic model. The previous study concluded that morphologies B and C are less unstable configurations (point 1 in Figure 2.4) and morphology A is a transitional state (Figure 2.12). The previous authors also concluded that an increase of the sea level will induce a global sedimentation of the lagoon. This way, the Óbidos lagoon also displays sedimentation behaviour in an idealized sea-level rise scenario.

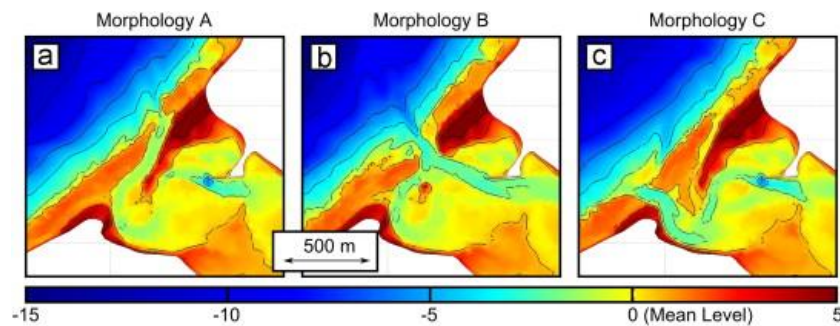


Figure 2.12 – Three possible configurations of the inlet (adapted from Bruneau *et al.* 2011)

2.2.3. Coastal management and human interventions in the last decades

The POOC – *Planos de Ordenamento da Orla Costeira* are the Portuguese legal instruments for coastal management and policy. These plans divide the Portuguese coast into ten sectors based on sedimentary cells and physiological characteristics. The study area of this dissertation is included in the Alcobaça-Mafra plan.

Beach plans (PP-Planos de Praia) are the building blocks of the POOC. The PP contain information on the urbanization and accesses (coastal management), the cliffs and shoaling risk lines (coastal dynamics) and the placement and disposition of the bathing areas and beach constructions (planning). Nevertheless, the area covered by the POOC (shadow zone in Figure 2.13) does not include the inner part of the Óbidos lagoon and so, the control of this lagoon is made by the municipalities.

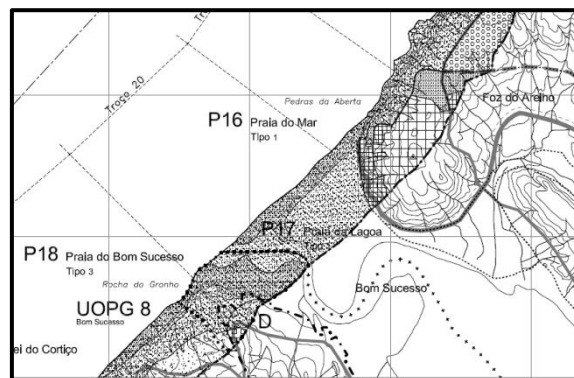


Figure 2.13 – Zoom in on the Óbidos lagoon inlet of the area covered by the POOC (shaded zone) (adapted from Consulmar, 2008)

From a coastal management perspective, structural coastal defense works are inexistent. Human occupation density along the shores is smaller than the ones near villages (Foz do Arelho and Nadadouro). An analysis of aerial photographs shows that agriculture is the most common use of the soil (Figure 2.10). This area also contains housing, tourist and leisure activities close to the lagoon area. Dunes with underbrush characterize a large part of the soil near the inlet and a large forest is present along the Óbidos lagoon, especially in the southern area (Figure 2.10). Rocky cliffs are found on the northern coast while sands constitute the inlet and the southern coast.

Anthropogenic interventions in Óbidos lagoon date back to the Middle Age. Freitas *et al.* (1992) referred that this lagoon was a cause of concern throughout the years, and often the inlet is opened to prevent agricultural lands from flooding since the 15th century.

The solutions implemented in the Óbidos lagoon have the following objectives (Oliveira *et al.*, 2006):

- 1) prevent the closure of the inlet transitional channel to allow the exchange of water between the coastal lagoon and the sea
- 2) protect the margin banks
- 3) stabilize the channels in the lower region.

Dredging was the most common engineering solution (red stars in Figure 2.14), used in 1995, 1998, 1999, 2001, 2003 and 2011-2012 (Fortunato and Portela, 2014). These dredging operations were not sufficient to prevent the transitional channel migration. In the winter of 1993/1994, this movement caused some damages to the buildings on the northern shore of the lagoon (Fortunato and Oliveira, 2007).

The first structural solution consists of a sheet-pile curtain construction near the northern shore of the lagoon in 1999. The inlet migration to the southern margin led to the implementation of emergency solutions, with sandbag placements (green stars in Figure 2.14) in 1998-1999, 2001, 2002 and 2004 (Fortunato and Portela, 2014). These sandbags had to be constantly replaced due to erosion in their foundations caused by tidal currents (Fortunato and Oliveira, 2007).

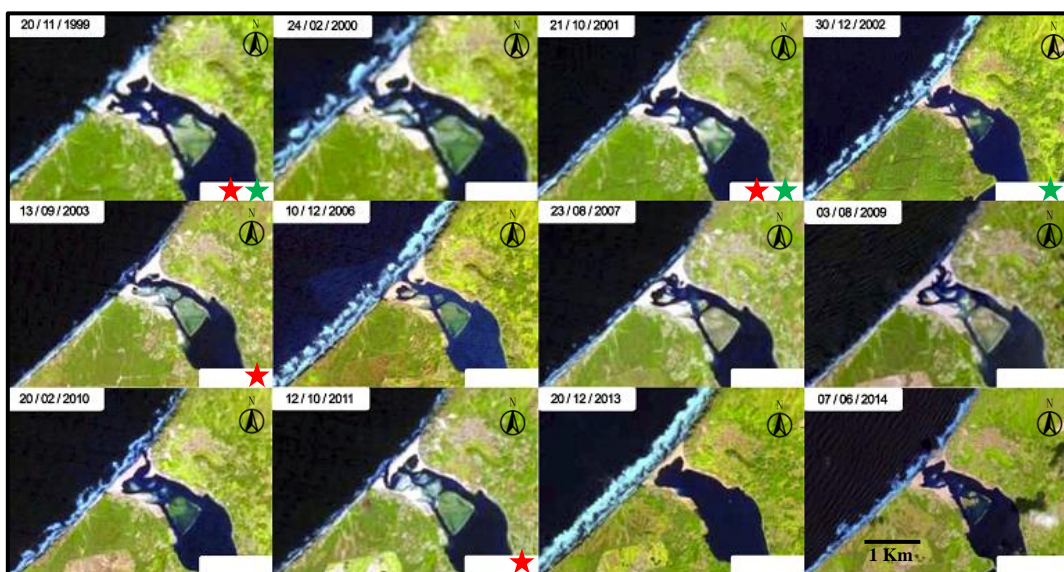


Figure 2.14 - Available Óbidos Lagoon morphological evolution since 1999 to present day (adapted from Landstat)

The high maintenance cost of these emergency solutions led to the development of studies of permanent solutions to ensure the inlet stability. Hidrotécnica Portuguesa (1991) proposed a traditional solution, of two parallel jetties perpendicular to the shoreline. This solution aims to reduce dredging operations in the lagoon. However, this solution was abandoned because the flood tide would transport water with a higher sediment concentration towards the lagoon, and the jetties solution had a high visual impact (Fortunato and Oliveira, 2007).

Theoretically, the increase of the lagoon surface area will enhance the ebb sediment transport capacity and increase the lagoon tidal prism. These concepts were the basis of the solution proposed by Vieira (2001) which would imply a large dredging operation plan ($15 \times 10^6 \text{ m}^3$). Although this solution promoted the ebb dominance, it did not guarantee the inlet stability (Fortunato *et al.*, 2002).

Fortunato and Oliveira (2007) proposed a third solution for the Óbidos lagoon (Figure 2.15). The solution comprehends a dredging plan with two main channels and seven transverse channels and a partially submerged guiding wall. The purpose of the transverse channels is to decrease the flood dominance, reducing therefore, the ebb period, optimizing the tidal asymmetry which improves the ability of the lagoon to flush out incoming sediments and to remove sediments from the sand bank (blue box in Figure 2.10).



Figure 2.15 - Solution proposed by Fortunato and Oliveira (2007)

According to Fortunato and Oliveira (2007), the main channels have a width of 30 m (southern channel) and between 70 to 100 m (northern channel). The lateral slopes of the cross-sections are 1:10 and they must be dredged to chart datum. The transverse channels have a width of 25 m and a depth of 2 m. The combined length of the seven transverse channels is 3 km and they are placed with an angle of 45° with the main channels to maximize the residual ebb sediment fluxes (Fortunato and Oliveira, 2004b). This dredging solution of two main channels and seven transverse channels represents a dredging volume of $5,1 \times 10^5 \text{ m}^3$ compared to the dredging plan proposed by DHI (1997) which corresponds to the initial bathymetry of July 2001.

The submerged guiding wall, with 300 m long and parallel to the northern main channel, improves the transitional channel stability and prevents the southward migration of the inlet. Another advantage of this solution is that this structure has a reduced visual impact when compared with traditional jetties.

In 2009 and 2010 it was necessary to relocate the inlet mouth and in January 2014, due to the high frequency of large waves in the winter, the inlet closed and the Portuguese Environment Agency (*Agência Portuguesa do Ambiente – APA*) proceeded to its reopening. Along all these years, several engineering solutions have been proposed and implemented. Yet, these solutions could not prevent the inlet closure at the beginning of 2014, for that reason, new engineering solutions are needed in order to stabilize the inlet and prevent the lagoon accretion during the winter period, thereby, preventing the inlet's closure.

3. The influence of external forcing on the Óbidos lagoon hydrodynamics

SUMMARY: This chapter concerns the effect of waves on the sea surface elevations and on the sediment dynamics in the Óbidos lagoon. To study the effect induced by waves, a fully coupled hydrodynamic and wave model was used in order to simulate the water circulation taking into account the wave-current interaction. The results of the following simulations are reported: sea surface elevations inside the lagoon, significant wave heights in front of the inlet and sediment discharge at the beginning of the flood delta. From these results, it was found that sea surface elevations increased for the north-west mean wave direction. Also, the consideration of the full interaction between waves and currents in the numerical model increased the sediment transport towards the lagoon during the flood up to 30%.

3.1. Brief review on the effect of waves in the hydrodynamics of two Portuguese coastal lagoons

Atmospheric (wind and rain) and marine (tides and waves) forcings cause rapid changes in the physical characteristics of coastal lagoons (Troussellier and Gattuso, 2007). The presence of waves can influence the sea-level variations in a quite significant way (Nielsen and Apelt, 2003) and this rise on the sea level due to waves is called “wave set-up”. Longuet-Higgins and Stewart (1960) relate this effect to the concept of radiation stresses, which are defined as the excess of flow momentum due to the presence of waves.

It is known that the radiation stresses increase the MSL outside the Óbidos lagoon and this increase outside propagates further inside (Malhadas *et al.*, 2009, Bertin *et al.*, 2009a). Also, a higher H_s will increase the SSE on the Óbidos lagoon transitional channel, due to the dependency of the radiation stresses on this variable. However, in what concerns T_p and Dir , the precise nature of their impact on the SSE in the Óbidos lagoon remains an open issue.

Coastal inlets, such as Óbidos lagoon inlet, induce the lag and damping of the tidal signal and the differences in the MSL, inside and outside the lagoon, produce tidal currents. The interaction between tidal currents and waves leads to complex current fields at the inlet entrance and adjacent beaches. Several wave experiments (Kemp and Simons, 1982, 1983; Klopman, 1994; Umeyama, 2005) and numerical simulations (Groeneweg and Klopman, 1998; Olabarrieta *et al.*, 2010; Teles *et al.*, 2013) were carried out to investigate the evolution of wave characteristics under the influence of currents.

Dodet *et al.* (2013) studied the wave-current interactions in the inlet of the Albufeira lagoon and experienced the wave blocking phenomenon. The wave blocking phenomenon is the blocking of wave-induced currents by the ebb currents. These authors found that if the wave model disregards wave-currents interaction the H_s will not drop quickly inside the lagoon as they verified in data. Also, the disregarding of this interaction attenuates the seaward sediment fluxes during the ebb. Therefore, this chapter also aims to investigate the wave blocking phenomenon at the Óbidos lagoon.

3.2. Methodology - a two-way coupling between SELFE and WWM-II

3.2.1. General scheme

The numerical model used in this chapter simulates the hydrodynamics of coastal areas taking into account the wave-current interactions (Figure 3.1). Ocean waves and water levels force the coupled

model at its boundaries. First, the wave model calculates the radiation stresses (S_{xx} , S_{yy} and S_{xy}), which will be the inputs for the circulation model. Then, the circulation model obtains the sea surface elevations (η) and the depth-averaged horizontal velocities (u and v). In the following time step, the wave model is updated with the SSE and the currents from the circulation model and this iterative cycle continues until an instant specified by the user.

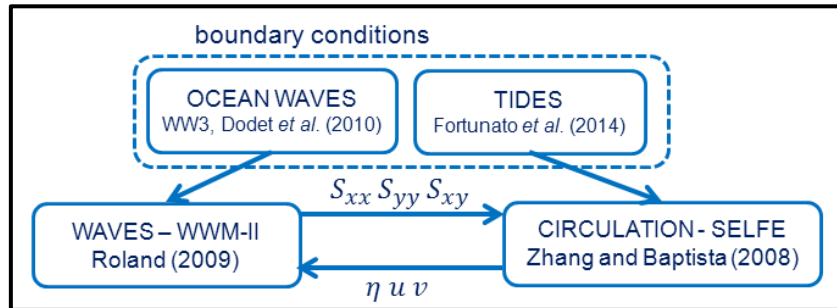


Figure 3.1 - Numerical scheme for model verification (adapted from Dodet, 2013)

In this dissertation, the two models will be verified against an ADCP and tidal gauges measurements recorded by the Portuguese Hydrographic Institute (IH, 2001) using the scheme of Figure 3.1. The following step corresponds to a sensitivity analysis on the wave characteristics and for that reason, the waves must be forced with a JONSWAP spectrum instead of the WW3 wave spectra. The wave blocking phenomenon will be also addressed with a JONSWAP spectrum in order to specify constant wave parameters (H_s , T_p and Dir) on the sea boundary.

3.2.2. The wave model – WWM-II

The spectral phase-averaged wave model WWM-II (*Wind Wave Model II* – Roland, 2009) solves the wave action density balance equation and was coupled with SELFE in quasi-steady mode to propagate the waves from the ocean to the coast. An explanation of the equations solved by WWM-II is presented in APPENDIX A.

This wave model uses an unstructured grid and has a parallel processing system. This feature renders it suitable to be coupled with SELFE, avoiding interpolation algorithms between structured and unstructured meshes (e.g., Bertin *et al.*, 2009a; Bruneau *et al.*, 2011). There are two coupling options between the models: the fully-coupled and partial coupled. The latter option causes the currents from SELFE to be disregarded by WWM-II.

The spectral discretization was done according to 12 frequencies ranging from 0,2 and 0,05 Hz. The directional discretization was done according to 12 directions in a window of 180° (0° - SW and 180° - NE) with a resolution of 15°. WWM-II was set to take into account the bottom friction based on the empirical JONSWAP parametrization (Hasselmann *et al.*, 1973) with the empirical coefficient C_b of 0,067 m^2s^{-3} , the triad wave-wave interaction, the full interaction between waves and currents (fully-coupled) and the dissipation by wave breaking with a constant wave breaking coefficient γ of 0,78.

In this dissertation the deep-water source-terms mechanisms (four-wave interaction, white capping and wave generation by wind forcing) were disregarded because the domain of this study area is small (10 km). These simplifications reduced the computational time by approximately 20 to 30%.

Significant wave height, period, direction, wavelength and orbital velocity outputted from WWM-II were used to compute gradients of radiation stresses to force the hydrodynamic model (SELFE). The radiation stresses were computed from the classical formulation of Longuet-Higgings and Stewart (1964) with a time step of 5 min, which is in accordance with Dodet (2013).

3.2.3. The hydrodynamic model - SELFE

The circulation model SELFE (*Semi-implicit Eulerian-Lagrangian Finite Element model* – Zhang and Baptista, 2008) uses an unstructured grid in a parallel processing system. The unstructured grid allows this numerical model to be suitable for estuaries and tidal inlets, ensuring flexibility for different scales and resolutions. The configuration used in this dissertation was a 2D barotropic model with the hydrostatic assumption and the Boussinesq approximation in Cartesian coordinates. SELFE solves the 2D shallow water equations through a semi-implicit time stepping scheme and uses a finite-element method for the spatial discretization. A review of the equations solved by SELFE is presented in APPENDIX A.

To account for bottom friction, a Manning coefficient (n_0) variable in space was used in SELFE based on the previous study of Bruneau *et al.* (2011) with a minor extension seawards (Figure 3.2) in order to reduce the large velocity vectors that appeared during the ebb (APPENDIX B). The Manning formulation threshold depth (variable H in Equation 17 - APPENDIX A) was set to 0,10 m. The water depth threshold for drying was set on 0,01 m, which is in accordance with Dodet *et al.* (2013). To account for horizontal turbulence, a constant horizontal eddy viscosity μ of $1 \text{ m}^2\text{s}^{-1}$ was set in the entire domain. The hydrodynamic time step was 20 s, producing outputs of sea surface elevations and velocities, saved every 30 min over the domain.

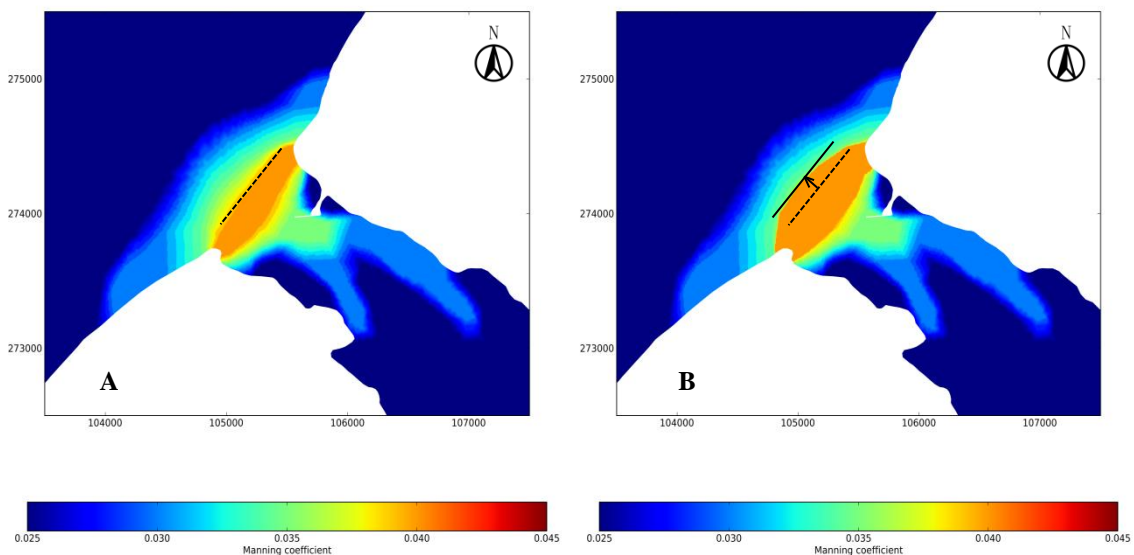


Figure 3.2 - A - Spatial values of the Manning coefficient ($\text{m}^{-1/3}\text{s}$) according to Bruneau *et al.* (2011); B - Extension of the Manning coefficient seawards

3.2.4. Boundary conditions and numerical grid

The hydrodynamic model was forced at its ocean boundary (white crosses in Figure 3.4) by 20 tidal constituents whose amplitude and phase were computed with the regional model of Fortunato *et al.*

(2014). The atmospheric pressure remained constant (1015 HPa with a standard deviation of 3 HPa) along the study period and the atmospheric pressure variations were disregarded (Bertin *et al.*, 2009a).

The upstream boundary in the lagoon was defined as a closed boundary since freshwater inflow is negligible (section 2.2.1). The MSL increased over the last decades. In Cascais tidal gauge it can be observed an increase of approximately 0,15 m relative to the hydrographic zero (HZ) since 1938, thereby, the amplitude of the harmonic constituent M_0 was set to 0,15 m. Finally, because the bathymetry was related to the MSL and data was recorded to the HZ, 2 m were subtracted to the data measurements in order to evaluate the results (Figure 3.3).

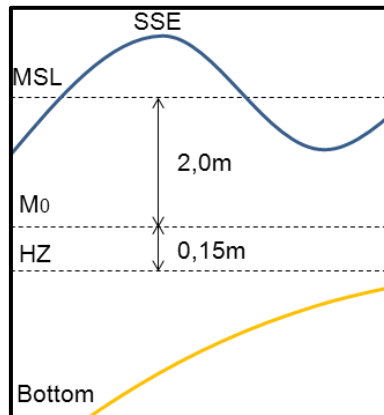


Figure 3.3 – Reference levels for the Óbidos lagoon

The offshore forcing was obtained from the previous study of Bruneau *et al.* (2011), provided by the regional application of the WW3 model which gives the 2D wave spectrum in each WW3 spectra point (white stars in Figure 3.4) with a time step of 3 h. The wave spectra are then linearly interpolated between the spectra points and the ocean boundary nodes (white crosses in Figure 3.4).

The unstructured grid used by both models (Figure 3.4) had 57880 elements and 29398 nodes that were set by previous studies (Bruneau *et al.*, 2011). The spatial resolution of the grid ranges from 1,6 km (open boundary) to 6 m (inlet).

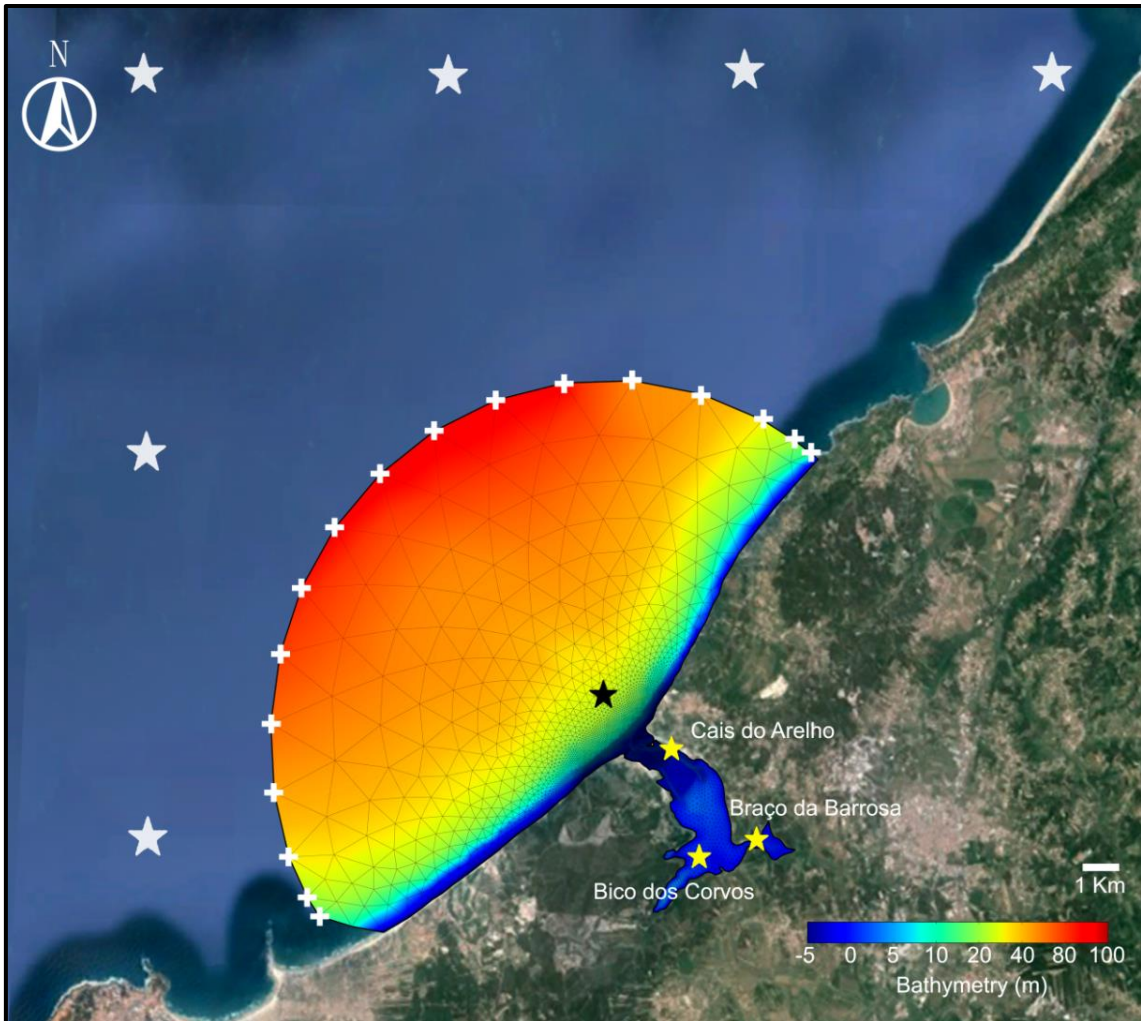


Figure 3.4 - Unstructured grid for the numerical model, bathymetry measured in July 2001, wave gauge (black star), tidal gauges (yellow stars), open boundary nodes (white crosses) and WW3 spectra points (white stars) (adapted from Google Earth)

3.2.5. Wave climate verification period and statistical error measures

The wave's verification period was from 4 July 2001 to 21 October 2001 (approximately three and a half months). The data recorded at the wave gauge (black star in Figure 3.4) were H_s , T_p and Dir . Wave parameters were obtained through an acoustic profiler located in the maritime station and placed at a bathymetric depth between 20 and 30 m (HZ). This acoustic Doppler profiler is a RDI Workhorse 600 kHz ADCP equipped with the module WAVES where the sample period was 10 minutes (IH, 2001). There are a few gaps in the time series of these three parameters in the four months of simulations. The gaps correspond to 5 hours in the last day of each month (July, August and September). As these interruptions are of the order of hours, a linear interpolation was performed. The recorded Dir also need to account with the magnetic and the geographic difference. In the report (IH, 2001), it does not refer that data were adjusted to this phenomenon. For the year of 2001 the magnetic difference in this lagoon was -4° W, and so, 4° were added in the measured Dir .

To evaluate the wave model performance, a comparison between the data, recorded at the buoy, and the model was made. It were chosen the following three error measures (Teles *et al.*, 2012): the root-mean

square error (RMSE) (Equation 3.1), the mean error (ME) (Equation 3.2) and the scatter index (SI) (Equation 3.3).

$$RMSE = \sqrt{\frac{1}{N} \sum_{i=1}^N (\Delta X_i)^2} \quad [3.1]$$

$$ME = \frac{1}{N} \sum_{i=1}^N \Delta X_i \quad [3.2]$$

$$SI = \frac{RMSE}{\bar{X}_o} \quad [3.3]$$

where $\Delta X_i = X_m - X_o$ is the difference between the modelled and the observed values, and \bar{X}_o is the mean value of the observed values.

3.2.6. Water levels verification period and statistical error measures

To verify the water levels inside Óbidos lagoon the simulated period was from 18 July 2001 to 6 August 2001. It represents the closest period to the dredging operations (May/June 2001) that has SSE records (IH, 2001) for the three tidal gauges (yellow stars in Figure 3.4). These tidal gauges have a sample period of 6 min.

The error measures used in this dissertation to evaluate the SSE were the root-mean-square error (RMSE) (Equation 3.4), the Skill error measure (Dias *et al.*, 2009) (Equation 3.5) and the ME (Equation 3.2). These are common statistical parameters to evaluate models behaviour (Bennett *et al.*, 2013). If the RMSE range between 5% to 10% of the measured SSE the model shows a very good agreement with the data (Dias *et al.*, 2009). Skill ranges between 0 and 1. The value 1 will be the perfect accordance between the model and the data. The ME assesses under or overestimation of any parameter.

$$RMSE = \sqrt{\frac{1}{n} \sum_{i=1}^n (y_i - \hat{y}_i)^2} \quad [3.4]$$

$$Skill = 1 - \frac{\sum_{i=1}^n (y_i - \bar{y})^2}{\sum_{i=1}^n (|y_i - \bar{y}| + |\hat{y}_i - \bar{y}|)^2} \quad [3.5]$$

where y_i are the observed values, \hat{y}_i the modelled values, \bar{y} the mean of the observed values and $\bar{\hat{y}}$ the mean of the modelled values.

3.2.7. Idealized scenarios to study the influence of waves on the Óbidos lagoon

A sensitivity analysis of the wave parameters was carried out after the model's verification. The analysis was based on a set of scenarios (Figure 3.5). The chosen values were based on the wave parameters measured in July. In this month the H_s had a mean value of 1,3 m, the T_p assume values between 4,5 to 12,5 s and Dir was confined to a sector defined by 327° N and 277° N.

To study the influence of the wave parameters on the SSE of the Óbidos lagoon four recording stations were chosen. Station 1 (Figure 3.5) recorded the SSE without the lag and damping effects caused by the inlet. Station 2 determined the increase on SSE caused by the wave setup. Station 3 verified the behaviour suggested by Bertin *et al.* (2009a), which corresponds to an increase of the SSE in the transitional channel for a higher H_s and so could act to validate the procedure for this methodology. Station 4 had the objective of understanding the differences in the SSE inside the lagoon caused by the change of the wave parameters.

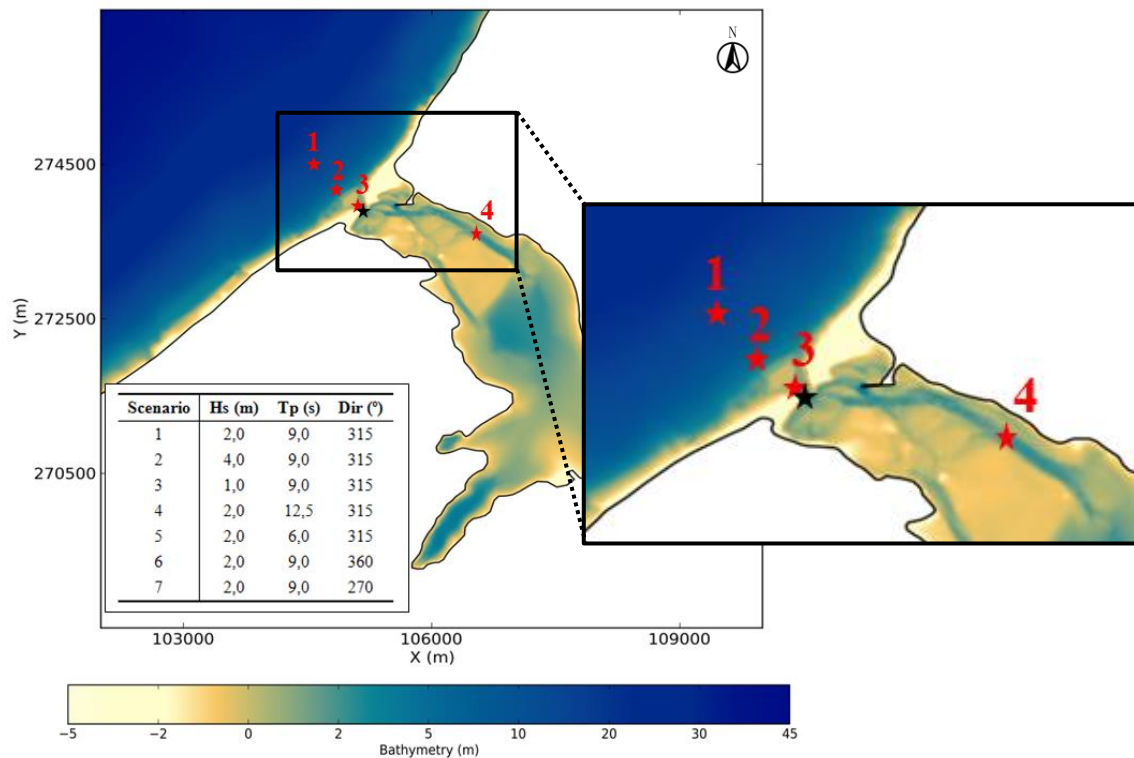


Figure 3.5 - Locations of the recorded model stations 1, 2, 3 and 4 (red stars from left to right) for the sensitivity analysis, station for wave blocking phenomenon verification (black star), bathymetry measured in July 2001 and simulated scenarios to study the influence of different wave parameters inside Óbidos lagoon

This study was conducted over twenty days and was divided into two phases. The first phase corresponded to spring tide while the second was associated to a neap tide. The bathymetry available represented the one measured in July 2001. The forcing agents of these models were the harmonic constituents of the synthetic tide (same as section 3.2.4) and the irregular JONSWAP spectrum. The Manning coefficient used was the same to the one on the verification step (section 3.2.3). A simulation without waves was performed in order to compare it to the simulated scenarios.

Finally, in order to study the wave blocking phenomenon in the Óbidos lagoon two simulations of five days were done using both the fully-coupled option and the partial coupled option. Except for the wave parameters, the values of the forcing parameters were the same as in the previous sensitivity analysis. The Hs had the value of 2,0 m, the Tp was 9 s and the Dir was set on 315° N. According to Dodet *et al.* (2013), the best location to record the wave blocking phenomenon was next to the flood delta (black star in Figure 3.5). The sediment discharge was computed with the Soulsby - Van Rijn formula (Soulsby, 1997). The specific sediment density was specified of 2,65, the d_{50} took the value of 1 mm and the bed slope (β_{SVR97}) was set in 0,1. Sea surface elevations, depth-averaged velocities and wave orbital velocities were obtain through the application of the coupled SELFE and WWM-II.

3.3. Analysis and discussion of the results

3.3.1. Comparisons between ADCP measurements and WWM-II simulations

Two comparison periods and the correspondent statistical errors are presented in Figure 3.6. The first period was from 4 July until 21 October (solid line boxes) and the second was from 4 of July until 30

September (dash line boxes). The modelled Hs was well reproduced by the model until October. In October, the peak of Hs was reproduced but some lags had appeared between the modelled and measured Hs. However, in what concerns the morphodynamic changes, the energy presented in the measured peak of Hs was simulated and these lags are of minor importance for the morphodynamic simulation since the morphological changes occurs in larger time scales than days.

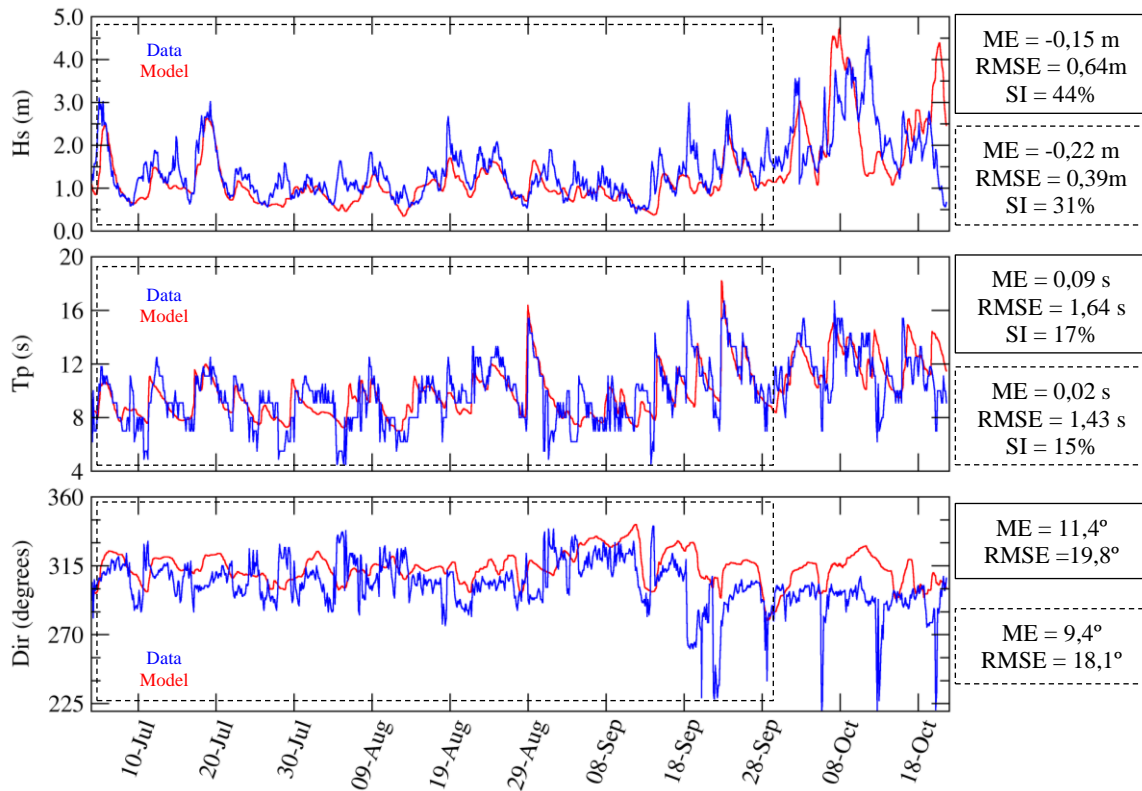


Figure 3.6 - Comparison between the model (red line) and data (blue line) for the time period between 4 July 2001 and 21 October 2001 at the maritime station

The modelled Tp displayed a very good statistical error measures. However, the wave model was not able to represent Tp less than 8 s during the verification period because the high frequency band (from 0,2 to 1 Hz) was truncated. In what concerns the Dir, the statistical errors can be partially attributed to the course discretization, therefore, the drastic changes in Dir during the last month were not simulated by the model. After a careful analysis on the WW3 spectra points, it can be concluded that the errors in Dir and Hs are due to the wave forcing at the open boundary while the errors in Tp are due to the wave model (APPENDIX F).

Dodet *et al.* (2013) and Teles *et al.* (2012) presented similar values for the statistical errors of Tp and Dir. However, the previous mentioned studies obtained a value for the Hs RMSE around 0,30 m and a SI of 15%. Previous simulations performed in the Óbidos lagoon also exhibited better values. Bertin *et al.* (2009a) were able to show the following statistical errors of 0,35 m for the Hs RMSE, 0,5 s for Tp and 5 ° for Dir during the period between 28 April 2001 and 6 June 2001. Bruneau *et al.* (2011) got a SI of: 18% for Hs and 14% for Tp. Concerning Dir, Bruneau *et al.* (2011) obtained a RMSE and a ME of 7° and -1,5° in the winter period, respectively.

3.3.2. Comparisons between tidal gauge measurements and coupled SELFE and WWM-II simulations

The model reproduces well the tidal SSE for two out of the three tidal gauge stations (Figure 3.7). In the third station, the data MSL (blue dashed lines) is below the data MSL in the other two stations. These results show that the reference level of this station (Bico dos Corvos) is incorrect.

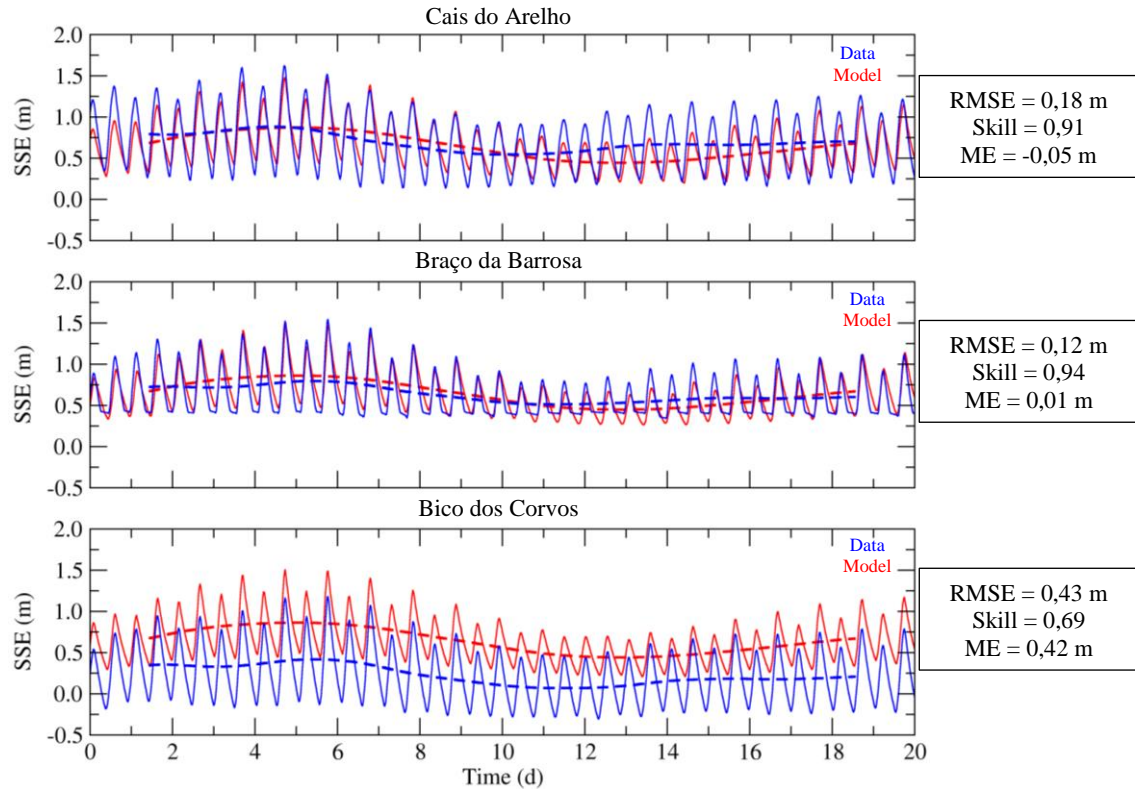


Figure 3.7 - Model (red lines) and data (blue lines) SSE and MSL (dash lines) for Cais do Arelho (top), Braço da Barrosa (middle) and Bico dos Corvos (bottom) from 18/07/2001 to 06/08/2001 referred to MSL

The statistical errors for the tidal gauge located at Bico dos Corvos were large. The reference level error in Bico dos Corvos tidal gauge could be due to previous settlement of the tidal gauge because the report (IH, 2001) states that occur some settling problems in the tidal gauges during the field campaign. Furthermore, this tidal gauge was located in the inner part of the lagoon and in this region the grid was coarser, therefore, the numerical errors start to increase.

Bertin *et al.* (2009a) obtained a RMSE of 0,05 m between 10/07/2001 and 30/07/2001 for the Cais do Arelho tidal gauge. The Manning coefficient used by those authors varied between 0,025 and 0,035 $m^{-1/3}s$. This reduction in the Manning coefficient could enlarge the tidal amplitude inside the lagoon and lead to better results in this period. However, in the winter period, the morphodynamic simulation is unfeasible due to large velocity vectors that increase the sediment transport and lead to spurious morphodynamic changes.

3.3.3. Sensitivity analysis of the wave parameters on the Óbidos lagoon hydrodynamics

The effect of the wave-currents interaction on the tidal SSE inside the Óbidos lagoon was investigated. The differences in SSE due to currents are insignificant for all the stations except Station 2. This station is located where wave breaking is the most intense. These differences can be explained by numerical

oscillations of the model (Figure 3.8). Since currents are irrelevant for the SSE inside the lagoon analyses, the partial coupling was used due to the computational costs (approximately 20% faster than fully-coupled).

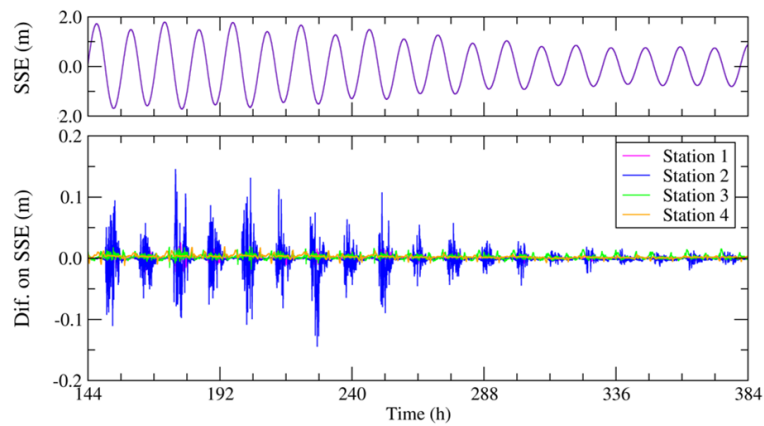


Figure 3.8 - Top: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE for simulations with and without the full interaction between waves and currents for the four stations in Figure 3.5

The analysis of scenarios 1, 2 and 3 in Figure 3.5 allows concluding that the SSE inside the lagoon increased for higher H_s (Figure 3.9). The radiation stresses are the explanation for the previous fact because when the wave height increases, more momentum will be transferred to the breaking zone and the SSE in that zone will increase. First, this phenomenon takes place outside (breaking zone) and then inside the lagoon, which is in accordance to Bertin *et al.* (2009a).

During spring tide the differences in SSE are smaller than at neap tide (Figure 3.9). The generation of even constituents inside the lagoon, such as MS_f , is responsible for the differences in SSE between spring and neap tides because when the water depth is higher (spring tide), the finite amplitude term relative importance and the velocities, decrease. When velocities decrease, the advection also decreases and in both cases, the amplitude of even constituents, inside the lagoon, lessens. The previous effect leads to a smaller difference when comparing the SSE, for the same H_s , between neap and spring tides.

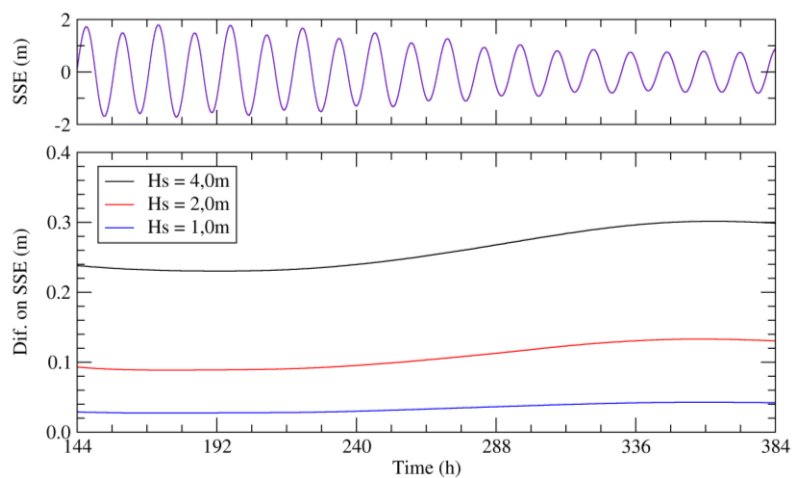


Figure 3.9 - Top: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE when comparing with a simulation without waves

The analysis of scenarios 1, 4 and 5 of Figure 3.5 allows concluding that for the same H_s (2,0 m) and Dir (315° N) at the boundaries, a higher T_p will increase the SSE inside the lagoon and also, the H_s in front of the lagoon were higher with growing T_p (Figure 3.10).

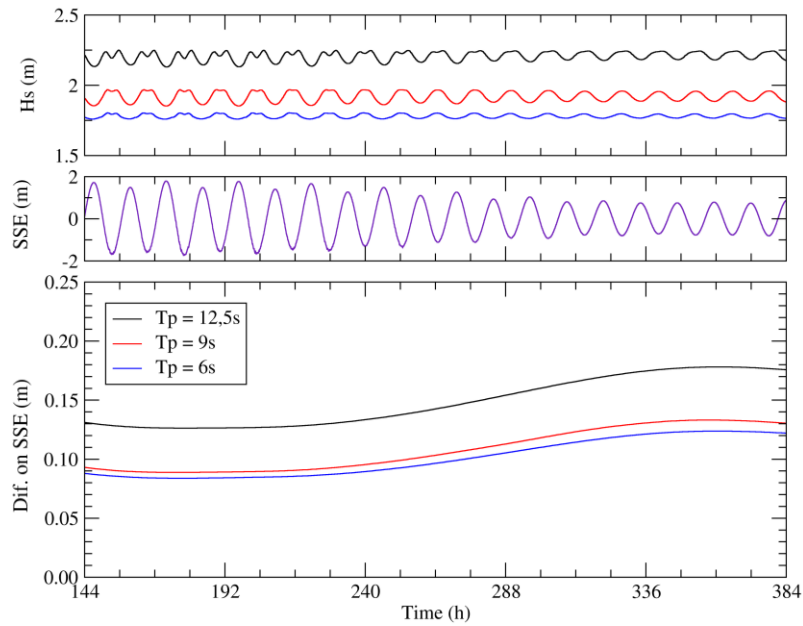


Figure 3.10 - Top: H_s at 5 m depth (black star in Figure 3.11); Middle: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE when comparing with a simulation without waves (Station 4)

An explanation for the increase of SSE is the group velocity, which corresponds to a term of the radiation stresses equations because a higher wave period will increase the group velocity ($c = LT$, where c is the wave celerity, L the wavelength and T the wave period). As group velocity increases, the radiation stresses increase too (Equations 11, 12 and 13 in **Appendix 1 - A**), and more momentum will be transferred to the water column. This transference explains why SSE increased for higher T_p .

Another explanation concerns the breaking type induced by different T_p . An analysis of the surf similarity parameter (Table 3.1) showed that for the same beach slope (β) and the same offshore wave height (H_0), the increase of the T_p will induce a different breaking type.

Table 3.1 - Wave breaking type analysis based on the surf similarity parameter varying the T_p

$T_p = 6s$				$T_p = 9s$				$T_p = 12,5s$			
β (°)	H_0 (m)	L_0 (m)	ξ_0	β (°)	H_0 (m)	L_0 (m)	ξ_0	β (°)	H_0 (m)	L_0 (m)	ξ_0
5	2	56.2	0.5	5	2	126.4	0.7	5	2	243.8	1.0
10	2	56.2	0.9	10	2	126.4	1.4	10	2	243.8	1.9
15	2	56.2	1.4	15	2	126.4	2.1	15	2	243.8	3.0

Where the surf similarity parameter is $\xi_0 = \tan \beta \sqrt{L_0/H_0}$ and the offshore wave length is $L_0 = gT^2/2\pi$. Often, the breaking type is plunging, but with the increase of the parameter T_p , it could be also collapsing/surging type. For higher T_p the radiation stresses become constrict to a smaller area, increasing

the wave setup and further, the SSE inside the lagoon. In Figure 3.11 is also shown the increase of energy in the breaking zone for waves with a T_p of 12,5 s.

The previous results concerning the study of the effect of T_p on the water levels should be carefully analysed. The coarse discretization and the truncation in the wave model for the high frequency band (from 0,2 to 1 Hz) can lead to the lack of wave energy for T_p smaller than 8 s. Therefore, it is expected that with a finer frequency resolution the differences between a T_p of 9 s compared to a T_p of 6 s would be smaller.

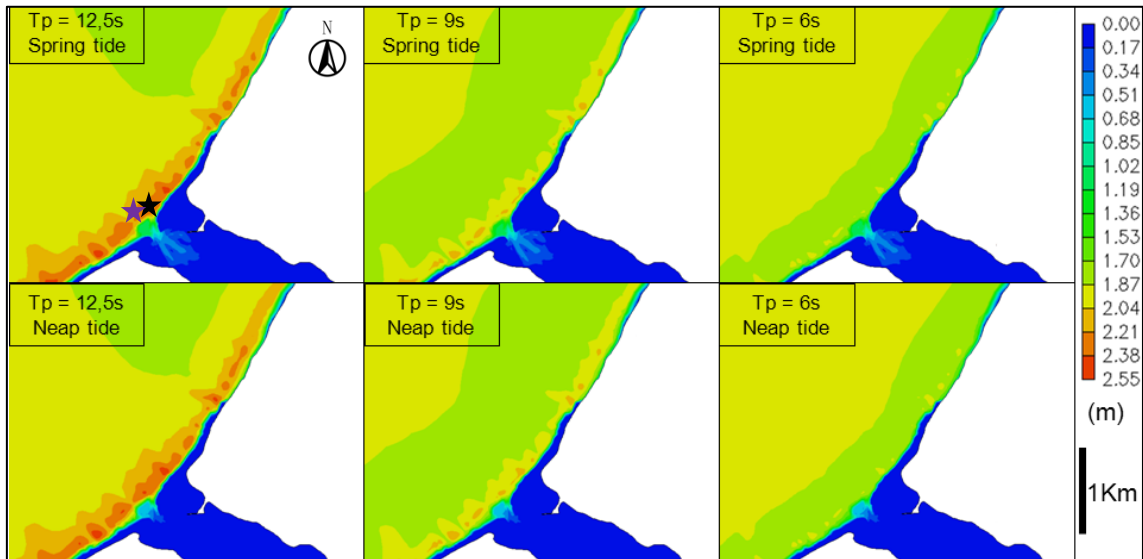


Figure 3.11 - H_s for a JONSWAP spectrum with constant H_s (2,0 m) and Dir (315° N) as forcing parameters

When the models were forced by waves with the same H_s (2,0 m) and T_p (9 s), which correspond to scenarios 1, 6 and 7 in Figure 3.5, the waves coming from 315° N had the higher H_s in the breaking zone and induced the higher increase on the SSE inside the lagoon, especially during neap tides (Figure 3.12).

When waves travel to the coastline, they are subjected to different phenomena. One of them is the wave refraction. The waves with a Dir of 315° N are less sensitive to the refraction phenomenon because more energy was concentrated in the breaking zone (Figure 3.13). This explains the higher H_s for waves with the 315° N angle of incidence.

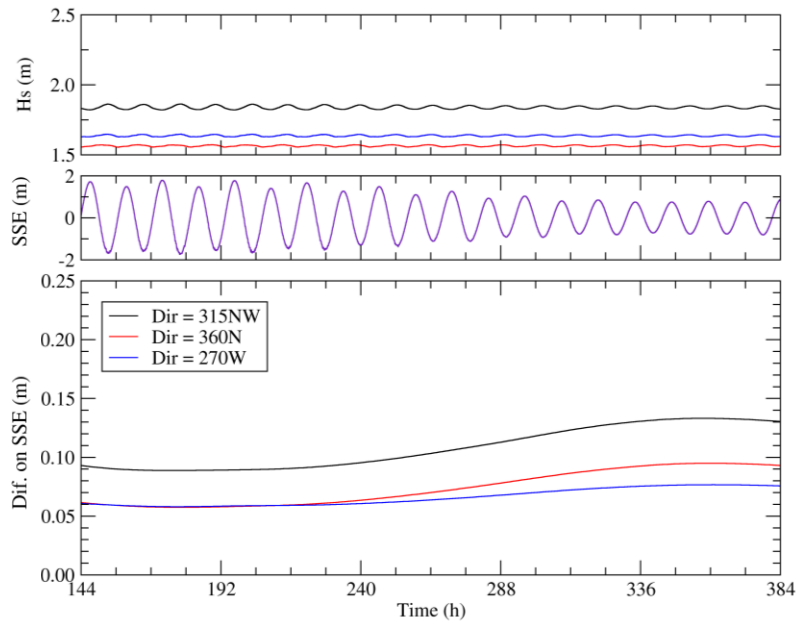


Figure 3.12 - Top: Hs at 10 m depth (black star in Figure 3.13); Middle: SSE for a simulation without waves in front of the inlet (purple star in Figure 3.11); Bottom: Differences on SSE when compared with a simulation without waves (Station 4)

The higher SSE inside the lagoon is explained due to the radiation stresses. The increase of energy in the breaking zone leads to an increase on the SSE inside the lagoon when compared with the other two Dir.

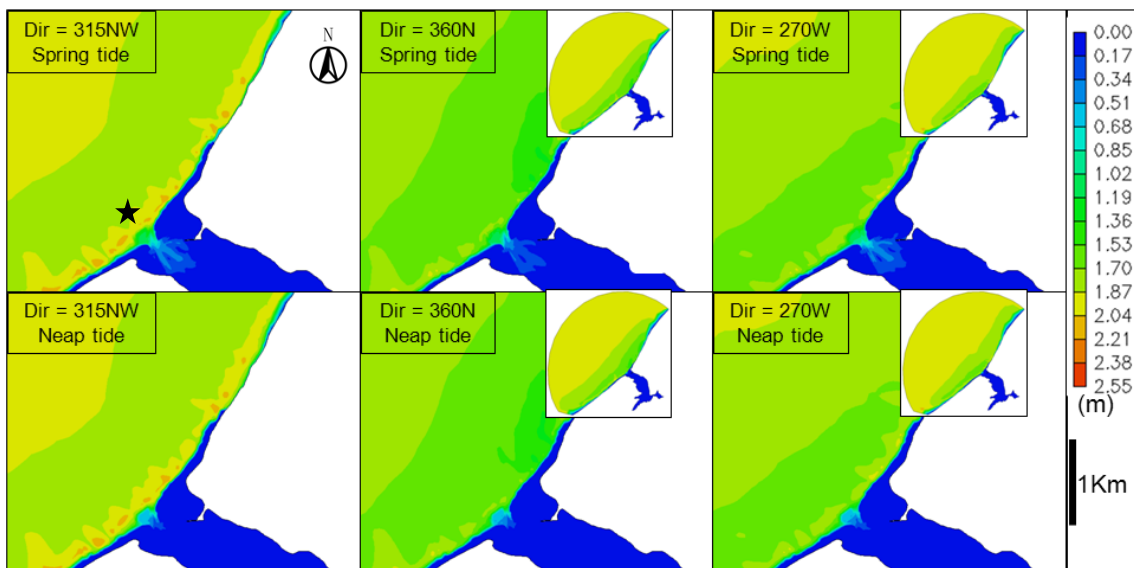


Figure 3.13 - Hs for a JONSWAP spectrum with constant Hs (2,0 m) and Tp (9 s) as forcing parameters

The ebb-delta explains the differences during spring and neap tides between N and W directions. During spring tide the water depth increases and the shoaling on the ebb-delta becomes smaller. When the water depth decreases (neap tide), the ebb-delta increases the wave shoaling for waves coming from W (Figure 3.12), and block the water movement towards the inlet, creating a difference of 1,5 cm in SSE between N and W Dir at neap tides.

3.3.4. The wave blocking phenomenon

This chapter 3 also aims to verify the wave blocking phenomenon experienced by Dodet *et al.* (2013). The H_s at the beginning of the flood delta (black star in Figure 3.5) increases during the flood and decreases during the ebb (first box in Figure 3.14). This happens because the water depths increase during the flood as more water could enter inside the lagoon. Even so, if the wave model takes into consideration the full interaction between waves and currents (hereafter FIBWC), the previous phenomenon is more evidence due to an increase of 9% in the H_s during the flood (first box in Figure 3.14). Also, the wave model did not take into account the white-capping effect due to the computational cost (Section 3.2.2). For that reason, it is expected that the H_s time series for the fully-coupled option will have a more asymmetric profile when compared with the time series of H_s for the partial coupled option.

The SSE and the depth-averaged horizontal velocities (DAHV) remained unaffected by the FIBWC effect (second and third boxes in Figure 3.14). The only difference when the wave model takes into account the FIBWC is the effective advection velocity (U_A in Equations 4, 5, 6 and 7 in APPENDIX A) which is not null. Therefore, the quantities obtained by the hydrodynamic model were unaffected (second and third boxes in Figure 3.14).

The wave orbital velocity (U_{rms}) during the flood with the FIBWC effect increased 33% when compared with the U_{rms} without the FIBWC (fourth box in Figure 3.14). During the ebb, the ebb-currents could drive the H_s to zero and the U_{rms} was also zero. The previous results are explained through the U_{rms} , because the U_{rms} formulations depend on the wave height in the numerator.

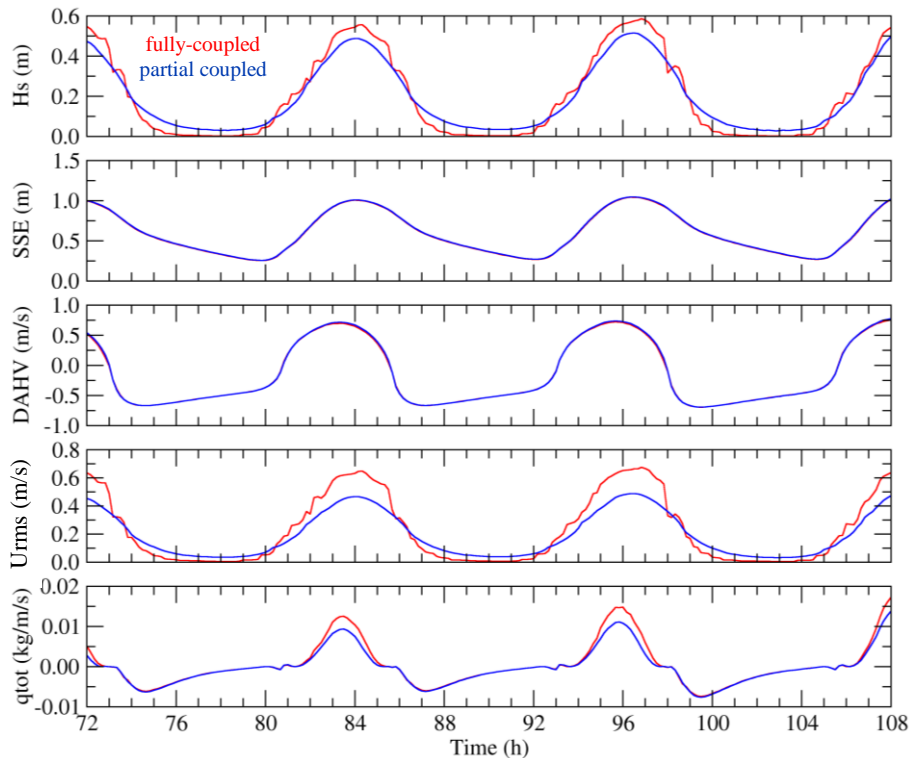


Figure 3.14 – Comparison between a simulation with the full wave-current interaction (red lines) and without (blue lines) for the H_s (first box), SSE (second box), depth-average horizontal velocities (third box), U_{rms} (fourth box) and sediment discharge (fifth box) between the third and fourth days on the black star (Figure 3.5)

The sediment transport was computed with the Soulby – Van Rijn (Soulby, 1997) formula. The results show that with the FIBWC effect, the sediment transport increased by 30% during the flood (fifth box in Figure 3.14). The U_{rms} explains the previous fact because, with the FIBWC effect, waves will put more sediment in suspension and then, the flood currents will transport them towards the lagoon. This way, the net sediment transport towards the lagoon increased over a tidal cycle, which is in accordance to the study made by Dodet *et al.* (2013).

4. Assessment of three dredging plans for the Óbidos lagoon

SUMMARY: This chapter concerns the assessment of three dredging plans to evaluate the improvements regarding the Óbidos lagoon inlet stability. A morphodynamic model which simulates the non-cohesive sediment transport taking into account the wave-current interaction was used to assess the morphological evolution during five months of three dredging plans. The results of the following simulations are reported: tidal prism, differences between ebb and flood durations, sediment volumes and fluxes through the inlet and application of empirical stability relationships to the Óbidos lagoon inlet. From these results, it is reasonable to assume that the new dredging plans improved the hydrodynamic conditions of the Óbidos lagoon. Morphologically, the higher dredging depths of the new dredging plans promoted the sediment accretion at the flood-delta. The addition of transverse channels improved the sediment export capacity during the studied period. This chapter ends with a summary which discusses the main advantages and drawbacks of each studied dredging plan in the Óbidos lagoon.

4.1. Review on seasonal closure, tidal distortions and dredging solutions in wave-dominated inlets

Understanding the dominant processes that control the morphodynamic changes in inlets is a key issue to develop sustainable engineering solutions. Ranasinghe and Pattiaratchi (2003) successfully pioneered in the simulation of the closure of a tidal inlet. They presented two mechanisms that are required in a morphodynamic model:

- 1) the interaction between inlet current and longshore currents
- 2) the effect of both inlet current and onshore sediment transport.

The first mechanism is described as the interruption of the longshore current by the inlet. As a consequence, this interaction will also interrupt the longshore sediment transport and induce the formation of a shoal up-drift of the inlet. The intensity of the longshore sediment transport will control the growth and the size of this shoal. If the inlet current (ebb-current) is strong enough to erode the sediments previously deposited, then the spit will not migrate downdrift. However, if the inlet current cannot erode the sediments in this shoal, a spit will continue to grow until the inlet closes or migrates.

The second mechanism can only dominate if the inlet current is small (less than 1 m/s) and in micro-tidal regimes. In such conditions, the tidal prism will be small. In the winter and under storm conditions, a linear bar is formed across the breaker zone. After the storm, the presence of long-period swell waves will continue to move the sand stored in the linear bar onshore. When the ebb currents are weak, the onshore sand transport, previously accumulated in the linear bar, will induce the inlet closure.

Bertin *et al.* (2015) have mainly focused their study on wave-related processes, which are the most important on wave-dominated inlets. They presented and verified the “bulldozer effect” studied in detail by Suastika (2004). Another wave-related process presented by Bertin *et al.* (2015) is the presence of lateral barotropic pressure gradients which accelerate the longshore flows towards the inlet.

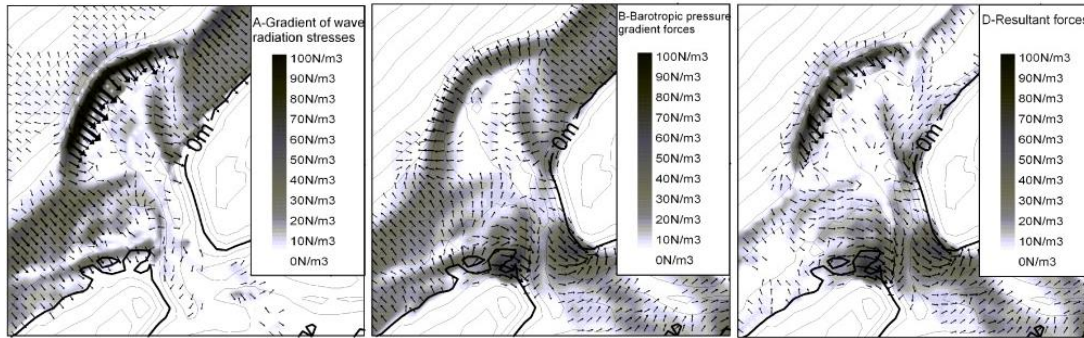


Figure 4.1 – Gradient of wave radiation stresses (left), barotropic pressure gradient forces (center) and resultant forces (right) at the Óbidos lagoon inlet for shore-normal waves of $H_s = 3,0$ m (adapted from Bertin *et al.*, 2009a)

The wave-induced radiation stresses gradients generate a barotropic pressure gradient on the northern and southern inlet's shores (Figure 4.1). The previous gradients are in equilibrium at the shores but not in front of the inlet, therefore, increasing the longshore sediment transport towards the inlet.

The same authors also show that the stronger tendency for the inlet closure on neap tides was partially due to the fortnightly compound tides¹ (e.g., MS₄) which reduce the ebb-dominance on neap tides. When tides propagate into the lagoon, the semidiurnal tidal constituents are damped, and quarter-diurnal and fortnightly nonlinear constituents start to develop (Mehta and Ozsoy, 1978).

Numerous authors have studied the nonlinear tidal interactions in shallow waters as in tidal inlets or estuaries (e.g., Speer and Aubrey, 1985; Friedrichs and Aubrey, 1988; Hoitink *et al.*, 2003). Tidal asymmetry is the difference in magnitude and duration between ebb and flood tidal currents (Dronkers, 1986). The tidal asymmetry occurs when the combination of multiple tidal constituents produces a tide with differences in the duration between the rise and the fall of the water elevation (Boon and Byrne, 1981). The phase difference between constituents creates what is known as flood or ebb dominance while the ratio of the constituents' amplitudes is responsible for the degree of distortion (Friedrichs and Aubrey, 1988). In frictionally-dominated estuaries with semi-diurnal tides, the main source of asymmetry is the interaction of the principal lunar semi-diurnal tide (M_2) with its first overtide², the lunar quarter-diurnal (M_4) (Speer and Aubrey, 1985). The physical mechanisms responsible for this distortion are the nonlinearities in the equations of motion (Parker, 1991).

Dredging solutions are used around the world, either to allow specific depths for navigation requirements (Dias *et al.*, 2000; Castelle *et al.*, 2007) or to reopen seasonally closed inlets (Vila-Concejo *et al.*, 2003; Oliveira *et al.*, 2006; Fortunato *et al.*, 2014). Usually, dredging solutions are expensive with prices having a wide spectrum of variability (according to the dredging project for the Óbidos lagoon in 2015 a dredging cost of 10 €/m³ can be estimated (Jornal das Caldas, 2014)). For that reason, dredging operations should be chosen with the knowledge needed to make optimal coastal management decisions (price, efficiency, effectiveness and footprint).

¹ Tides that result from the interaction between two major tidal constituents (e.g., M_2 and S_2 interaction result in MS₄).

² Tidal constituent whose frequency is an exact multiple of its major tidal constituent frequency (e.g., M_2 has a frequency of $1,4 \times 10^{-4}$ rad/s and its first overtide, M_4 , has a frequency of $2,8 \times 10^{-4}$ rad/s)

Dredging interventions in wave-dominated inlets are also frequent. Castelle *et al.* (2007) studied a stabilized wave-dominated inlet (Currumbin Creek, Australia), in which the infilling sediment problem continues to happen. Shaeri *et al.* (2014) stated for the previous inlet that “currently, an annual dredging campaign is needed to reduce the risk of flooding”. In the Óbidos lagoon the dredging interventions have been common (section 2.2.3) and some studies were already made (Oliveira *et al.*, 2006; Fortunato and Oliveira, 2007a; Bertin *et al.*, 2009a; Bruneau *et al.*, 2011). However, the morphological evolution for different dredging plans and the impact of transverse channels in terms of flushing capabilities remain open issues.

4.2. The morphodynamic model and the implemented methodology

4.2.1. Morphodynamic model general outline

The morphodynamic model used in this dissertation simulates the morphodynamics of non-cohesive sediments in coastal areas taking into account the wave-current interactions (Figure 4.2). Similar to section 3.2.4, ocean waves and tides force the model at its open boundary. First, the wave model calculates the wave parameters and the radiation stresses to compute the gradients of radiation stresses (S_{xx} , S_{yy} and S_{xy}) in the hydrodynamic model. After the calculation of η , u and v through the hydrodynamic model, the sediment transport module determines the sediment fluxes (Q_x and Q_y) taking into account the velocities and the water levels from the circulation model and the wave parameters from the wave model. These fluxes are the inputs of the bottom evolution model, which then solves the Exner equation. The bottom evolution model obtains the bed update (Δh) and shares the new bathymetry with the other models for the next time step.

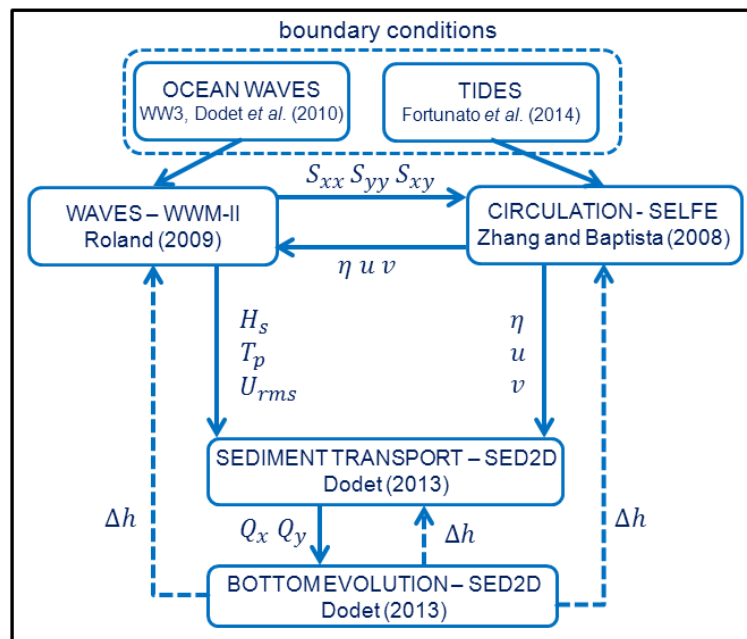


Figure 4.2 – Morphodynamic numerical scheme (adapted from Dodet, 2013)

The first two models (SELFE and WWM-II) were already described in section 3.2. The module SED2D (Dodet, 2013), which is a sediment transport and bottom evolution model, is coupled with the other two to form a full morphodynamic model. All these models share the same unstructured grid and work in parallel computations.

4.2.2. The sediment transport and bottom evolution model – SED2D

The sediment transport model has several alternative sediment transport formulations, ranging from the traditional ones (e.g., Eugelund and Hansen, 1967; Ackers and White, 1973) to the most recent, which include the sediment transport induced by waves (e.g., Camenon and Larsson, 2005 and 2008). In this dissertation, the Van Rijn (2007a,b) sediment transport formulation was used. This formulation takes into account both the bedload and the suspension transport, and is applicable to a wide spectrum of particle sizes (between 0,05 mm to 2 mm). Also, it includes the effects of both waves and currents. Usually, sediment transport formulations are derived for flat bottom conditions. To overcome this limitation, Lesser *et al.* (2004) proposed a formulation to include the bed-slope effect and SED2D takes this effect into consideration.

SED2D solves the Exner equation through a node-centred finite volume technique to update the bottom bathymetry every 20 s. In this dissertation, three filters were used due to the oscillations inherent to morphodynamic models:

- 1) an extrema filter
- 2) a slope filter
- 3) a diffusive filter.

Fortunato and Oliveira (2000) implemented the extrema filter to eliminate local extremes in the bathymetry after each morphodynamic time step. This filter is volume-conservative and introduces small numerical diffusion (Oliveira and Fortunato, 2002).

Avalanching was taken into account using a slope filter. After each morphodynamic time step, this filter limits the local slopes to a user-specified threshold. Roelvink *et al.* (2009) studied a beach profile and defined these values as 1 for dry and 0,3 for wet areas. Dodet (2013) used 1 for dry and 0,2 for wet areas in the Albufeira lagoon, despite the fact that this author was able to stabilize the sediment model without using this filter. In this dissertation, the slope filter took the value of 1 for dry and 0,2 for wet areas.

The diffusion filter (Fortunato and Oliveira, 2007b) aims to prevent the development of spatial oscillations through a modification on the Exner equation (Equation 26 in APPENDIX A). The value of the coefficient ε was set to 2, which is in accordance with Fortunato *et al.* (2014). All the details regarding the equations solved by SED2D are in APPENDIX A.

To summarize, the sediment fluxes (Q_x and Q_y) are calculated with one of the available formulations (Van Rijn, 2007a,b) taking into account the bed-slope effect (Lesser *et al.*, 2004) and the modified sediment fluxes (Fortunato and Oliveira, 2007b). Then, the calculation of the Exner equation with a node-centred method is used to update the bottom bathymetry. Before computing the next time step, the local extremes in the bathymetry were smoothed with the extrema filter (Fortunato and Oliveira, 2000) and the slope filter ensures that all element slopes are lower than a defined threshold. Finally, the model shares the bottom update (Δh) with the wave and the hydrodynamic model and proceeds to the next time step.

The sediment median grain diameter (d_{50}) was specified on SED2D (*d50.gr3*) and was adapted from Bruneau *et al.* (2011). Due to numerical oscillations on the northern and southern coasts, the d_{50} needs to be increased by approximately 0,4 mm in order to prevent oscillations (Figure 4.3). Previous simulations

of Bertin *et al.* (2009a) and Bruneau *et al.* (2011) have experienced difficulties to represent the down shore evolution of the inlet's northern barrier island until November 2001. To circumvent these previous difficulties, the d_{50} was reduced by 0,1 mm in the end of the transitional channel, near the flood-delta.

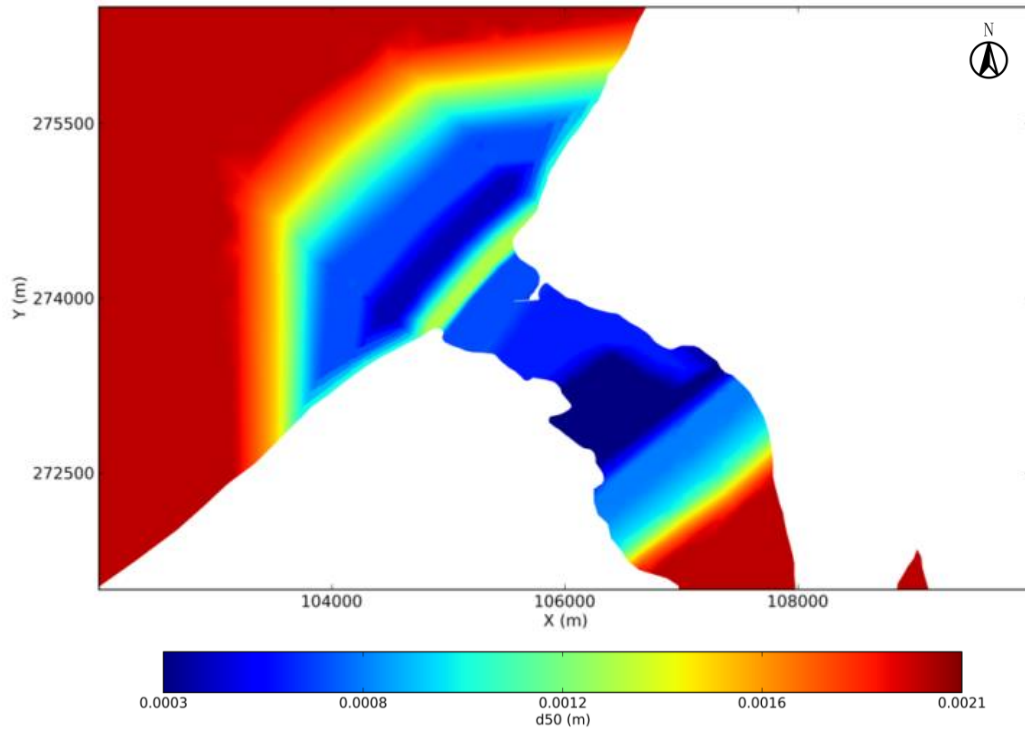


Figure 4.3 – Spatial distribution of the median grain diameter – d_{50} (m)

4.2.3. Sediment model verification period and statistical error measures

The verification period of the morphodynamic model was from July to November 2001. In November 2001 an intense field campaign was carried out by the Portuguese Hydrographic Institute in order to measure the bathymetry inside the lagoon.

The bathymetry measured in November will be compared to the simulated one, assessing its validation through two error measures: the brier skill score (BSS) (Equation 4.1) proposed by Sutherland *et al.* (2004) and the RMSE used by Fortunato *et al.* (2014). The BSS error measure compares three bathymetries for each node: the simulated bathymetry, the measured bathymetry and the initial bathymetry. The error values could go from negative values to 1.

$$BSS = 1 - \frac{MSE(Y,X)}{MSE(B,X)} = 1 - \frac{\sum_{i=1}^n (Y-X)^2}{\sum_{i=1}^n (B-X)^2} \quad [4.1]$$

where X is the simulated bathymetry, Y is the measured bathymetry, B is the initial bathymetry and n the total number of nodes in the grid. A negative value of this error measure means that the initial bathymetry is closer to the measured bathymetry than the simulated bathymetry. A value of 1 represents a perfect agreement between the simulated and measured bathymetries. Van Rijn *et al.* (2003) proposed a displacement error (δ) for the BSS error due to inaccuracies in the measured bathymetries (Equation 4.2). The value of the displacement error used in this dissertation was 0,3 m, similar to Bertin *et al.* (2009a).

$$BSS_{vR} = 1 - \frac{\sum_{i=1}^n (|Y-X|-\delta)^2}{\sum_{i=1}^n (B-X)^2} \quad [4.2]$$

4.2.4. Boundary conditions of the sediment model

As explained in section 3.2.4, the forcing agents at the boundaries were tides and waves. All the parameters mentioned previously for the hydrodynamic and wave model were the same (as in sections 3.2.2, 3.2.3 and 3.2.4).

In the case of large waves (with H_s higher than 3 m) the coupling between the wave and the hydrodynamic models generated spurious velocities near the closed boundaries. As waves reach the closed boundaries they are reflected and this generates large velocity vectors that cause numerical oscillations. WWM-II does not have an option to specify different reflection coefficients and to prevent this phenomenon the bathymetry near the northern and southern coast was increased to 5 m above the MSL (Figure 4.4). This procedure guarantees that the waves break before reaching the boundaries.

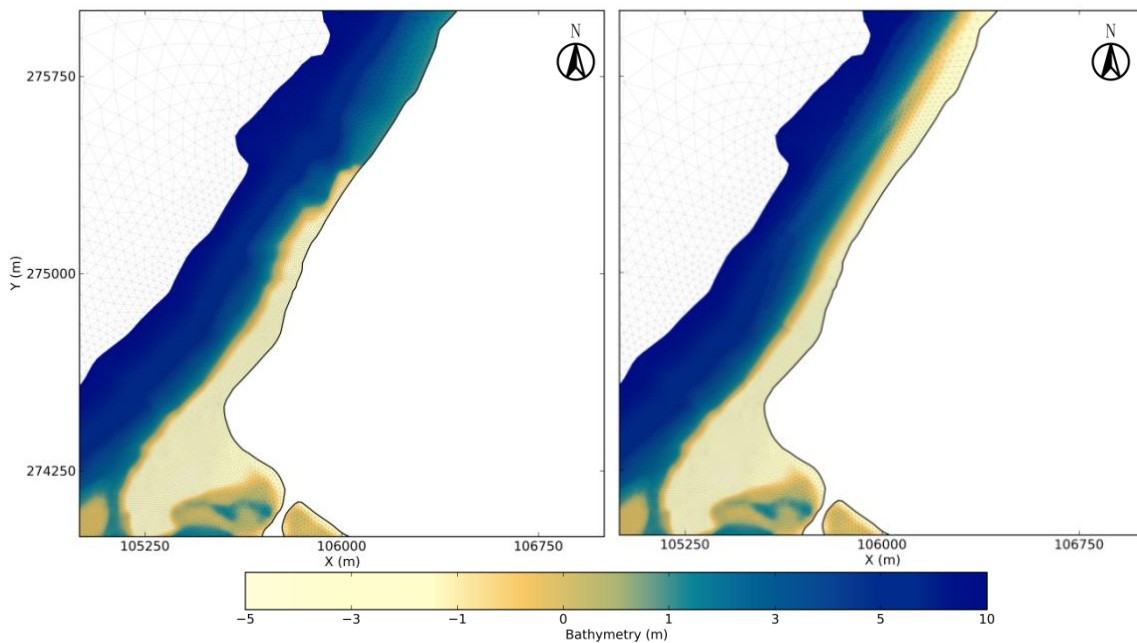


Figure 4.4 - Comparison between the initial (left) and the changed (right) bathymetry on the northern margin

Even after the increase of 0,4 mm in the d_{50} on the southern and northern shores (Figure 4.3) and the changes in the bathymetry on the southern and northern coasts (Figure 4.4), the morphodynamic model continued to have a number of instabilities after 5 months because the waves are reflected against the grid boundaries and the development of spurious bathymetries occur. In physical terms, at the wet and dry interface on beaches, the bottom depth remains constant in time because the eroded sand from the wet area is supplied with sand from the dry area. Therefore, the previous difficulty was avoided by maintaining constant the bottom depth at the grid boundaries, thereby inducing the wave breaking inside the domain and preventing further wave reflection. Further details are presented in the APPENDIX C.

4.2.5. The simulated dredging plans under study

Three dredging plans were chosen in order to study the morphodynamic changes in the Óbidos lagoon induced by the dredging interventions:

- 1) the dredging plan implemented in the last decades proposed by DHI (1997) (DP1) which corresponds to the initial bathymetry of July 2001
- 2) the dredging plan proposed by Fortunato and Oliveira (2007a) (DP2)

3) the dredging plan similar to (2) but with only one transverse channel (DP3)

In order to understand the initial bathymetry for the different dredging plans please consult the Figure 4.6 for the DP1, APPENDIX D for the DP2 and APPENDIX E for DP3. In terms of dredging volumes, the DP2 represents an additional dredging of $5,1 \times 10^5 \text{ m}^3$ and the DP3 of $3,5 \times 10^5 \text{ m}^3$ when both compared to DP1.

Due to the actual computational time (approximately 8 to 10 days), the simulations were restricted to five months (from July to November 2001). APA major concerns were the stability of the inlet and the direct impact of a dredging plan on the Óbidos lagoon. The results obtained with the three different dredging plans will be discussed quantitatively and qualitatively. A series of images representing the five-month evolution are presented, together with the evolution of the tidal prism, the differences in ebb-flood duration for two specific points inside the lagoon and at the inlet, and the comparison between the initial and final bathymetries. Further, the sediments fluxes that crossed the inlet were calculated, together with the application of empirical stability relationships to determine the stability of the Óbidos lagoon inlet.

4.3. The verification of the bathymetry simulated until November 2001

4.3.1. The qualitative and quantitative morphodynamic model simulation

As mentioned above, the period of the simulation extended for 5 months. At the end of this period the model shows a very good agreement with data (Figure 4.5). Bertin *et al.* (2009a) obtained a value of 0,2 for BSS with a displacement error of 0,3 m. With the same displacement error, the model shows a BSS error of 0,60. These values for BSS show that the model accurately simulates the morphodynamic changes that occur in the Óbidos lagoon. According to Sutherland *et al.* (2004), a model with a BSS of 0,47, without any displacement error, can be considered in a “good agreement” with the measured bathymetry. In terms of RMSE, Fortunato *et al.* (2014) reached an average RMSE of 1,0 m during the 7 months of their morphodynamic simulation in the Albufeira lagoon. The calculated RMSE for the morphodynamic model used in this dissertation was 0,82 m. This previous value reflects the good accuracy of the model to predict morphological changes.

Qualitatively, the model could reproduce the southward migration of the transitional channel and its final north-western orientation. This migration is also verified on the measured data. Nevertheless, the model slightly underestimates the migration of the transitional channel (100 m on measurements against 90 m on the simulation).

The two main channels (C1 and C2) inside the lagoon were also accurately reproduced. In the beginning, there are two channels. After 2 months, the meanderization closes one of them and enlarges the other (Figure 4.6). This morphodynamic change can be seen both on data collected and on the simulation (Figure 4.5).

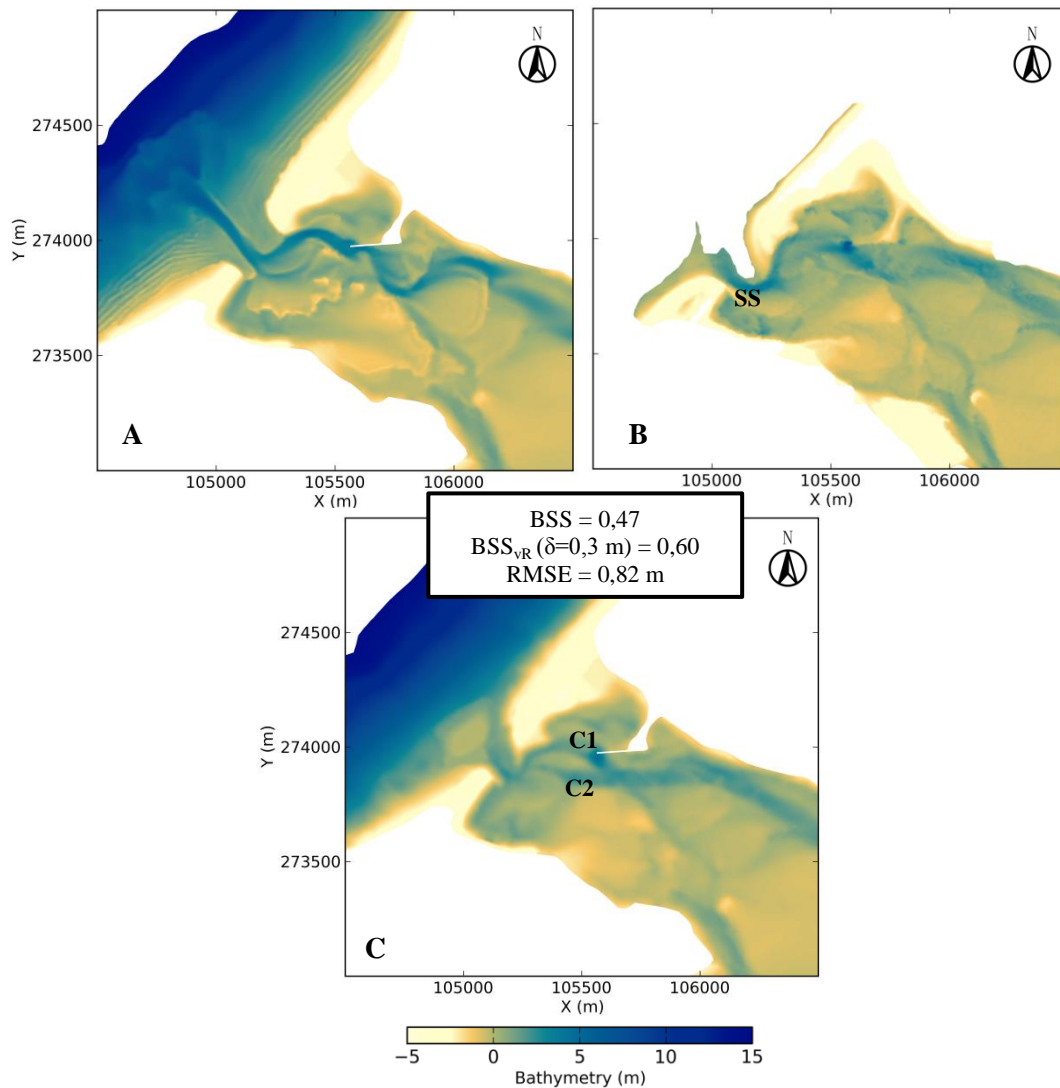


Figure 4.5 – A – Simulated bathymetry at the end of November; B – Measured bathymetry during November; C- Initial bathymetry in July 2001; Statistical error measures

The model could not catch the full shoaling of the southern and the northern margins, neither the formation of a sand spit (SS) in the southern margin (Figure 4.5). The formation of this spit and also the underestimated shoaling could be due to the model disregards the swash and the eolic transports and the high d_{50} on the northern margin needed to stabilize the model could influence the longshore sediment transport capacity towards the inlet.

The initial bathymetry has an ebb-delta that was breached after 3 months (Figure 4.6). This phenomenon was described in section 2.1.3. The wave-induced currents will move the swash bars that are formed by the longshore sediment transport landwards, narrowing the ebb channel. Then, the ebb delta breaches due to the overtopping phenomenon and the transitional channel was relocated southwards in September (Figure 4.6).

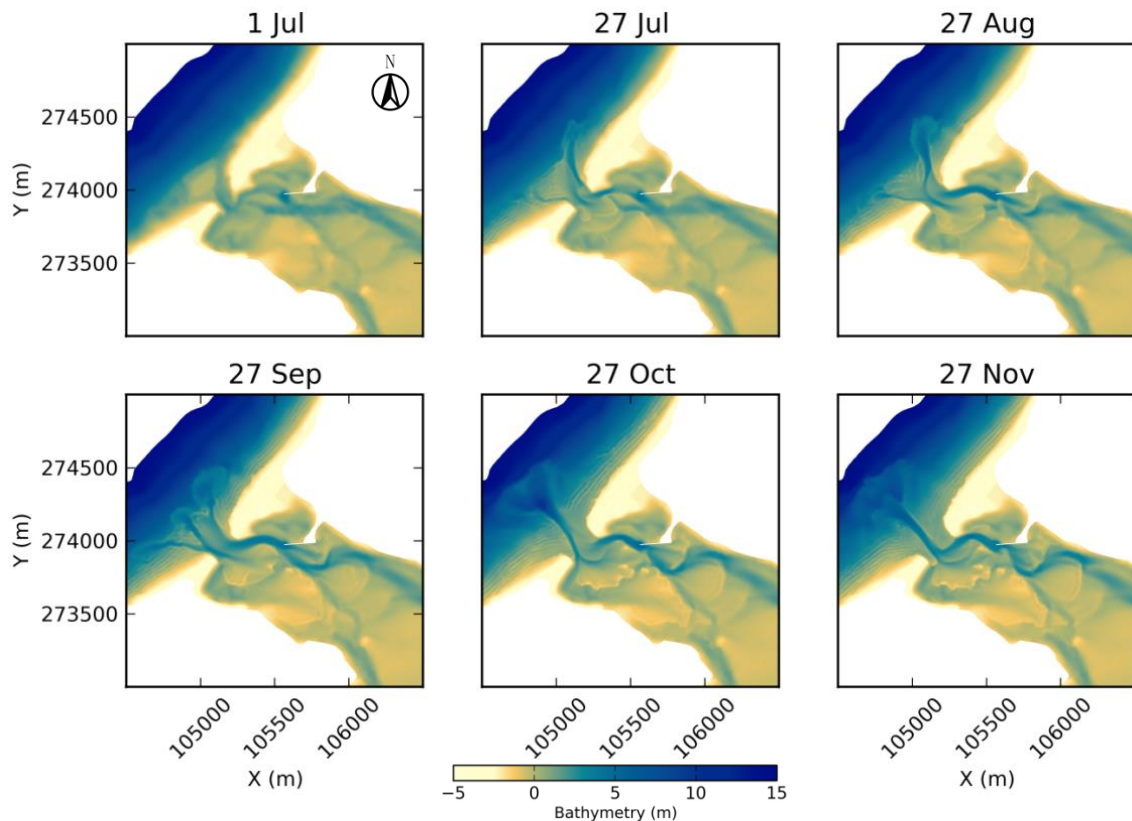


Figure 4.6 –Morphodynamic evolution of the Óbidos lagoon between July and November 2001 for DP1

The effect of the sheet-pile curtain was also displayed during this 5 months period. This curtain will reduce the flow area, increasing the velocity. The channel C1 (Figure 4.5 - C) starts to become deeper and turns into a preferred channel to ebb-currents wash out the sediments of the inlet. This effect of the curtain prevents the channel C1 migration northwards (Figure 4.6).

4.3.2. The evolution of the M_2 harmonic constituent amplitude

According to Bertin *et al.* (2009a) the growth of the tidal amplitude inside the lagoon can characterize the enlargement and infilling of the wave-dominated inlets. Figure 4.7 shows the data (solid line) and model (dashed-dot line) evolution of the harmonic constituent amplitude M_2 in Cais do Arelho tidal gauge based on the study of Bertin *et al.* (2009a) and simulated by the model used in this dissertation (dashed blue line).

The evolution of the M_2 constituent amplitude was computed with WORLD TIDES (Boon, 2013) using sequential intervals of 29 days with a sliding window of 15 days. The confidence intervals and the RMSE of the previous model varied between 75% and 90%, and 0,20 and 0,12 m, respectively.

The morphodynamic model can clearly represent the inlet enlargement until the end of September and after this period, the infilling until November. The measured amplitude (solid line) grows from 0,38 to 0,53 m and the simulated (dashed line) from 0,35 to 0,54 m. After September, the measured amplitude dropped to 0,43 m and the simulated to 0,47 m. However, the model simulates the peak of amplitude (dashed line) approximately 15 days after the amplitude peak measured in data (solid line) and overestimates the value by 0,01 m. This lag of 15 days can be related to the uncertainty of the precise moment that bathymetry was measured because the bathymetry measurements can take from few days to some weeks.

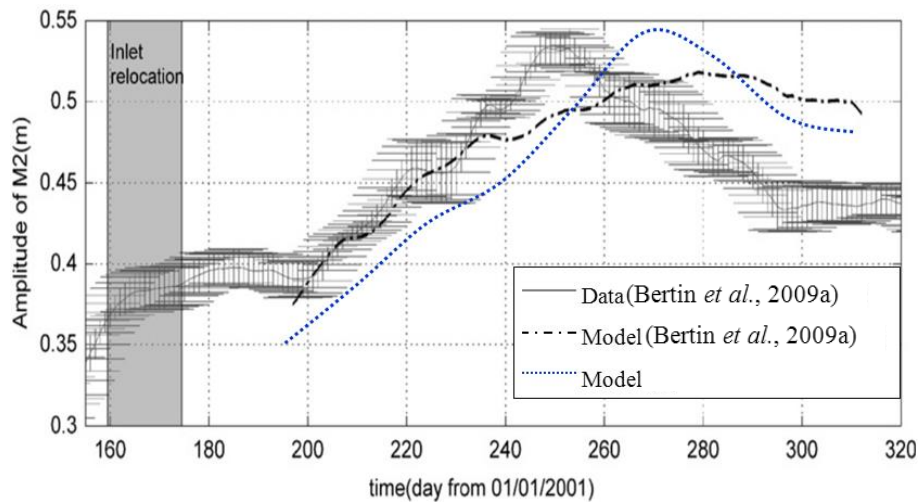


Figure 4.7 – Evolution of the M_2 amplitude in Cais do Arelho from July to November 2001 (adapted from Bertin *et al.*, 2009a)

4.4. The impact of dredging plans on the inlet’s behaviour

4.4.1. Qualitative evolution of the inlet

The DP2 clearly changes the behaviour of the inlet after 5 months (Figure 4.8). The cross-section of the transitional channel becomes wider and deeper. In average terms, the transitional channel depth, width and area for DP2 and DP3 become 25, 40 and 60% higher when compared with the DP1. These evidences reflect the characteristics of an ebb-dominated inlet, where the ebb-currents are the main agents that transport sediments seawards. Channel C1 (Figure 4.5 - C) also becomes deeper and wider with the DP2 and was less meandered since the channel did not reach the sheet-pile curtain in the northern margin.

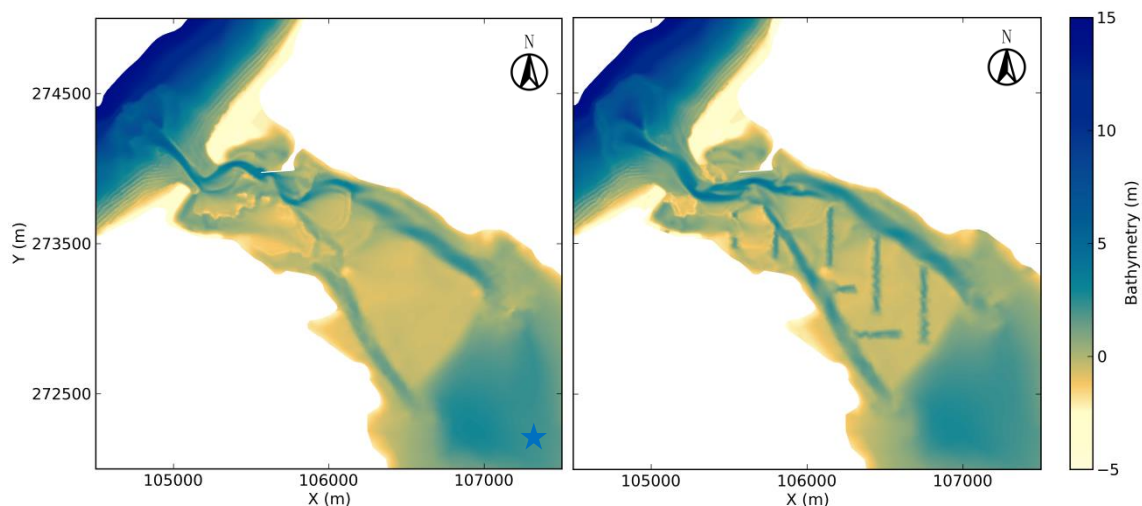


Figure 4.8 – Qualitative comparison between the DP1 (left) and the DP2 (right) at the end of November 2001

Overall, the morphodynamic evolution of the DP3 was very similar to the DP2 (Figure 4.9). The inlet becomes more ebb-dominant with a wider and deeper transitional channel. With DP2 and DP3, the amount of sand that enters near the southern margin was limited to a smaller area than with the DP1. Such an action confirms the more stable and less meandered configuration for the transitional channel of DP2 and DP3.

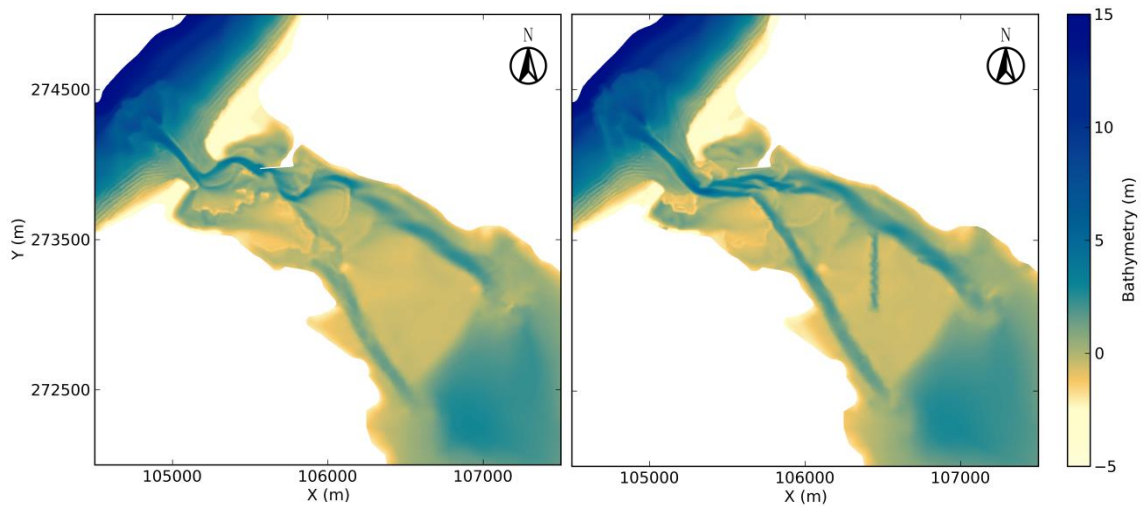


Figure 4.9 - Qualitative comparison between the DP1 (left) and the DP3 (right) at the end of November 2001

The transverse channels addition promoted an increase of the sediment transport capacity in Channel C2 (southern channel) because, when comparing DP2 with DP3 (Figure 4.8 and Figure 4.9), more sediments were accumulated at the end of Channel C2. In terms of hydrodynamics, it is not clear that the transverse channels would largely improve the stability of the inlet. Also, the importance of the western south transverse channel in DP2 is quite small considering that this channel started to meander after 3 months (APPENDIX D). Overall and based on Bruneau *et al.* (2011), the results suggest that the inlet is in a stable configuration after 5 months for the three dredging plans (Figure 2.12).

4.4.2. Evolution of the tidal prism

The tidal prism is the volume of water that enters inside a lagoon during half of the tidal cycle and was calculated multiplying the area of the lagoon by the tidal amplitude in order to compare quantitatively the three dredging plans (Figure 4.10). Diurnal tidal constituents in the tidal prism time series caused some oscillations (dashed lines). Therefore, a 6-day filter was applied (solid lines) in order to improve the visualization of the results.

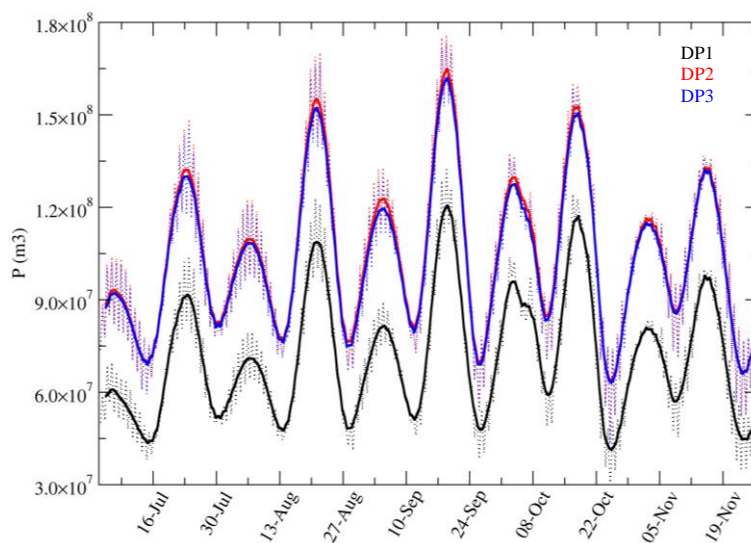


Figure 4.10 – Tidal prism of DP1 (black line), DP2 (red line) and DP3 (blue line) from July to November 2001

The tidal prisms of DP2 and DP3 in average terms were 32 and 31% higher than DP1, respectively. This behaviour shows the improved capability of water renewal with DP2 and DP3. As more water enters the lagoon, more sediments are washed out by the ebb-currents. The tidal prism is also associated with the inlet stability through the PA relationship (O'Brien, 1969) and the results suggest that the DP2 and DP3 promoted the inlet stability comparing to DP1.

Figure 4.11 displays the ratio between the tidal prisms for DP2 and DP3 with the DP1 during the studied period. The differences between both ratios are quite small with the largest (0,03) in the beginning of September. Between the middle of September until the end of October the ratios reduced because a severe wave climate, with a H_s larger than 5,0 m, was presented in the beginning of October.

The regression analysis made on the data clearly shows a decreasing behaviour of both solutions which converge to the tidal prism of the DP1. Assessing a constant trend, the tidal prism of DP2 and DP3 will be equal to the DP1 after 742 and 740 days, respectively. In other words, after 18 months the DP2 and DP3 solutions will have the same tidal prism as the DP1.

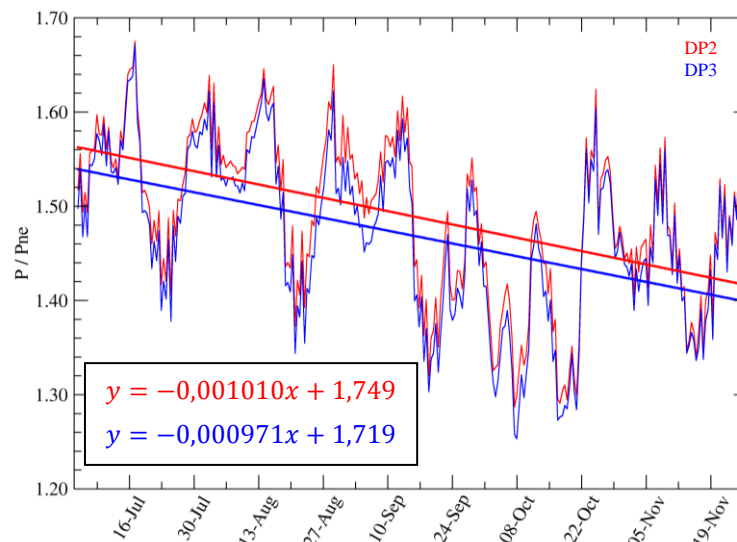


Figure 4.11 – Ratio between the tidal prism of the DP2 (red lines) and the DP3 (blue lines) with the tidal prism of the DP1 and the linear regression analysis for DP2 and DP3 (strong lines) with the respective coefficients.

4.4.3. Differences in ebb-flood durations

Sediments deposited inside the lagoon are mainly transported seawards by the ebb currents. The differences between ebb and flood durations give a quantitative measure of the magnitude of the ebb currents because these currents will be stronger if the ebb durations decreased. The analysis on the differences in ebb and flood durations shows that DP2 and DP3 improve the hydrodynamic characteristics of the Óbidos lagoon (Figure 4.12 and Figure 4.13). The previous dredging plans reduced the differences between ebb and flood durations, thereby promoting faster ebb when compared to the DP1 in two different locations.

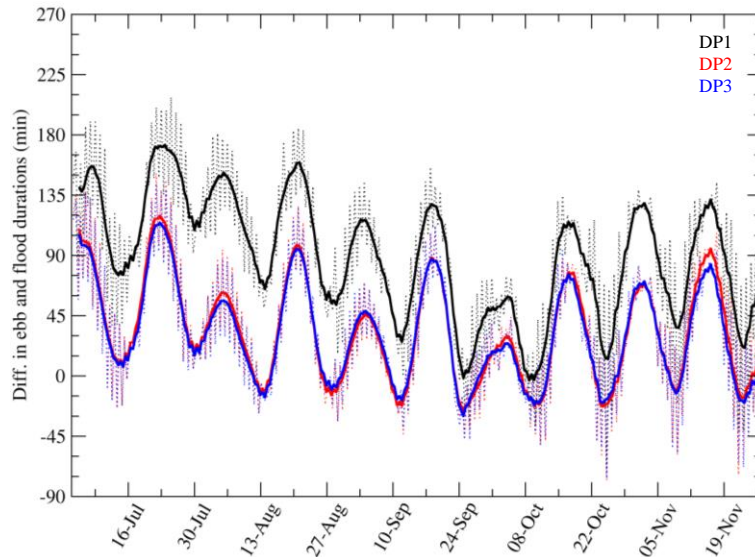


Figure 4.12 - Difference in ebb-flood durations near Cais do Arelho for the DP1 (black line), the DP2 (red line) and the DP3 (blue line) from July to November 2001

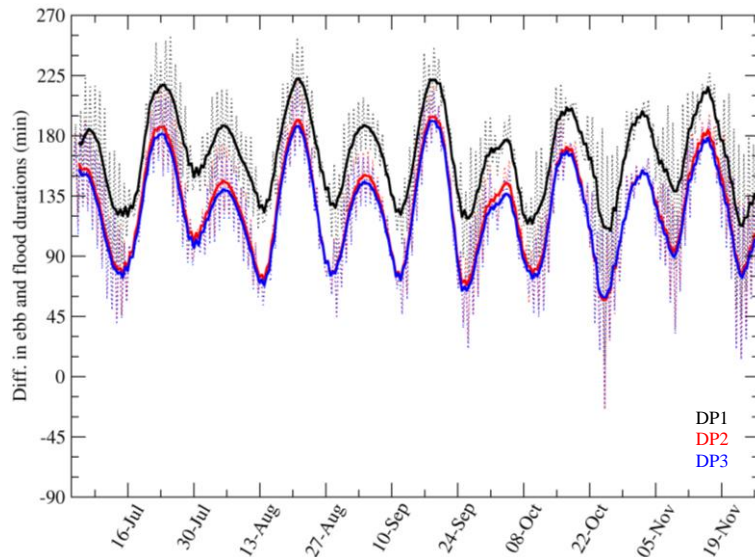


Figure 4.13 - Difference in ebb-flood durations inside the lagoon (blue star in Figure 4.8) for the DP1 (black line), the DP2 (red line) and the DP3 (blue line) from July to November 2001

Table 4.1 summarizes the differences between ebb and flood durations for the different scenarios near Cais do Arelho and inside the Óbidos lagoon. In average, the dredging solutions DP2 and DP3 decreased the differences near Cais do Arelho by 60% and inside the lagoon by 25%, compared with the DP1. The previous results suggest that the dredging solutions promote a less flood-dominant behaviour of the Óbidos lagoon.

Table 4.1 – Average differences between ebb and flood durations for two points inside the Óbidos lagoon

<i>Differences (min)</i>			
	<i>DP1</i>	<i>DP2</i>	<i>DP3</i>
Cais do Arelho	89	33	32
Central zone of the Óbidos lagoon	165	126	123

4.4.4. The Óbidos lagoon ability to flush out sediments

Figure 4.14 and Figure 4.15 show the differences between the initial and final bathymetries and the spatial distributions of the eroded (blue) or accreted (red) zones of sediments. The DP2 could export 5% more sediments, while the DP3 exports 1% less when both are compared to the DP1. In terms of deposition, the DP2 and the DP3 solutions import 20% and 12% more sediments than DP1, respectively. The previous values of the increased sediment importing capacity displayed by DP2 and DP3 in the Óbidos lagoon suggest that the accretion rate will decrease asymptotically until an equilibrium depth is reached, as suggested by Vicente and Uva (1985) and verified in APPENDIX G.

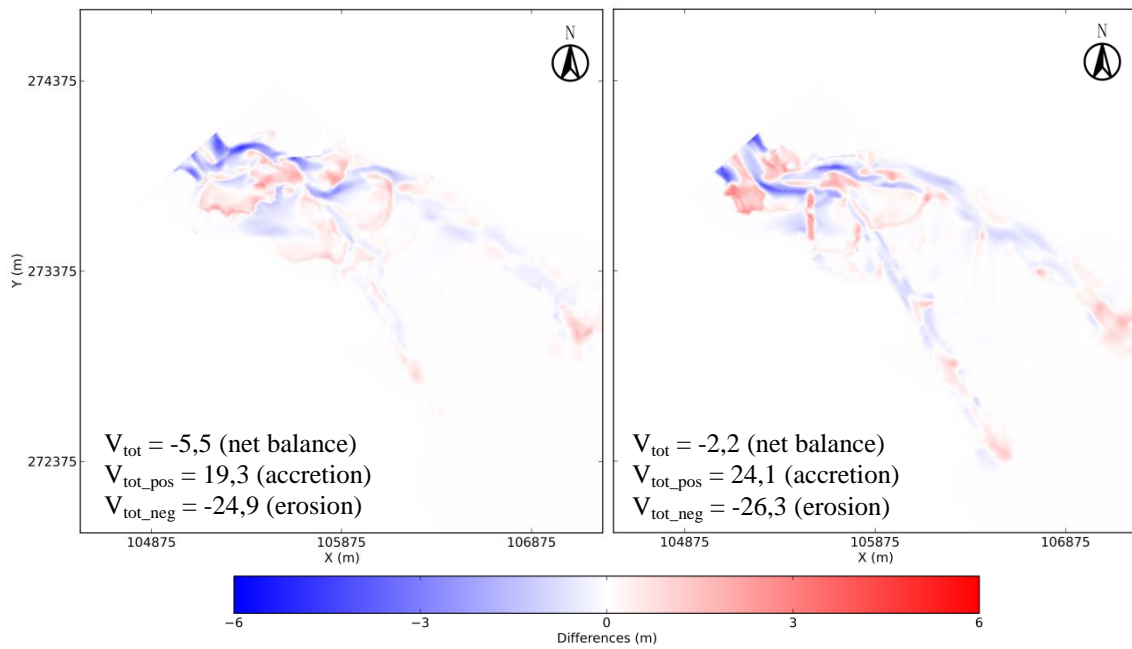


Figure 4.14 - Differences between the initial and 5 month simulated bathymetries for the DP1 (left) and the DP2 (right) and the sediment volumes that crossed the inlet ($\times 10^4 \text{ m}^3$)

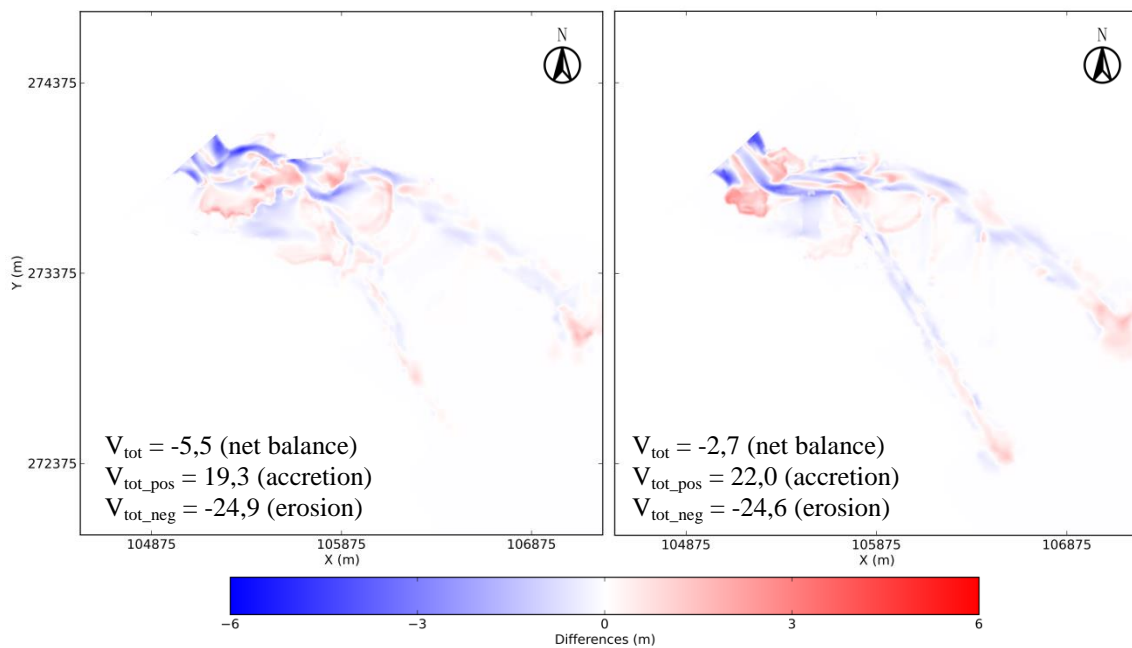


Figure 4.15 – Differences between the initial and 5 month simulated bathymetries for the DP1 (left) and the DP3 (right) and the sediment volumes that crossed the inlet ($\times 10^4 \text{ m}^3$)

Dredging is a human intervention which changes the lagoon's natural equilibrium. With an intervention similar to the DP2, the Óbidos lagoon will be away from its natural equilibrium (dredging $5,1 \times 10^5 \text{ m}^3$), this way, the accretion rate will be higher in the beginning and decreases until it reaches the DP1, which is a solution with smaller dredging depth.

The comparison between DP2 and DP3 shows that the impact of the transverse channels in dredging solutions was reproduced through the increased capacity in 6% of exporting sediments. The larger differences of the accreted zones were in the transverse channels and near the inlet, where the ebb currents are weaker and become preferred zones to sediment accretion.

In order to quantify the sediment fluxes and the differences between ebb and flood at the inlet, six simulations were made only with the hydrodynamic model SELFE. Each simulation represents the initial and final bathymetry for each dredging plan. Simulations last 30 days with one ramp-up day and the model parameters were same as section 3.2.3 except for harmonic constituents. At the open boundaries, the regional tidal model of Fortunato *et al.* (2014) computed the amplitude and phase of the seven harmonic constituents which have the higher amplitudes in semi-diurnal regimes: M_0 , M_2 , S_2 , M_4 , MS_4 , M_6 and MS_6 . Further, a harmonic analysis during 29 days was made on the velocities and on the sea surface elevations over the domain. Finally, the section of the inlet chosen to calculate the sediment fluxes and ebb and flood durations was parallel to the coastline and centred in the middle of the barrier islands (APPENDIX D).

As already mentioned, the differences between ebb and flood durations can determine the ability of the Óbidos lagoon to flush out sediments because a shorter ebb will increase the ebb-currents, thereby increasing the sediment transport seawards. Table 4.2 presents the differences between ebb and flood durations for each dredging plan at the inlet. The three dredging plans exhibit similar asymmetries for the initial bathymetry of July 2001, although a decrease of 23% can be seen for the final bathymetry of November 2001 when comparing the DP2 and DP3 with the DP1. Previous results enhance the stronger ebb-currents which take place for the DP2 and the DP3.

Table 4.2 - Differences between ebb and flood durations for the initial and final bathymetries and for each dredging plan

<i>Differences (min)</i>		
	<i>July 2001</i>	<i>November 2001</i>
DP1	191	132
DP2	195	104
DP3	194	100

Further, the sediment fluxes were computed at the inlet with the velocities and water levels obtained by the previous harmonic analysis. The empirical formulation used was Engelund and Hansen (1967) (hereafter EH67) which only accounts with the water depths, d_{50} and velocities for the sediment transport. The d_{50} was set on 1 mm and the sediment fluxes were integrated over half of the tidal cycle and over the inlet cross-section.

The results suggest the opposite in what concerns the ebb currents based on the differences between ebb and flood durations (Table 4.3). Previously, the ebb currents were stronger for DP2 and PD3 because the

differences in ebb and flood durations lessened, although, the sediment fluxes for DP2 and DP3 decreased during the ebb and increased during the flood. However, the EH67 formulation does not account for the wave orbital velocities neither the current refraction and for that reason, the higher depths of DP2 and DP3 lead to smaller velocities which were power up by five, thereby decreasing the sediment fluxes at the inlet for DP2 and DP3.

Table 4.3 – Ratios between the sediment fluxes integrated in half of the tidal cycle (ebb or flood) and in the inlet cross-section of DP2 and DP3 by DP1 for the initial and final bathymetries

<i>Ratios (-)</i>				
	<i>July 2001</i>		<i>November 2001</i>	
	<i>Ebb</i>	<i>Flood</i>	<i>Ebb</i>	<i>Flood</i>
DP2	0,16	0,75	0,65	1,23
DP3	0,15	0,66	0,81	1,18

4.4.5. The application of empirical stability relationships to the Óbidos lagoon inlet

Several empirical relationships are used to predict or to evaluate the stability of an inlet based, for example, on the tidal prism and on the area of the inlet cross-section. These relationships were already used over the world, and some of them were presented in section 2.1.2. In this dissertation, it was decided to perform the application of the PA relationship (Equation 2.2), the Escoffier (1977) diagram and the Bruun and Gerritsen (1960) criterion to evaluate the stability of the Óbidos lagoon inlet for the different dredging plans during the studied period.

The PA relationship relates the tidal prism at spring tides with the area of the inlet cross-section. For each of the six simulations performed in section 4.4.4, the maximum tidal prism during the 29 days and the area of the inlet cross-section (APPENDIX D) were calculated. The stability curve presented in Figure 4.16 was made with the Jarret (1976) empirical coefficients of $5,02 \times 10^{-4}$ for C_{PA} and 0,84 for n .

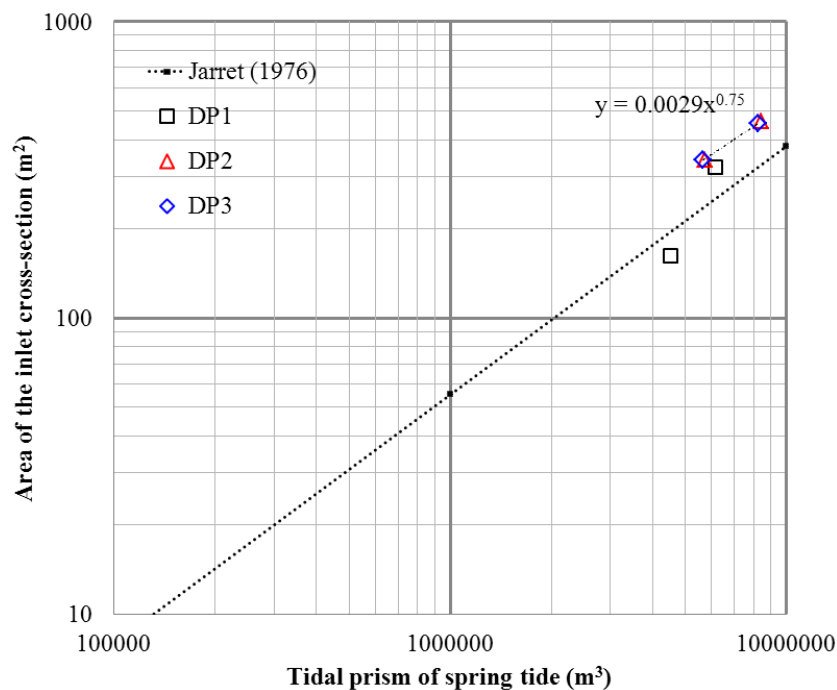


Figure 4.16 – Tidal prism-cross-section relationships for the three dredging plans

The persistent behaviour displayed by the DP2 and DP3 with the same distance to the Jarret's stability curve and also the displacement of the DP1 dot in November 2001 suggest that the equilibrium curve for the Óbidos lagoon inlet is above the one drawn with the coefficients of Jarret (1976). In this dissertation, is therefore proposed an equilibrium curve (PA relationship) for the Óbidos lagoon inlet with the values of $2,9 \times 10^{-3}$ for C_{AB} and 0,75 for n .

The Escoffier (1977) diagram relates the maximum velocity on the transitional channel with the area of the inlet cross-section. This diagram gives a qualitative measure of the inlet stability as described in section 2.1.2. The results present the three dredging plans where the left dot is the bathymetry of July 2001 and the right dot is the bathymetry of November 2001 (Figure 4.17).

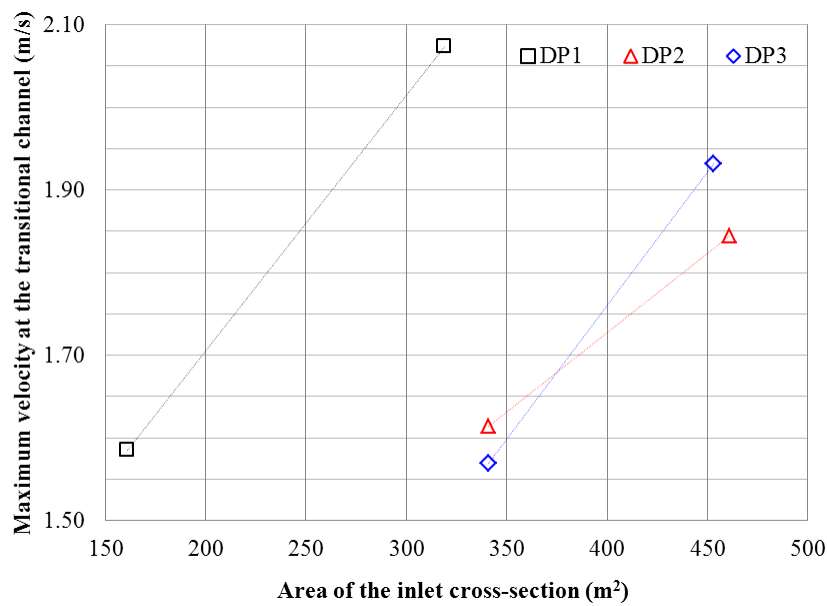


Figure 4.17 – Escoffier (1977) diagram for each dredging plan

Due to the positive slope of the lines for both dredging plans, it can be concluded that the dredging plans are within the left tail of the Escoffier (1977) diagram. Therefore, the results suggest that all the dredging plans are within an unstable equilibrium. This unstable equilibrium is reproduced by the inlet shoaling southwards in DP1 and by the enlargement of the cross-sectional area and the associate increase on the maximum velocities for DP2 and DP3. However, the smaller slope presented for DP2 suggest that this dredging plan is the less unstable.

Finally, the application of the Bruun and Gerritsen (1960) criterion is presented in (Table 4.4). The values of the spring tidal prisms were the same used in the PA relationship application and the value of M_{tot} was $1,6 \times 10^6 \text{ m}^3/\text{year}$. The previous value for the gross annual drift was derived from the previous work of Nahon (unpublished data). This author estimated the gross annual drift between 1953 and 2010 in the Óbidos lagoon with the empirical expression that relates the wave height at breaking, the angle of incidence and the grain diameter to determine the potential longshore sediment transport. Overall, the results suggest an inlet far away from stability since the ratio P/M_{tot} is below 20.

Table 4.4 - Bruun and Gerritsen (1960) stability criterion application to the three dredging plans

	<i>P/M_{tot}</i>	
	<i>July 2001</i>	<i>November 2001</i>
DP1	3	4
DP2	4	5
DP3	4	5

4.5. Summary

After the careful analysis regarding the previous sections, it was decided to summarize the major relative merits and drawbacks of all dredging plans. The most stable channel for DP2 and DP3 when compared to DP1 was the northern main channel. This channel remained stable from the deepest part of the lagoon until the sheet-pile curtain (APPENDIX D and APPENDIX E). However, the simulation with the DP1 shows two meanders that were formed in the northern main channel between August and November (Figure 4.6). Therefore, the design (width, bottom depth and margins slope) of the northern main channel for future dredging works must be equal with the proposed by Fortunato and Oliveira (2007a).

The southern main channel did not accrete for the DP2 and DP3 simulations, although, this channel displayed the formation of a small meander in southern main channel for DP2 due to the sediments transported in the transverse channels (Figure 4.8 and Figure 4.14). The combination of the southern and northern main channels with the higher depths when compared to DP1 provided an increase of the tidal prism and a decrease of the differences between ebb and flood durations. For that reason, these two channels must be held for future dredging works in the Óbidos lagoon.

The transverse channels addition allows to increase the sediment transport seawards by 6% when compared DP2 with DP3. Also, the transverse channels were able to decrease the flood-dominance (Figure 4.12), to reduce the tidal asymmetry and to remove sediments from the sandbank located near the southern margin (Figure 4.14 and Figure 4.15) as stated by Fortunato and Oliveira (2007a). However, the dredging of transverse channels in the southern sand bank led to an increase of the sedimentation zones near the southern main channel (junctions of the transverse channels with the main southern channel in Figure 4.14). Also, the westernmost transverse channel started to disappear after three months (APPENDIX D). On the other hand, the comparison between the final and initial bathymetries of the DP2 and DP3 shows that the northern transverse channels increased the erosion in the northern main channel when compared to DP1 (Figure 4.14 and Figure 4.15). Due to the previous reasons, three transverse channels should be dredged in the northern main channel with the configuration proposed by Fortunato and Oliveira (2007a).

The analysis on the sediment fluxes and the sediment volumes inside the Óbidos lagoon suggests an asymptotically decrease of the shoaling rate for all the dredging plans and also, an increase on the sediment transport capacity towards the lagoon for the DP2 and DP3. As a higher dredging depth increases the dredging cost, it is rationale to decrease the dredging depths of all channels (main and transverse). However, the dredging vessels need a certain water depth to be able to dredge and to avoid practical issues the bottom depths of all channels must be in accordance to Fortunato and Oliveira (2007a)

except at the inlet where the bottom depth must be +0,0 m (HZ) in order to hamper the sediment transport towards the lagoon.

Based on the application of the PA relationship to all the dredging solutions, in the case of implementing the dredging plan proposed by DHI (1997) (DP1), it is also suggested to increase the inlet cross-section by 100 m², maintaining the bottom depth and increasing the width in order to have a PA relationship closer to the proposed stability curve (Figure 4.16).

Overall, the proposed dredging plan for the Óbidos lagoon should have two main channels and three transverse channels with the configuration proposed by Fortunato and Oliveira (2007a) except the bottom depth at the inlet which must be dredged to +0,0 m (HZ). Afterwards, the application of the Escoffier (1977) diagram and the Bruun and Gerritsen (1960) criterion highlighted the fact that without a permanent solution to stabilize the Óbidos lagoon inlet, this lagoon must be dredged frequently because its shallow depths and the very energetic wave regime will often drive the inlet to closure.

5. Conclusions and future research

SUMMARY: This final chapter summarizes the conclusions achieved in chapters 3 and 4 and draws some future work directions for further research on the dissertation purpose.

The conclusions obtained in this dissertation will be divided in two groups because both chapters have relevant conclusions.

In chapter 3 the effect of the wave parameters (H_s , T_p and Dir) on the SSE in the Óbidos lagoon was investigated, and the wave blocking phenomenon was also identified. Figure 5.1 summarizes the most relevant contributions of this dissertation and sketches out a couple of future lines of research. Here, the main conclusions are enumerated:

- 1) Inside the lagoon the tidal SSE has its maximum when waves are coming from the 315° N - up to 5 cm when compared with from N or W incidences. This increase is explained by wave refraction. Waves with north-western mean wave direction are less sensitive to wave refraction and the radiation stresses in front of the inlet increase.
- 2) The wave blocking phenomenon in the Óbidos lagoon was verified. This phenomenon enhances the importance of the numerical models to take into consideration the full interaction between waves and currents, which leads to an increase of 30% on the sediment transport towards the inlet during the flood in the Óbidos lagoon.

The effect of the wave parameters on currents in the Óbidos lagoon remains an open issue, as well as the effect of locally-generated wind in the SSE inside the Óbidos lagoon and also in the sediment dynamics in the transitional channel, therefore, the wind effect must be investigated.

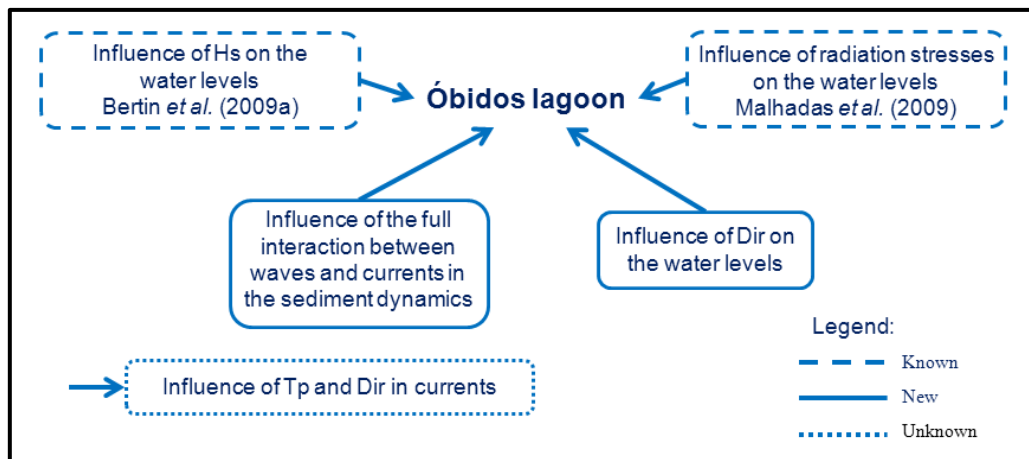


Figure 5.1 – Conclusions scheme of chapter 3

The morphodynamic evolution of the Óbidos lagoon inlet was addressed in chapter 4 for three different dredging plans: the dredging plan used in the last two decades proposed by DHI (1997) (DP1), the dredging plan proposed by Fortunato and Oliveira (2007a) (DP2) and the dredging plan similar to (DP2) but with only one transverse channel (DP3). The results obtained and future research directions are shown in Figure 5.2. The major findings are the following:

- 1) The application of the PA relationship to the different dredging plans suggest the equilibrium curve to the Óbidos lagoon has a value for C_{PA} of $2,9 \times 10^{-3}$ and 0,75 for n , and in the case of implementing the proposed dredging plan by DHI (1997), inlet cross-section should be increased by 100 m^2 to be closer to the previous proposed curve.
- 2) The higher depths of the new dredging plans (DP2 and DP3) increased the tidal prism and reduced the differences between ebb and flood durations, mainly due to the northern and southern main channels. Therefore, these channels must be included in future applications in the Óbidos lagoon with the configurations proposed by Fortunato and Oliveira (2007a) except for the bathymetric depth at the inlet which must be +0,0 m (HZ).
- 3) The transverse channels in the southern main channel increased the accretion while the transverse channels located in the northern main channel enhanced the erosion. Hence, three transverse channels must be placed in the northern main channel with the configuration proposed by Fortunato and Oliveira (2007a).
- 4) The application of the Escoffier (1977) diagram and the Bruun and Gerritsen (1960) criterion highlights the importance of permanently stabilize the Óbidos lagoon inlet to avoid the inlet's closure.

In terms of future research, the proposed dredging solution with two main channels and three transverse channels should be simulated and assessed, and the simulation period should comprehend the maritime winter. Also, new environmental friendly and reduced visual impact solutions need to be further exploited and they should have a more permanent nature. A partially submerged guiding wall with a clothoid on its end in the southern margin of the inlet could be a good engineering solution to permanently stabilize the Óbidos lagoon inlet (Fortunato and Oliveira, 2007a).

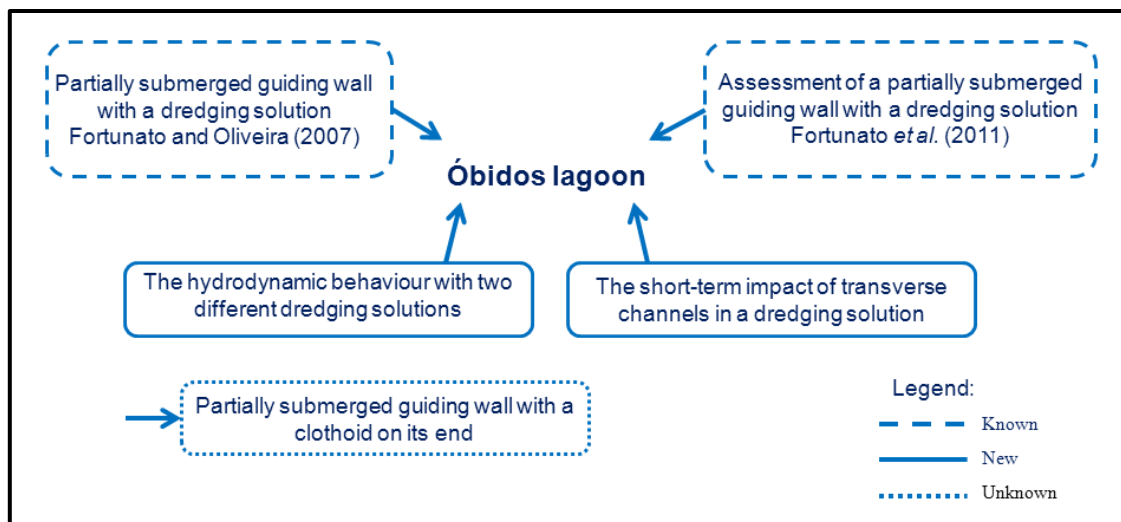


Figure 5.2 – Conclusions scheme of chapter 4

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APPENDIX A - Governing equations of the numerical models

WWM-II (Roland, 2009)

A scheme to explain the variables notation used in the numerical models is presented in Figure 1.

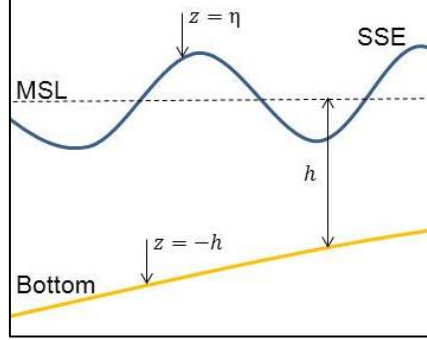


Figure 1 – Coordinates used in the numerical models

The spectral wave model WWM-II solves the wave action equation that describes the evolution of the wave spectrum in time, geographical, and spectral spaces in a quasi-steady mode. The wave action density balance equation can be written as:

$$\frac{\partial}{\partial t} N + \nabla(\dot{\mathbf{X}}N) + \frac{\partial}{\partial \sigma} (\dot{\theta}N) + \frac{\partial}{\partial \theta} (\dot{\sigma}N) = S_{tot} \quad [1]$$

$$N(t, \mathbf{X}, \sigma, \theta) = \frac{E(t, \mathbf{X}, \sigma, \theta)}{\sigma} \quad [2]$$

$$S_{tot} = S_{nl4} + S_{in} + S_{nl3} + S_{br} + S_{bf} + S_{ds} \quad [3]$$

where $N(t, \mathbf{X}, \sigma, \theta)$ is the wave action density, σ is the relative wave frequency and θ is the wave direction. S_{in} is the energy increase due to wind, S_{nl4} and S_{nl3} represent the interactions between four waves (deep waters) and three waves (shallow waters), respectively, S_{ds} and S_{br} are the energy dissipation terms due to white capping (deep waters) and the depth-induced wave breaking (shallow waters), and S_{bf} represents the loss of energy due to bottom friction. Advection velocities are obtained through the geometrical optics approximation (e.g., Keller, 1958).

$$\dot{\mathbf{X}} = \mathbf{c}_X = \frac{d\mathbf{X}}{dt} = \frac{d\omega}{dk} = \mathbf{c}_g + \mathbf{U}_{A(k)} \quad [4]$$

$$\dot{\theta} = c_\theta = \frac{1}{k} \frac{\partial \sigma}{\partial H} \frac{\partial H}{\partial m} + \mathbf{k} \cdot \frac{\partial \mathbf{U}_{A(k)}}{\partial s} \quad [5]$$

$$\dot{\sigma} = c_\sigma = \frac{\partial \sigma}{\partial H} \left(\frac{\partial H}{\partial t} + \mathbf{U}_{A(k)} \cdot \nabla_{\mathbf{X}} H \right) + c_g \mathbf{k} \frac{\partial \mathbf{U}_{A(k)}}{\partial s} \quad [6]$$

where s represents the coordinate along the wave propagation direction, and \mathbf{m} perpendicular to it. $\dot{\mathbf{X}}$ is the Cartesian coordinates vector (x, y) in the geographical space, H is the total water depth ($H = \eta + h$) obtained from the hydrodynamic model, \mathbf{k} is wave number vector (with $k = |\mathbf{k}|$), c_g is the wave group velocity, ω is the wave absolute frequency and $\nabla_{\mathbf{X}}$ is the gradient operator in geographical space. The effective advection velocity \mathbf{U}_A is approximated by the depth-averaged current (Dodet, 2013).

$$\omega = \sigma + \mathbf{k} \mathbf{U}_A \quad [7]$$

The neglected source mechanisms in this dissertation were the four-wave interactions, the white capping and the wave generation by wind forcing. The domain of this study area is small (approximately 8 to 10 km) and the wind input could be neglected. Therefore, the source terms become:

$$S_{tot} = S_{nl3} + S_{br} + S_{bf} \quad [8]$$

Information is shared between SELFE and WWM-II through the radiation stresses induced by waves. These stresses are obtained with the Longuet-Higgins and Stewart (1960) classical formulation:

$$R_{sx} = -\frac{1}{\rho_0 H} \left(\frac{\partial S_{xx}}{\partial x} + \frac{\partial S_{xy}}{\partial y} \right) \quad [9]$$

$$R_{sy} = -\frac{1}{\rho_0 H} \left(\frac{\partial S_{yy}}{\partial y} + \frac{\partial S_{yx}}{\partial x} \right) \quad [10]$$

In Equations 9 and 10, the S_{xx} , S_{yy} , S_{yx} e S_{xy} correspond to the tensor components of the radiation stresses which, for an irregular wave spectrum (Battjes, 1974) can be written as:

$$S_{xx} = \int_0^\infty \int_0^{2\pi} N(\sigma, \theta) \sigma \left[\frac{c_g(\sigma)}{c_p(\sigma)} (\cos^2 \theta + 1) - \frac{1}{2} \right] d\theta d\sigma \quad [11]$$

$$S_{yy} = \int_0^\infty \int_0^{2\pi} N(\sigma, \theta) \sigma \left[\frac{c_g(\sigma)}{c_p(\sigma)} (\sin^2 \theta + 1) - \frac{1}{2} \right] d\theta d\sigma \quad [12]$$

$$S_{xy} = \int_0^\infty \int_0^{2\pi} N(\sigma, \theta) \sigma \frac{c_g(\sigma)}{c_p(\sigma)} \sin \theta \cos \theta d\theta d\sigma \quad [13]$$

where c_p is the wave phase velocity. WWM-II solves Equation 1 with a Fractional Step Method (Yanenko, 1971). Propagation terms are solved with a Residual Distribution Scheme (Abgrall, 2006). Spectral advection is solved with the Ultimate Quickest Scheme (Leonard, 1991). A TVD third-order Runge-Kutta scheme is used for wave breaking and the bottom friction dissipation terms. Triad-wave interaction uses a dynamic method. For more details on the numerical methods the reader should consult Roland (2009). For a fully explanation on waves in coastal areas please consult Holthuijsen (2007).

SELFE (Zhang and Baptista, 2008)

The hydrodynamic model SELFE in the depth-averaged version solves the free surface elevation η and the 2D velocities (u, v) and, in a barotropic configuration, the governing equations in Cartesian coordinates are:

$$\frac{\partial \eta}{\partial t} + \nabla \cdot (\mathbf{u}H) = 0 \quad [14]$$

$$\frac{D\mathbf{u}}{Dt} = -\mathbf{g}\nabla\eta + \nabla \cdot (\mu\nabla\mathbf{u}) - \frac{\boldsymbol{\tau}_b}{\rho_0 H} - \mathbf{f}(\mathbf{v}, -\mathbf{u}) + \mathbf{R}_s \quad [15]$$

where $\nabla = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right)$ is the horizontal gradient operator, $\mathbf{u} = \mathbf{u}(x, y, t)$ is the depth-averaged horizontal velocity with Cartesian components (u, v) , η is the free surface elevation, $\frac{D}{Dt}$ is the total derivative, μ is the horizontal turbulent eddy viscosity, $\boldsymbol{\tau}_b = \boldsymbol{\tau}_b(\tau_{bx}, \tau_{by})$ is the bottom stress, f is the Coriolis factor, ρ_0 is the water specific mass and \mathbf{R}_s are the radiation stresses. To review the full derivation of the shallow water equations please consult Vieira (1984). The bottom stress is computed with a quadratic law:

$$\boldsymbol{\tau}_b = (\tau_{bx}, \tau_{by}) = \rho_0 C_D \sqrt{u^2 + v^2} \mathbf{u} \quad [16]$$

$$C_D = \frac{gn_0^2}{H^{1/3}} \quad [17]$$

where C_D is the bottom drag coefficient and n_0 is the inverse of the Manning coefficient.

SELFE solves Equations 14 and 15 simultaneously through a semi-implicit scheme, dealing with the advection term through an Eulerian-Lagrangian Method. With the Galerkin Finite-Element Method, a positive, symmetric and sparse matrix is obtained. This matrix is solved with the Jacobian Conjugate Gradient Method. Then, the verification of Equation 14 is done with a Finite Volume Method, ensuring this way the mass continuity. For more details on the numerical methods the reader should consult Vreugdenhil (1994). For a fully explanation on the shallow water flows please consult Stewart (2008).

SED2D (Dodet, 2013)

The sediment transport and bottom update model – SED2D – used the Van Rijn (2007a,b) sediment transport formula which takes into account both the bedload (Equation 18) and the suspended sediment transport (Equation 19), and is applicable to a wide range of sediments (diameters between 0,05 to 2 mm). Also, this formula considers the interaction between waves and currents.

$$q_b = 0,015uH \left(\frac{d_{50}}{H}\right)^{1,2} M_e^{1,5} \quad [18]$$

$$q_s = 0,012uH d_{50} M_e^{2,4} D_*^{-0,6} \quad [19]$$

$$M_e = \frac{u_e - u_{cr}}{\sqrt[2]{(s-1)gd_{50}}} \quad [20]$$

where M_e represents the mobility parameter (Equation 20), $u_e = u + \gamma U_w$ is the actual speed with $\gamma = 0,4$ in the case of irregular waves and $\gamma = 0,8$ for regular waves, $U_w = \pi H_s / [T_p \sinh(kH)]$ is the peak orbital velocity, H_s is the significant wave height, T_p is the peak period, $u_{cr} = \beta u_{cr,c} + (1 - \beta)u_{cr,w}$ with $\beta = u / (u + U_w)$, where $u_{cr,c}$ is the critical velocity for currents based on Shields (initiation of motion) and $u_{cr,w}$ is the critical shear velocity for waves based on Komar and Miller (1975). Critical shear velocities for waves and currents depend on the sediment median grain diameter (d_{50}) as:

If $0,05 < d_{50} < 0,5$ mm:

$$u_{cr,c} = 0,19(d_{50})^{0,1} \log(4H/d_{90}) \quad [21]$$

$$u_{cr,w} = 0,24((s-1)g)^{0,66} d_{50}^{0,33} (T_p)^{0,33} \quad [22]$$

If $0,5 < d_{50} < 2$ mm:

$$u_{cr,c} = 8,5(d_{50})^{0,6} \log(4H/d_{90}) \quad [23]$$

$$u_{cr,w} = 0,95((s-1)g)^{0,57} d_{50}^{0,43} (T_p)^{0,14} \quad [24]$$

The bottom update is done through the Exner equation, which is a sediment continuity equation (Equation 25). The Exner equation provides the bottom evolution based on the divergence of the sediment flux field:

$$\frac{\partial h}{\partial t} - \frac{1}{1-\lambda} \nabla \mathbf{q} = 0 \quad [25]$$

where $\mathbf{q} = \mathbf{q}(q_x, q_y)$ represents the depth-integrated volumetric sediment flux ($\text{m}^3\text{s}^{-1}\text{m}^{-1}$) and λ is the porosity. After integrating in time with a Finite Volume Method, Equation 25 can be solved with a node-centred method or an element-centred method. The second method estimates the bed changes with a higher resolution, but the interpolation of bathymetric changes between the nodes is not mass-conservative, and may lead to some inaccuracies (Dodet, 2013).

In this dissertation, three filters were used: an extrema filter, a slope filter and a diffusive filter.

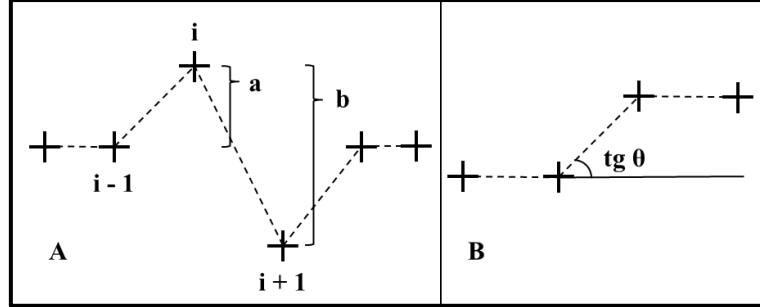


Figure 2 – A – Formulation used in the extrema filter; B – Formulation used in the slope filter

Fortunato and Oliveira (2000) implemented the extrema filter to eliminate local extremes in the bathymetry after each morphodynamic time step (Figure 2 - A). This filter is volume-conservative and introduces small numerical diffusion (Oliveira and Fortunato, 2002). This filter will compute the average in node i with the large amplitude between a and b .

The slope filter reproduces the gravity effects on steep bottoms. After each morphodynamic time step, this filter smooths the local slopes that are higher than a certain threshold. The formulation used in this filter comprehends the calculation of the angle's tangent and then, the verification with the maximum threshold specified by the user for dry or wet elements (Figure 2 - B).

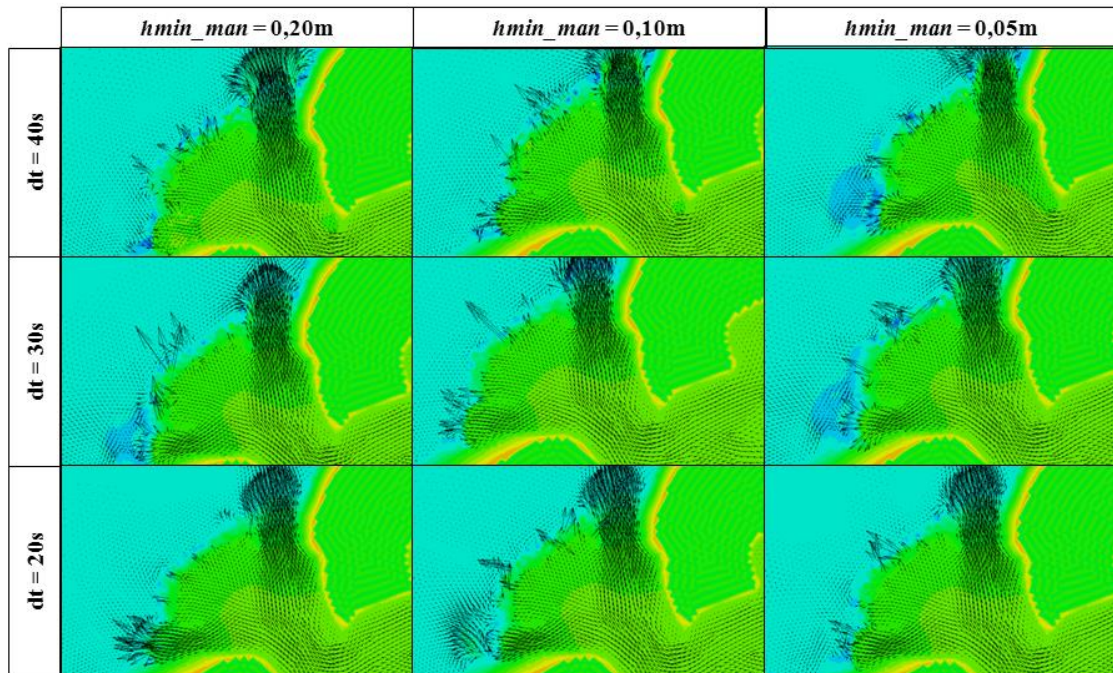
The diffusion filter aims to prevent the development of spatial oscillations similar to the filter mentioned above. This filter uses a modified expression of the diffusion-like term in the Exner equation (Equation 26) (Fortunato and Oliveira, 2007b).

$$\mathbf{Q}^* = \mathbf{Q} + \varepsilon(1 - \lambda) \left(|Q_x| \frac{\partial h}{\partial x}, |Q_y| \frac{\partial h}{\partial y} \right) \quad [26]$$

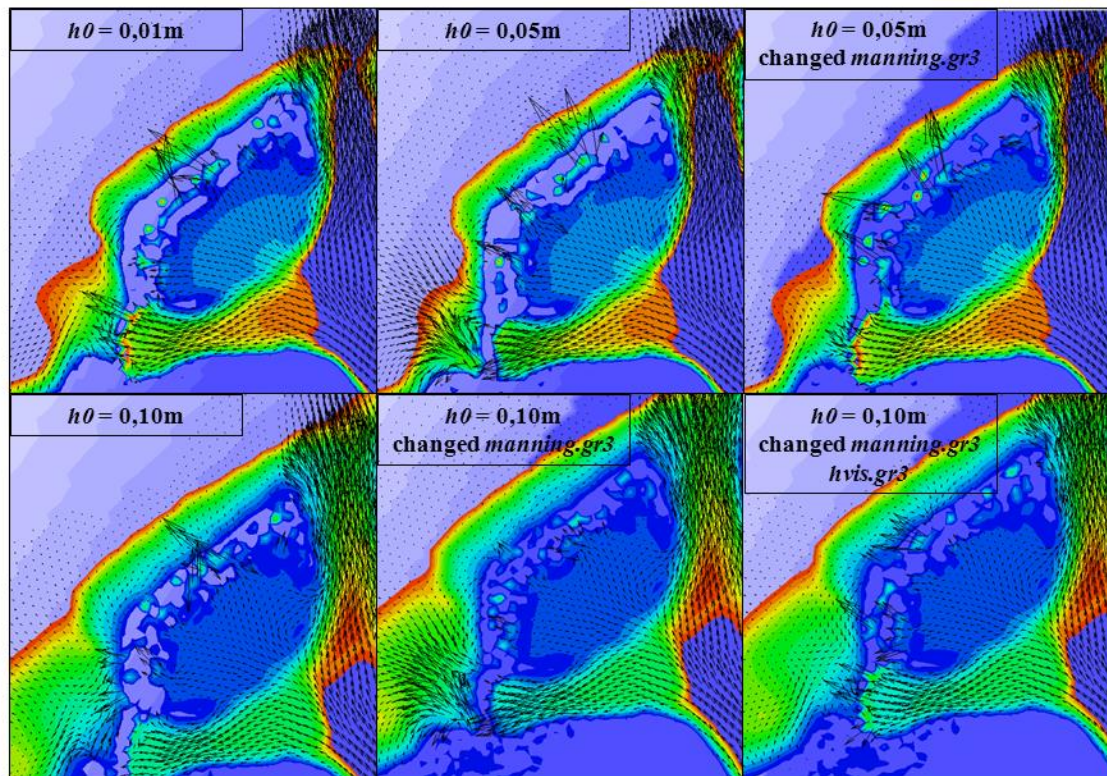
$$\mathbf{Q} = \int_{t_i}^{t_{i+1}} \mathbf{q} dt \quad [27]$$

where ε is a dimensionless coefficient, and \mathbf{Q}^* is the modified sediment flux \mathbf{q} integrated in time (Equation 27).

APPENDIX B - Velocity fields during the ebb and the changed parameters in SELFE



Velocity field vectors during the ebb at the ebb-delta for different $hmin_man$ (Manning formulation threshold depth) and Δt (geographic, color and velocity scales were disregarded because is intended to be a qualitative figure)



Velocity field vectors during the ebb at the ebb-delta for different $h0$ water depth threshold for drying), *manning.gr3* (grid of the Manning coefficient) and *hvis.gr3* (grid of the horizontal eddy viscosity) (geographic, color and velocity scales were disregarded because is intended to be a qualitative figure)

APPENDIX C - Modified boundary conditions in SED2D

In this appendix is presented the rearrangement of the Exner equation in a clear manner to introduce the modifications made in the SED2D source code. It consists of three steps: 1) integration in time the Exner equation, 2) integration over a defined control volume and 3) discretization of each side of the Equation 3.

$$\frac{\partial d}{\partial t} = \frac{1}{1-\lambda} \nabla \mathbf{q} \quad [1]$$

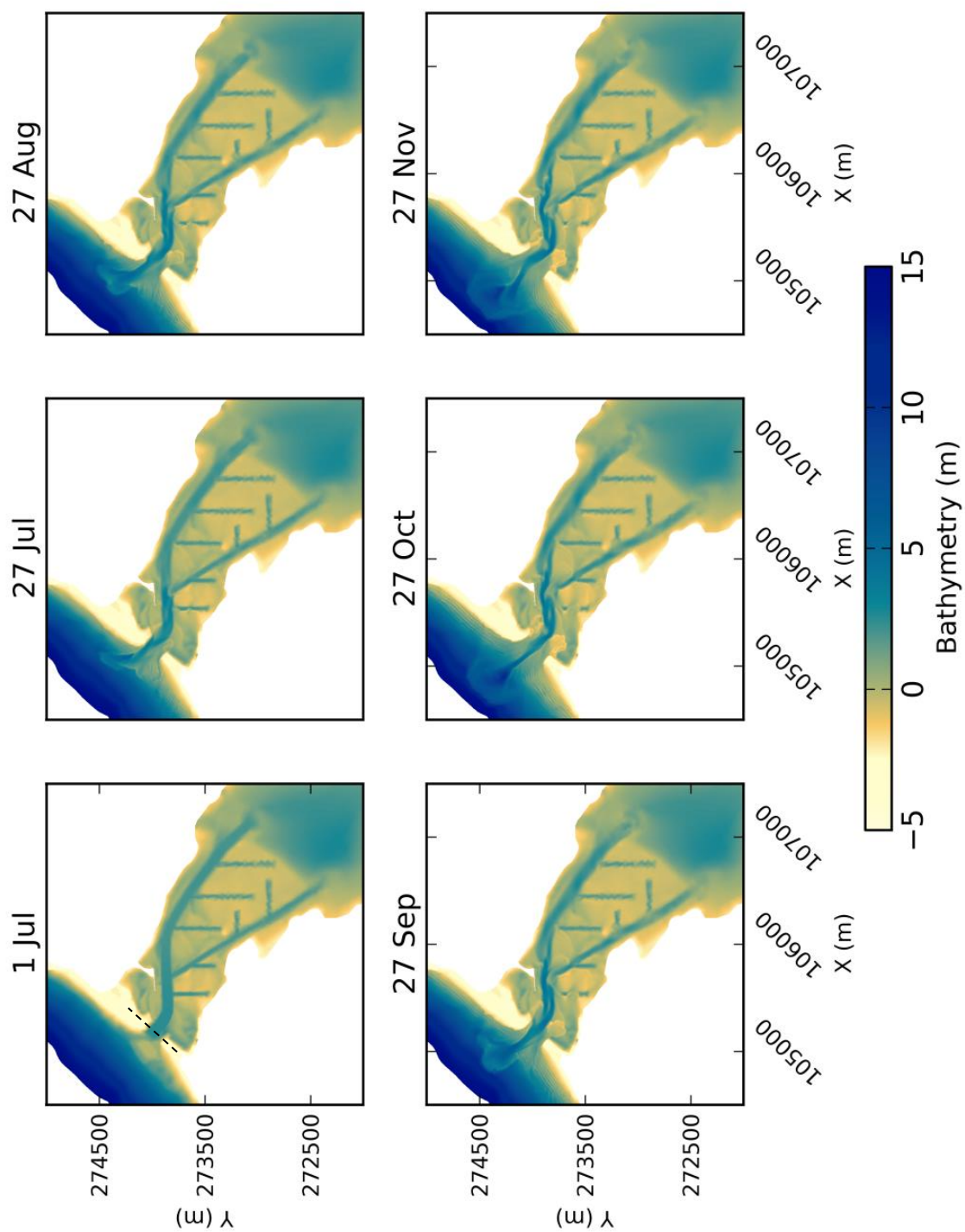
$$\Delta d^i = \frac{1}{1-\lambda} \nabla \cdot \mathbf{Q}^i \quad [2]$$

$$\int_{\Omega} \Delta d \, d\Omega = \frac{1}{1-\lambda} \int_{\Gamma} \mathbf{Q} \cdot \mathbf{n} \, d\Gamma \quad [3]$$

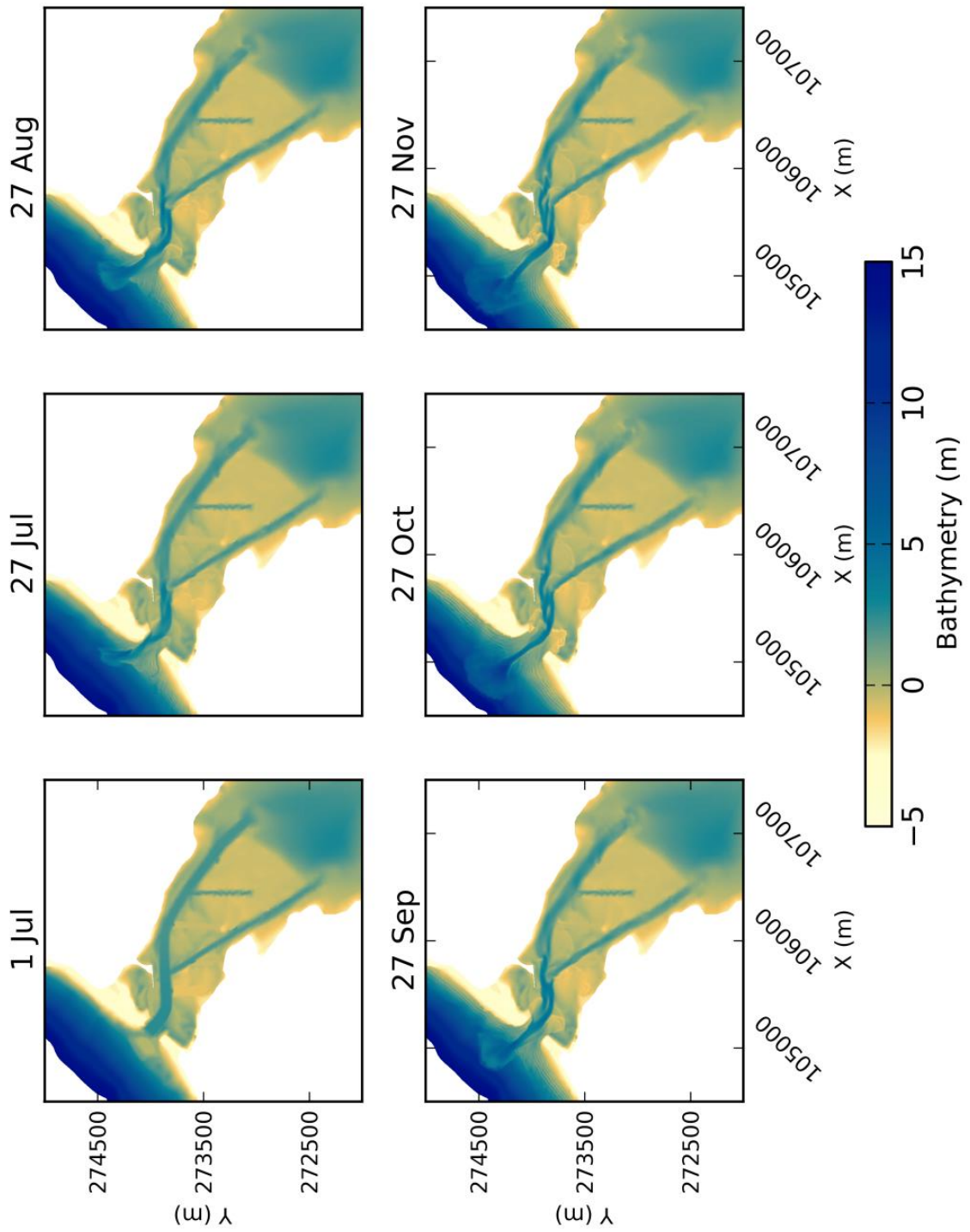
$$\sum_{i=1}^{nel} (\sum_{j=1}^3 \Delta h_{ij} C_i A_{ni}) = \frac{1}{1-\lambda} \sum_{i=1}^{nel} (\sum_{j=1}^2 Q_{ix} (y_{ij} - y_{i0}) - Q_{iy} (x_{ij} - x_{i0})) \quad [4]$$

The procedure adopted in this dissertation consisted of two steps: put the mass matrix in the Exner equation to zero ($C_i A_{ni}$ coefficients) if the row represents a boundary node and set the fluxes in SED2D to zero (Q_{ix} and Q_{iy}) if the element was on a boundary. In other words, the boundary conditions were changed by imposing a Dirichlet boundary condition in which the bathymetric depth for each boundary node remains equal to the initial bathymetry along the simulation. Doing so, the waves will break before reach the boundaries because the boundary nodes will not be updated and the morphodynamic model stays stable without spurious morphological developments.

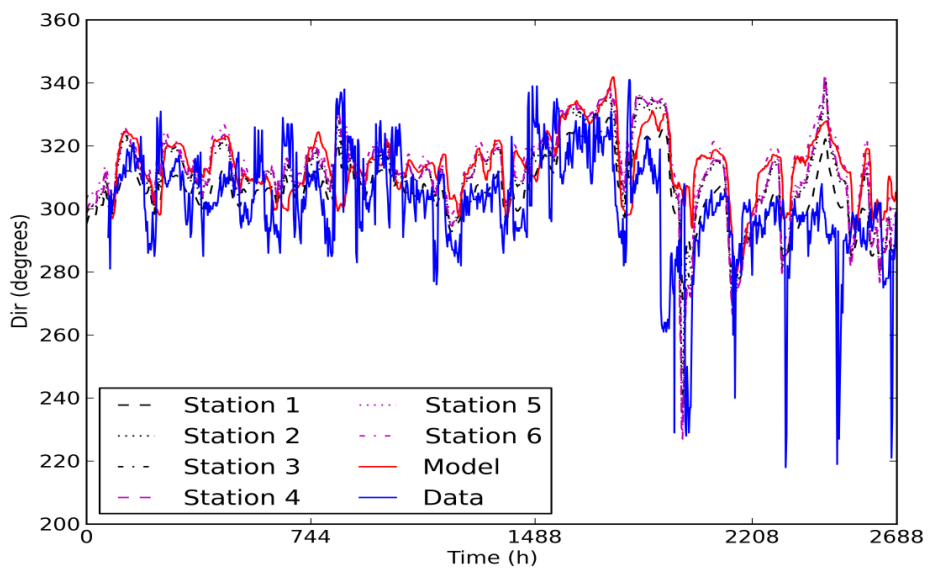
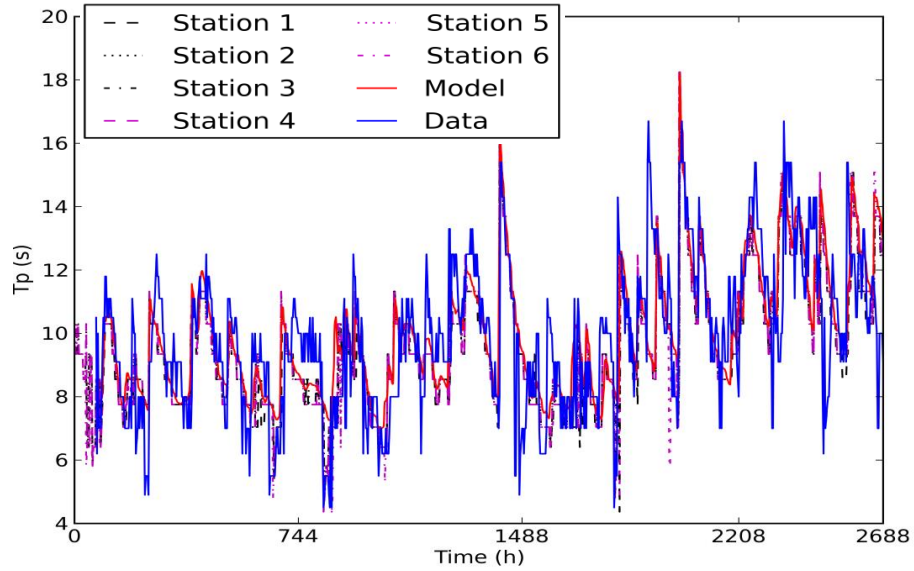
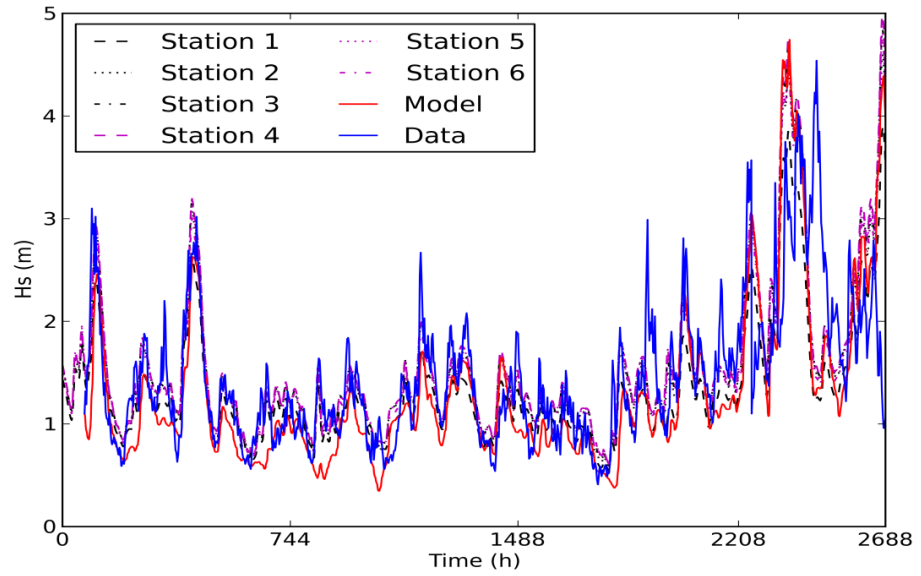
APPENDIX D - Five-month evolution of the dredging plan proposed by Fortunato and Oliveira (2007a) (DP2)



APPENDIX E - Five-month evolution of the dredging plan proposed by Fortunato and Oliveira (2007a) with only one transverse channel (DP3)



APPENDIX F - Comparison between the modelled and measured on the wave gauge and WW3 wave parameters for each WW3 station (numbering from right to left and top to bottom)



APPENDIX G – Asymptotically evolution of the bottom depth calculated for the three measured bathymetries available (July 2001, November 2001 and April 2002) at the flood delta

